

A TRIPLE PROTOSTAR SYSTEM FORMED VIA FRAGMENTATION OF A GRAVITATIONALLY UNSTABLE DISK

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Binary and multiple star systems are a frequent outcome of the star formation process^{1;2}, and as a result, almost half of all sun-like stars have at least one companion star³. Theoretical studies indicate that there are two main pathways that can operate concurrently to form binary/multiple star systems: large scale fragmentation of turbulent gas cores and filaments^{4;5} or smaller scale fragmentation of a massive protostellar disk due to gravitational instability^{6;7}. Observational evidence for turbulent fragmentation on scales of >1000 AU has recently emerged^{8;9}. Previous evidence for disk fragmentation was limited to inferences based on the separations of more-evolved pre-main sequence and protostellar multiple systems^{10;11;12;13}. The triple protostar system L1448 IRS3B is an ideal candidate to search for evidence of disk fragmentation. L1448 IRS3B is in an early phase of the star formation process, likely less than 150,000 years in age¹⁴, and all protostars in the system are separated by <200 AU. Here we report observations of dust and molecular gas emission that reveal a disk with spiral structure surrounding the three protostars. Two protostars near the center of the disk are separated by 61 AU, and a tertiary protostar is coincident with a spiral arm in the outer disk at a 183 AU separation¹³. The inferred mass of the central pair of protostellar objects is $\sim 1 M_{\odot}$, while the disk surrounding the three protostars has a total mass of $\sim 0.30 M_{\odot}$. The tertiary protostar itself has a minimum mass of $\sim 0.085 M_{\odot}$. We demonstrate that the disk around L1448

IRS3B appears susceptible to disk fragmentation at radii between 150 AU and 320 AU, overlapping with the location of the tertiary protostar. This is consistent with models for a protostellar disk that has recently undergone gravitational instability, spawning one or two companion stars.

L1448 IRS3B is located in the Perseus molecular cloud at a distance of ~ 230 pc¹⁵ and contains three protostars out of the six that collectively make up L1448 IRS3^{13;14}, spanning 0.05 pc. L1448 IRS3B is a Class 0 protostar system¹⁶, which signifies an early phase of the star formation process when the protostars are deeply enshrouded in an envelope of accreting material¹⁷. The three protostars in L1448 IRS3B (denoted -a, -b, and -c) have a hierarchical configuration; the central-most protostar, IRS3B-a, has projected separations from IRS3B-b and IRS3B-c of 61 AU and 183 AU, respectively¹³. The new observations of L1448 IRS3B conducted with the Atacama Large Millimeter/submillimeter Array (ALMA) at a resolution of $0''.27 \times 0''.16$ ($62 \text{ AU} \times 37 \text{ AU}$) provide images at 1.3 mm of the dust and gas emission surrounding the three protostars with $10\times$ higher sensitivity and $2\times$ higher resolution than previous studies.

The ALMA 1.3 mm image of L1448 IRS3B is shown in Figure 1, revealing dust emission toward each of the three distinct protostars identified in previous Karl G. Jansky Very Large Array (VLA) observations¹³. The ALMA images also reveal a disk with substructure surrounding the entire system, extending to a radius of ~ 400 AU. The disk appears to have a dominant one or two-armed spiral that links IRS3B-a and IRS3B-b with the more widely separated IRS3B-c, which is embedded in the outermost arm. The disk geometry and rotation profile (see below) place the system center of mass near the close pair, IRS3B-a and IRS3B-b, rather than toward the more distant (and brighter at 1.3 mm and 8 mm) IRS3B-c. A VLA 8 mm image is also shown in Figure 1, having comparable resolution to the ALMA data. The 8 mm emission toward the protostar positions may be a combination of both concentrated dust emission and free-free emission from the base of the jets driven by the protostars¹⁸. However, the weak 8 mm emission associated with the disk and spiral arms surrounding the three protostars is from dust only. The weaker emission from the disk at 8 mm is unsurprising, because thermal dust emission at sub-millimeter and longer wavelengths decreases in brightness with the inverse square of wavelength, modified by the dust opacity spectral index (see Methods Section 3).

The resolved disk in the ALMA observations also enables us to estimate the inclination of the system to be 45.4° , allowing us to calculate the deprojected separations of the companions (see Methods section 2). The deprojected separations of IRS3B-b and IRS3B-c from IRS3B-a are 78 AU and 254 AU, respectively. We estimate the mass of the entire disk using the dust continuum emission (with IRS3B-c removed). Assuming the dust and gas are well mixed, we find $M_{\text{disk}} = 0.30 M_\odot$ using the integrated 1.3 mm flux density of 0.52 ± 0.002 Jy. Finally, the mass of IRS3B-c in the outer disk is estimated to be $0.085 M_\odot$ from the 1.3 mm flux density of 0.19 ± 0.003 Jy; this is a lower limit to the gas mass given that IRS3B-c is likely opaque at 1.3 mm (see Methods Section 3). The surface density profile of the disk is subsequently derived from the 1.3 mm continuum data (see Methods Section 4).

The velocity structure of the disk is shown in Figure 2, using the Doppler shift of the C^{18}O molecular line emission as the gas orbits the center of mass in the system. The distinct separation between the blue and red-shifted emission is indicative of a rotation pattern centered on IRS3B-a and IRS3B-b. The rotation pattern reinforces the conclusion from

the dust continuum that IRS3B-a and IRS3B-b comprise the dominant mass in the system. Furthermore, the overall mass of IRS3B-c must be low relative to IRS3B-a and IRS3B-b, despite it being the brightest source at 1.3 mm and 8 mm. We estimate that IRS3B-a and IRS3B-b have a combined mass of $\sim 1.0 M_{\odot}$ from the disk rotation profile (see Methods Section 5). We are unable to make an estimate of the mass for IRS3B-c from molecular line kinematics, but the $0.085 M_{\odot}$ from the dust emission can be considered a lower limit. Note that the C^{18}O emission does not trace the entire disk structure seen in the dust because any emission at velocities within $\pm \sim 1 \text{ km s}^{-1}$ of the system velocity ($\sim 4.5 \text{ km s}^{-1}$) is filtered-out along with the large, surrounding envelope.

With these physical parameters in hand, we can now assess whether L1448 IRS3B hosts a gravitationally unstable disk that recently fragmented to form IRS3B-c. Gravitational instabilities arise in protostellar disks when the self-gravity of the disk becomes significant compared to the gravity of the central protostar(s). While disk fragmentation is the non-linear outcome of a complex hydrodynamic instability, the canonical analytic estimate for when instability occurs is that the Toomre Q parameter¹⁹ is of order unity (see the Methods section 6 for details). For Keplerian disks, Toomre’s Q can be expressed as:

$$(1) \quad Q \approx 2 \frac{M_* H}{M_d r}$$

where $H = c_s/\Omega$ is the disk scale height, c_s the disk sound speed, and Ω the Keplerian angular velocity. Using the masses derived above, we can immediately see that for typical disk temperatures of 30 K ($c_s = 0.3 \text{ km s}^{-1}$) at a radius of 200 AU ($\Omega \sim 7 \times 10^{-11} \text{ s}^{-1}$), $Q \approx 1$. While suggestive of instability, a more robust model is desirable. In the Methods Section 6, we describe a simple analytic disk model²⁰ illustrating that the IRS3B disk likely underwent gravitational instability in the recent past. Our analytic model is consistent with many previous numerical simulations of fragmenting disks that replicate this type of disk morphology (e.g., spiral arms) for parameters comparable to the observed quantities^{21;22}.

We show the result of our analysis in Figure 3, where we plot the range of possible Q values for different inferred mass accretion rates, as a function of disk radius. Given the mass of the central protostars (both assumed to be $\sim 0.5 M_{\odot}$) and the luminosity of $3.2 L_{\odot}$ ²³, a current mass accretion rate of $\sim 10^{-7} M_{\odot} \text{ yr}^{-1}$ is most consistent with the model. For a wide range of accretion rates, the disk around IRS3B is marginally unstable ($Q \approx 1$) at radii between 150 AU and 320 AU. This conclusion is consistent with both the continued presence of the spiral arms and IRS3B-c forming near its current deprojected separation of 254 AU. The evolutionary phase of IRS3B-c is also consistent with it forming within the last few dynamical times ($\sim 4000 \text{ yr}$ assuming Keplerian rotation around IRS3B-a/b). Prior to the formation of IRS3B-c, the outer disk is likely to have been more unstable, and the subsequent accretion of disk mass onto IRS3B-c is now stabilizing the disk. A fragment formed by instability in this disk would likely begin with a mass of roughly $10^{-2} M_{\odot}$, consistent with IRS3B-c undergoing accretion to arrive at our measured mass of at least $0.085 M_{\odot}$ (see Methods Section 6).

In addition to IRS3B-c, IRS3B-b may have formed from an earlier episode of disk fragmentation²¹. IRS3B-b could have formed near its current location, if the disk was more unstable in the past. It could also have formed in the outer disk prior to IRS3B-c and migrated inward to its current location over a few 10s of orbital timescales²⁴. Fragments formed early in the evolution of the disk can grow to comparable mass as the primary star²¹. Finally, we note that if all three protostars were on roughly circular, co-planar, Keplerian

orbits with semi-major axes equal to their deprojected separations, they would be marginally dynamically unstable²⁵. Thus we expect the orbits to evolve on rapid timescales (with respect to the expected stellar lifetime), especially as the disk dissipates. A natural outcome of this dynamical instability is the formation of a more hierarchical system with a tighter (few AU) inner pair and wider (100s to 1,000s AU) tertiary²⁶, consistent with observed triple systems²⁷.

Our results demonstrate that protostellar disks can experience gravitational instability at young ages, leading to the formation of hierarchical multiple systems. Such disks may also show spiral structure associated with the linear growth phase of gravitational instability. The spiral structure in L1448 IRS3B is consistent with simulated ALMA observations of unstable disks that assume a comparable ratio of disk mass to stellar mass²⁸. Spiral structure has been detected previously in more-evolved disks from observations of near-infrared scattered light originating from the disk upper layers²⁹. However, those disks are not sufficiently massive to generate the spiral structure by gravitational instability, nor is the spiral structure apparent in observations of the disk midplanes. All these cases are more consistent with spiral arms triggered by a low-mass stellar, brown dwarf, or planetary mass companion³⁰.

Although the frequency of fragmenting disks like IRS3B is still unknown, we do not expect them to be intrinsically rare. Our conclusion is based on the typical multiplicity and separation distributions observed in the parent star forming region, Perseus. Close multiples such as the IRS3-a/b or IRS3B-a/c pairs are frequently found in protostar systems^{13;10;12}. This discovery, together with the recent work supporting turbulent fragmentation at larger scales (1,000s of AU)^{8;9}, provide direct observational support for the two dominant theoretical models for binary/multiple star formation. We predict that systems similar to IRS3B will be discovered in larger samples with facilities like ALMA, and their frequency will help benchmark the relative contribution of disk fragmentation to the population of binary/multiple stars.

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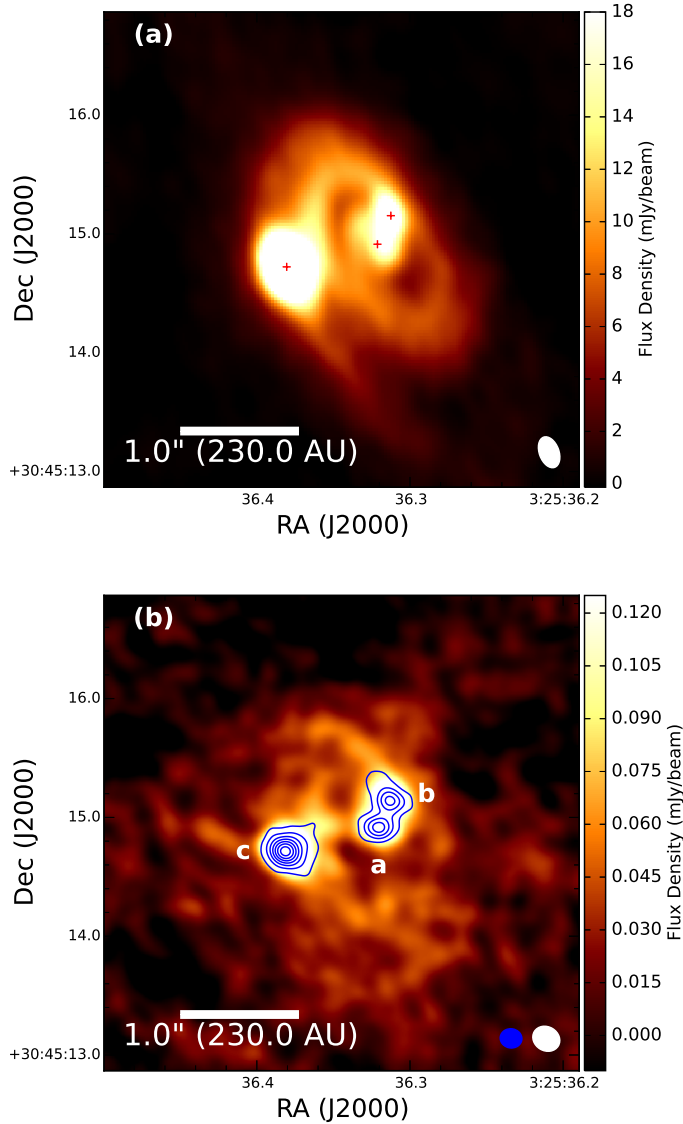


FIGURE 1. ALMA and VLA images of the disk and triple protostar system L1448 IRS3B. (a) ALMA 1.3 mm image of the extended disk, showing an evident bright source on the left (IRS3B-c) in the outer disk and another blended source on the right near the disk center (IRS3B-a and IRS3B-b). (b) VLA 8 mm image smoothed to a similar resolution as the ALMA image, capturing some of the faint, extended disk at longer wavelengths. The contours in panel (b) are from a higher-resolution VLA 8 mm image¹³ clearly showing the individual protostars with corresponding designations. All three protostars are embedded within apparent spiral arms that emerge from IRS3B-a/IRS3B-b and extend to IRS3B-c in the outer disk. The positions of the three protostars identified from the VLA data are shown by red crosses in panel (a). The contours start at and increase with 5σ , where $\sigma = 0.009 \text{ mJy beam}^{-1}$. The resolution of each image is shown with an ellipse(s) drawn in the lower right corner, corresponding to $0''.27 \times 0''.16$ (62 AU \times 37 AU) for the ALMA image in panel (a), $0''.24 \times 0''.20$ (55 AU \times 46 AU) for the VLA image in panel (b), and $0''.18 \times 0''.16$ (41 AU \times 37 AU; blue ellipse) for the contour image in panel (b).

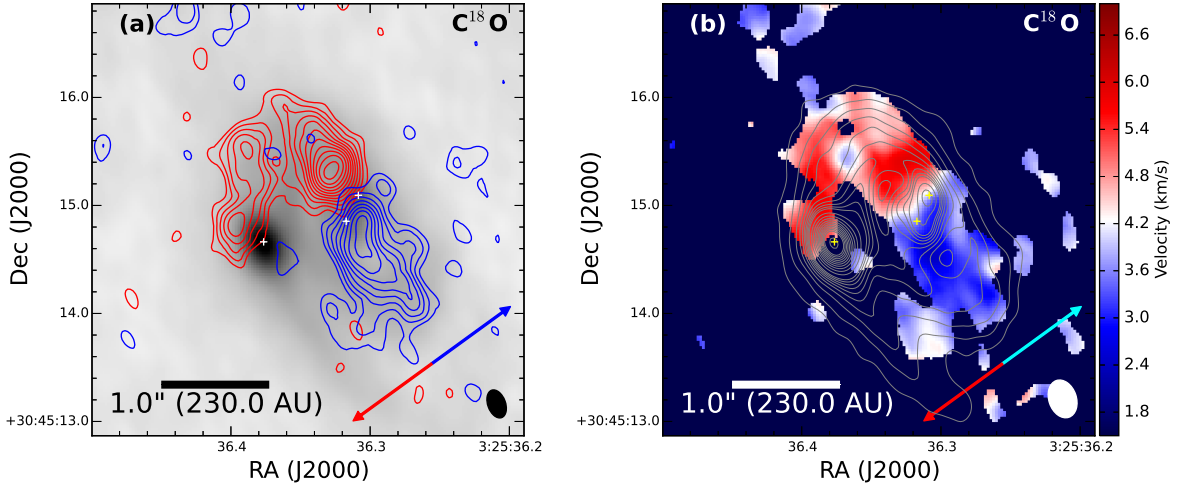


FIGURE 2. Images of the $C^{18}O$ emission and its corresponding velocity maps from the disk around L1448 IRS3B showing a rotation signature. (a) Red and blue-shifted $C^{18}O$ ($J = 2 \rightarrow 1$) emission are overlaid on the ALMA 1.3 mm continuum image (grayscale) as red and blue contours. (b) Line-center velocity map of the $C^{18}O$ emission with 1.3 mm continuum contours overlaid in gray. The $C^{18}O$ traces higher-velocity emission near IRS3B-a and IRS3B-b, consistent with the system center of mass being near these protostars. Lower velocity gas is found to be associated with the outer spiral arm detected in dust emission. The molecular line emission does not fully trace the disk due to spatial filtering of emission with velocities close to that of the system ($\sim 4.5 \text{ km s}^{-1}$). The source positions are marked with white or yellow crosses. The outflow direction¹⁴ is denoted by the blue and red arrows. The angular resolution of these data is given by the ellipse in the lower right corners, $0''.36 \times 0''.25$ ($83 \text{ AU} \times 58 \text{ AU}$). The contours in panel (a) start at 4σ and increase in 1σ intervals. The $C^{18}O$ emission was integrated over $1.25 - 4.0 \text{ km s}^{-1}$ and $5.5 - 7.0 \text{ km s}^{-1}$ for the blue and red-shifted maps, respectively. The noise levels for $C^{18}O$ are $\sigma_{Blue} = 2.25 \text{ K km s}^{-1}$ and $\sigma_{Red} = 1.65 \text{ K km s}^{-1}$. The continuum (gray) contours in panel (b) start at and increase by 10σ , at 100σ the levels increase in steps of 30σ , and at 400σ the levels increase by 100σ steps; $\sigma = 0.14 \text{ mJy beam}^{-1}$.

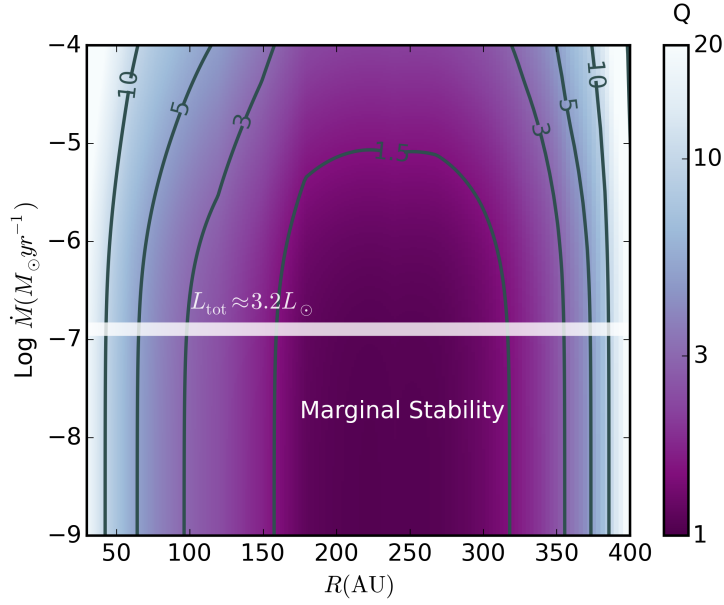


FIGURE 3. Plot of Toomre Q (instability criterion; $Q \approx 1$ or less being unstable), versus mass accretion rate and radius. The y-axis is the mass accretion rate in the disk and the gray contours are Q at a given radius and accretion rate. The luminosity and estimated $1.0 M_{\odot}$ mass of the inner pair indicates a mass accretion rate of $\sim 10^{-7} M_{\odot} \text{ yr}^{-1}$ (thick horizontal line). Q approaches ~ 1 at radii between 150–300 AU indicating that the outer disk is marginally unstable. At radii > 350 AU the surface density of the disk declines, such that it becomes stable again. However, the outer radius of ~ 400 AU, (800 AU ($3''5$) diameter), is comparable to the spatial scale at which our observations no longer recover all emission.

METHODS

1. OBSERVATIONS

L1448 IRS3B was observed with the Atacama Large Millimeter/submillimeter Array (ALMA) during Cycle 2 on 27 September 2015 with 33 antennas sampling baselines between 32 - 2000 meters. The observations were executed within a 1.8 hour block and the total time spent on L1448 IRS3B was ~ 2.9 minutes. The precipitable water vapor was ~ 0.7 mm throughout the observing session. The phase calibrator was J0319+4130 (3C84), the bandpass calibrator was J0237+2848, and the amplitude and absolute flux calibrator was the monitored quasar J0238+166. The correlator was configured to observe a 2 GHz continuum band centered at 232.5 GHz, and with 60 MHz bands centered on the following molecular transitions: ^{12}CO ($J = 2 \rightarrow 1$), ^{13}CO ($J = 2 \rightarrow 1$), C^{18}O ($J = 2 \rightarrow 1$), SO ($J_N = 6_5 \rightarrow 5_4$), and H_2CO ($J = 3_{03} \rightarrow 2_{02}$). The flux calibration accuracy is expected to be better than 10%.

The raw data were manually reduced by the North American ALMA Regional Center staff using CASA³¹ version 4.5.0. We performed self-calibration on the continuum data to increase the signal to noise ratio by correcting for short timescale phase and amplitude fluctuations. The phase and amplitude solutions from the continuum self-calibration were also applied to the spectral line bands. The data were imaged using the *clean* task within CASA 4.5.0; the ALMA images shown in Figure 1 were generated using Briggs weighting with a robust parameter of 0.5. Within the *clean* task, we only include data at uv -distances > 50 k λ to mitigate striping in the images from large-scale emission detected on the shortest baseline that could not be properly imaged. The molecular line images shown in Figure 2 (main text) and Extended Data Figures 1 and 2 were imaged with natural weighting, tapering at 500 k λ , and with data having uv -distances > 50 k λ . Tapering reduces the weight of longer baseline data in the deconvolution to facilitate the detection of larger, lower surface brightness structures. The beam size of the continuum image is $0''.27 \times 0''.16$ (62 AU \times 37 AU) and the beam size of the molecular line data is $0''.36 \times 0''.25$ (83 AU \times 58 AU). The resultant noise in the 1.3 mm continuum was 0.14 mJy/beam and 15 mJy/beam in 0.25 km s $^{-1}$ channels for the spectral line observations.

We show the ^{12}CO ($J = 2 \rightarrow 1$) and H_2CO ($J = 3_{03} \rightarrow 2_{02}$) integrated intensity maps in Extended Data Figure 1. The ^{12}CO map clearly shows an outflow(s) originating from the region of the three protostars. The red-shifted emission shows a clear, wide outflow cavity, with another red-shifted feature within it. We suggest that this secondary red-shifted feature within the outflow cavity is a jet-like outflow from IRS3B-c and that the arc-like shape is due to the orbital motion of the source. The blue-shifted outflow is much more diffuse and not as well-recovered in our data, but there appears to be emission associated with all three protostars. Nonetheless, it is clear that the outflow map toward these three sources is quite complex and made even more complex by the two more widely separated systems (IRS3A and IRS3C)³² that also drive outflows that are extending across this map. The larger-scale outflow has been examined more thoroughly at lower resolution¹⁴.

We also detect faint H_2CO emission from the vicinity of L1448 IRS3B, shown in Extended Data Figure 1. The emission has low intensity, but also shows a velocity gradient in the same direction as ^{13}CO and C^{18}O . SO ($J_N = 6_5 \rightarrow 5_4$) emission was not detected toward L1448 IRS3B. H_2CO is expected to trace the inner envelope and disk around the protostars³³ and it is also sometimes present in outflows³⁴. SO has been found to trace both outflow shocks³⁵

and possibly a centrifugal barrier, highlighting a transition between the disk and infalling envelope of protostar^{36;37}.

The ^{13}CO ($J = 2 \rightarrow 1$) emission traces kinematics similar to that of the higher-velocity C^{18}O , with the blue and red-shifted emission concentrated around IRS3B-a and IRS3B-b. The integrated intensity maps of the red and blue-shifted emission are shown overlaid on the ALMA 1.3 mm continuum image in Extended Data Figure 2, along with the velocity map derived from the ^{13}CO emission. The ^{13}CO does not extend toward IRS3B-c because of spatial filtering; ^{13}CO is more severely affected by spatial filtering because it is more abundant and has higher opacity than C^{18}O , leading to a wider velocity range that is filtered-out. There is some outflow emission detected in the ^{13}CO , but the apparent rotation of the inner disk(s) around IRS3B-a and IRS3B-b dominates the velocity field for the detected ^{13}CO . Previous studies have also found that ^{13}CO can be a good tracer of the disk kinematics with sufficient spatial resolution to rule-out outflow contamination^{12;38;39}.

2. SUBTRACTING THE TERTIARY, IRS3B-C

To characterize the circumtriple disk, it is necessary to remove the bright protostar IRS3B-c located in the outer arm of the disk. IRS3B-c has the largest peak intensity in the system and its presence potentially masks some of the underlying disk structure, in addition to skewing the derivation of system geometry (e.g., inclination, position angle). Furthermore, IRS3B-c appears internally heated and is likely optically thick and the process of removing it from the data enables us to characterize its properties, independently of the disk that it is embedded within.

To remove IRS3B-c as cleanly as possible, we fit a two-component Gaussian with a constant zero-level offset. The two-component Gaussian is comprised of a point-source component superposed with a more extended Gaussian, and the zero-level offset is used to preserve emission from the disk in the vicinity of IRS3B-c source. The fitting region is also restricted to a $0''.7 \times 0''.6$ ellipse around IRS3B-c, ensuring that the Gaussian fit to IRS3B-c is not strongly affected by the emission in the surrounding, extended disk.

We used the *imfit* task in CASA 4.5.0 to fit the Gaussian components. We then used these components to construct a model image of IRS3B-c within CASA. Next, we used the *setjy* task to generate visibility data from the model image via a Fourier transform, and then filled the model column of the CASA measurement set with the visibility data. Finally, we use the *uvsb* task to subtract the model column from the data column, producing a dataset with IRS3B-c subtracted. Images with IRS3B-c removed were then generated with *clean*, and we show the resultant images in Extended Data Figure 3.

With IRS3B-c removed, we fit the major and minor axes of the disk to estimate the inclination with respect to the line of sight and position angle. To do this, we first generated an image from the dataset with IRS3B-c subtracted using the CASA *clean* task, but taper (e.g., smooth, see Section 1) the visibility data to have a resolution of $0''.43 \times 0''.36$; this image smooths over the spiral structure such that a single Gaussian can better fit the disk emission from the entire source. We use the *imfit* task in CASA to fit a single Gaussian, with the center position fixed at the location of IRS3B-a, which appears to be the most centrally-located source in the system.

The deconvolved major and minor axes of the resulting Gaussian were $1''.78$ and $1''.25$ ($409 \text{ AU} \times 288 \text{ AU}$). Assuming symmetry, this corresponds to an inclination angle of $45.4^{\circ}_{-12.8}^{+10.4}$ ($\arccos(\theta_b/\theta_a)$) and the position angle is 29.5° east of north. The uncertainty in the inclination results from the disk being asymmetric, as shown by the Figure 1 in the main text and Extended Data Figure 3 (see Section 4). We estimate that there could be a systematic uncertainty of 20% in the minor axis of the disk because of the disk asymmetry.

3. DISK MASS

The dust emission from L1448 IRS3B can be used to estimate the mass of the surrounding disk and IRS3B-c itself. Assuming that the dust around L1448 IRS3B is optically thin, isothermal, and that the gas and dust are well-mixed, we can calculate the mass with the equation

$$(2) \quad M_{dust} = \frac{D^2 F_{\lambda}}{\kappa_{\lambda} B_{\lambda}(T_{dust})}.$$

F_{λ} is the integrated flux density at 1.3 mm, B_{λ} is the Planck function, and $D = 230 \text{ pc}^{15}$ is the distance to the Perseus molecular cloud. We adopt $T_{dust} = 30 \text{ K}$ as a typical dust temperature at a radius of 100 AU, derived from radiative transfer models^{40;41}. The disk does have a temperature gradient, but 30 K is a reasonable average value for the outer disk where most of the mass resides. The remaining term κ_{λ} is the dust opacity at the observed wavelength, which is adopted from dust opacity models. While there are a multitude of possible models to consider⁴²; we adopt two limiting cases for simplicity. The first value is $\kappa_{1.3mm} = 0.899 \text{ cm}^2 \text{ g}^{-1}$ ^[43], appropriate for dense cores in molecular clouds like L1448 IRS3B. The second value we adopt is $\kappa_{1.3mm} = 2.3 \text{ cm}^2 \text{ g}^{-1}$, which is typically used for proto-planetary disks^{44;45}. The second value is derived from a parametrization of the dust opacity to have a value of $10.0 \text{ cm}^2 \text{ g}^{-1}$ at $\lambda = 0.3 \text{ mm}$ and $\kappa_{\lambda} = 10.0 \times (\lambda/0.3 \text{ mm})^{-\beta}$, where $\beta = 1$ ^[46;44;45]. We then adopt a standard dust to gas mass ratio of 1:100^[47] and assume that this ratio is constant throughout the disk. We note that the dust to gas ratio may be changing in more-evolved proto-planetary disks^{48;49;45}, but we assume the constant, canonical value for this much younger system. Additionally, young protostellar disks may not be optically thin at 1.3 mm, and thus our mass estimates using Equation 2 are likely lower limits. Note that a higher disk mass would only strengthen our conclusion that the disk is unstable. We show below that analytic models of the disk suggest that the inner parts are marginally optically thick at somewhat shorter wavelengths.

Keeping the assumptions and caveats from the previous paragraph in mind, we can now calculate the masses of the disk surrounding the three protostars. We calculate a total mass of $0.39 M_{\odot}$ for the disk surrounding L1448 IRS 3B, assuming $\kappa_{1.3mm} = 0.899 \text{ cm}^2 \text{ g}^{-1}$. Of the total, $0.085 M_{\odot}$ is contained within the concentrated dust around IRS3B-c; note that we adopt $T_{dust} = 40 \text{ K}$ for IRS3B-c since the region directly around it should be warmer and the peak brightness temperature (T_B) in the data is $\sim 40 \text{ K}$ toward this source. This high peak T_B is indicative of increased opacity in IRS3B-c because $T_B = T$ when the emission source is optically thick. Thus, the mass estimate of $0.085 M_{\odot}$ is also likely to be a lower limit because of optical depth effects. If we instead adopt $\kappa_{1.3mm} = 2.3 \text{ cm}^2 \text{ g}^{-1}$, assuming that the dust has already grown substantially from ISM size as in older proto-planetary disks, we find a total mass of $0.15 M_{\odot}$ (disk plus IRS3B-c) and $0.03 M_{\odot}$ for IRS3B-c alone.

To conclude, the higher mass estimates using $\kappa_{1.3mm} = 0.899 \text{ cm}^2 \text{ g}^{-1}$ are favored because that opacity model is meant to reflect the dust typically found in dense molecular clouds. Since L1448 IRS3B is still in a very early stage of the star formation process, the bulk of its dust should reflect the molecular cloud content rather than the dust in proto-planetary disks that has been able to undergo significant dust evolution^{50,42}.

4. DISK SURFACE DENSITY

The asymmetry and complexity of the dust continuum around L1448 IRS3B prohibit us from directly fitting symmetric disk models to the data to determine the underlying density structure. However, given that the disk is well-resolved, we can empirically determine the azimuthally averaged surface density profile as a function of radius from the data. We first deprojected the visibility data using the inclination of 45.4° and position angle of 29.5° , determined from the Gaussian fitting in Section 2. The uv-data are deprojected using standard methods⁵¹ that adjust the u and v coordinates based on the geometric projection. We then generate deprojected images using the *CASA clean* task in the same manner as we made images for the non-deprojected data (see Section 1). The deprojected image is shown in Extended Data Figure 3.

With the deprojected image, we measure the flux density in a series of circular annuli, centered on IRS3B-a; each annulus has a width of half the beam-size. Afterwards, we divide the flux density in each annulus by its surface area and converted the surface brightness to a surface density using the methodology described in Section 3, assuming a radial temperature profile that is described in Section 6. The surface density (Σ) inferred from the observed intensity profile, and the assumed temperature profile are shown in Extended Data Figure 4.

5. PROTOSTAR MASSES

It is essential to know the masses of the central protostars (IRS3B-a and IRS3B-b) to quantify the stability of the disk. The angular resolution of the observations is not sufficient to determine the masses of each component of the central pair, and our measurement will be limited to the combined mass of both components. Figure 2 of the main text shows the integrated intensity and velocity maps of the C^{18}O ($J = 2 \rightarrow 1$) emission and the ^{13}CO ($J = 2 \rightarrow 1$) is shown in Extended Data Figure 2. The maps show clear evidence of rotation, centered on IRS3B-a and IRS3B-b, based on the spatial separation of the blue and red-shifted components. Furthermore, the C^{18}O emission is also tracing some of the structure observed in dust emission in the outer disk.

To examine the molecular line kinematics, we extract position-velocity (PV) diagrams along the major axis of the disk, centered on IRS3B-a and IRS3B-b. The PV diagram for C^{18}O is shown in Extended Data Figure 5. The signal-to-noise ratio of the data is not high enough to employ the techniques that have been used to quantitatively measure protostellar masses in other works^{38;52;53;34;54;37}. Nevertheless, we can estimate the protostar mass by drawing a Keplerian curve on the PV diagram, accounting for the 45.4° system inclination. We find that the combined mass for both IRS3B-a and IRS3B-b is likely $\sim 1.0 M_\odot$. We also compare this estimate to a PV diagram of a rotating thin disk model⁵⁵ with the same central mass and inclination, shown in Extended Data Figure 4. The thin disk model occupies a very

similar domain in PV space as the data. Notice that the line we draw in the PV diagrams does not go through the middle of the data, but instead traces the edge of the highest velocity emission for a given position. This is because observations of a Keplerian disk will measure multiple velocities superposed at a given position and the Keplerian velocity at a given radius only corresponds to the highest velocities⁵³. Thus, the total protostellar mass cannot be substantially larger than $1.0 M_{\odot}$ without being inconsistent with the data.

6. STABILITY ANALYSIS

The simplest estimate for when disks are subject to gravitational instability is provided by the Toomre Parameter¹⁹:

$$(3) \quad Q = \frac{c_s \kappa}{\pi G \Sigma},$$

where c_s is the sound speed within the disk, κ is the epicyclic frequency (which is Ω , the angular rotation frequency, for Keplerian disks), Σ is the disk surface density, and G is the gravitational constant. The quantities c_s , Ω , and Σ are all functions of radius within a disk. If Q reaches order unity, the disk becomes susceptible to the growth of spiral density waves. Disk fragmentation due to gravitational instability is thought to be the non-linear outcome of this hydrodynamic instability²⁰. The evolution of disks subject to the instability, and in particular their ability to create long-lived, bound objects, is a function of many parameters, including infall from the environment and radiative heating and cooling. The importance of the latter effect can be reasonably well encapsulated by the disk cooling time:

$$(4) \quad t_{\text{cool}} = \beta \Omega^{-1} = \frac{4}{9\gamma(\gamma-1)} \frac{\Sigma c_s^2}{\sigma T^4} f(\tau).$$

where

$$(5) \quad f(\tau) = \tau + \frac{1}{\tau},$$

where the optical depth $\tau = \kappa \Sigma / 2$ is calculated using the Rosseland mean opacity. $f(\tau)$ reasonably captures how the accretion energy diffuses from the midplane in the optically thick and thin regimes⁵⁶.

For fragments to survive, the gas must be able to contract by radiating away heat generated by shocks on roughly the orbital timescale⁵⁷. β is often used in the literature to parametrize cooling via a single dimensionless number to probe disk behavior in different regimes. Although there is some debate about the critical value of β below which fragmentation successfully produces bound objects^{58;59;60}, fragment formation is robust for $\beta < 10$.

To assess whether IRS3B is a good candidate for disk fragmentation, we generate a set of one-dimensional analytic disk models. Specifically, we calculate the Toomre parameter and cooling time as a function of disk radius. We use the surface density profile derived for the disk after the removal of IRS3B-c (see Section 4), and treat IRS3B-a and IRS3B-b as equal mass $0.5 M_{\odot}$ protostars (see Section 5). We use the youngest age available from the Baraffe et al. models⁶¹ (1 Myr) to calculate the stellar radii and intrinsic luminosity. To estimate the disk temperatures, a two component temperature model is used^{62;21}. The disk midplane temperature is found by balancing heating and cooling terms:

$$(6) \quad \sigma T^4 = \frac{3}{8} f(\tau) F_{\text{acc}} + F_{*}.$$

Here F_* represents the energy flux associated with irradiation from the protostars,

$$(7) \quad F_* = \sigma T_*^4 = f \frac{L_*}{4\pi r^2}$$

where $f \approx 0.1$ is based on ray tracing calculations that show that much of the stellar flux impacts the infalling envelope⁶³ and is reprocessed back down onto the disk in embedded sources⁶⁴. L_* includes both the model stellar luminosity and the accretion luminosity, which are comparable to each other. F_{acc} is the energy flux due to dissipation of accretion energy within the disk:

$$(8) \quad F_{\text{acc}} = \frac{3}{8\pi} \dot{M} \Omega^2$$

We consider two models for disk opacity: a temperature independent opacity of $\kappa = 0.24 \text{ cm}^2 \text{ g}^{-1}$ [⁶⁵], and an interpolated model, where roughly $\kappa \propto T^2$ [⁶⁶]. The former model is more consistent with grain growth beyond mm sizes, while the latter should capture more ISM-like grains. For the disk models presented here the opacity has negligible effect on the disk temperature because it is dominated by stellar irradiation, which sets the disk surface temperature, and therefore is unaffected by optical depth. This is typical of unstable protostellar disks²⁰. When the midplane heating is relatively weak, the disk becomes nearly vertically isothermal⁶⁷. Note that the numerical values for these two models are distinct from those used to derive the disk mass in Section 2 because calculation of equilibrium temperatures analytically requires a mean opacity, rather than the 1.3 mm opacity.

We first evaluate the disk properties using temperature independent opacities, and explore the impact of a range of disk accretion rates. Recall that the disk accretion rate contributes both to the local viscous heating through F_{acc} and the stellar irradiation through F_* . Once we have calculated the disk temperature, we can then compute the Toomre parameter as a function of radius. We also calculate the disk cooling time to demonstrate that if the disk is unstable, fragments that form could cool quickly and likely remain bound^{57;68}. We find that the cooling time is well below the critical threshold of $\sim 10\Omega^{-1}$ throughout the entire disk. The discrepancy between our calculated β values and those quoted in the literature derives mostly from the difference in radius between this disk and standard models with $r_d < 100\text{AU}$. Reaching $Q = 1$ at larger radii requires lower values of Σ ; t_{cool} scales roughly as Σ^2 in the optically thick regime relevant for such disks. In addition, the dimensionless cooling time scales linearly with Ω , which also declines at large radii. In Figure 3 of the main text, we show contours of Q as a function of disk radius, and model accretion rate. The disk is marginally unstable between 150 – 320 AU.

While we do not have a direct constraint on \dot{M} , the observed total luminosity provides a benchmark: In a typical viscous accretion disk, the total luminosity is:

$$(9) \quad L_{\text{tot}} = L_* + \frac{GM\dot{M}}{R_*}$$

where the second term represents the total accretion luminosity, half of which is liberated within the accretion disk, and half at the stellar surface. Thus we can derive the expected accretion rate assuming that the bolometric luminosity of $3.2 L_\odot$ ¹⁶ is equivalent to the the internal luminosity. This implies an accretion rate of $\sim 10^{-7} M_\odot \text{ yr}^{-1}$, which is consistent with star formation models⁶⁹ and observations of young stars^{70;71}. This band is highlighted in Figure 3 of the main text. We account for the contributions from both protostars in the

irradiation models, assuming that the accretion luminosity is divided equally between the two. Note that the internal luminosity is not precisely equivalent to the bolometric luminosity; the bolometric luminosity is calculated from integrating the observed spectral energy distribution from the near-infrared to the millimeter. To relate the bolometric luminosity to the internal luminosity, a correction must be applied to account for the geometric projection of the disk and outflow cavities. For an inclination of 45.4° the correction factor is ~ 0.9 ^[72]. Furthermore, external irradiation can add an additional $0.1 - 0.5 L_\odot$ to the bolometric luminosity (depending on the local environment); this amount should be subtracted from the bolometric luminosity to relate it to the internal luminosity. However, we do not apply these small corrections given that the uncertainty in the bolometric luminosity is $\sim 10\%$ or greater depending on wavelength coverage⁷⁰, and these corrections would only lower the internal luminosity (and by inference the accretion rate), causing Q to decrease and imply a greater degree of instability.

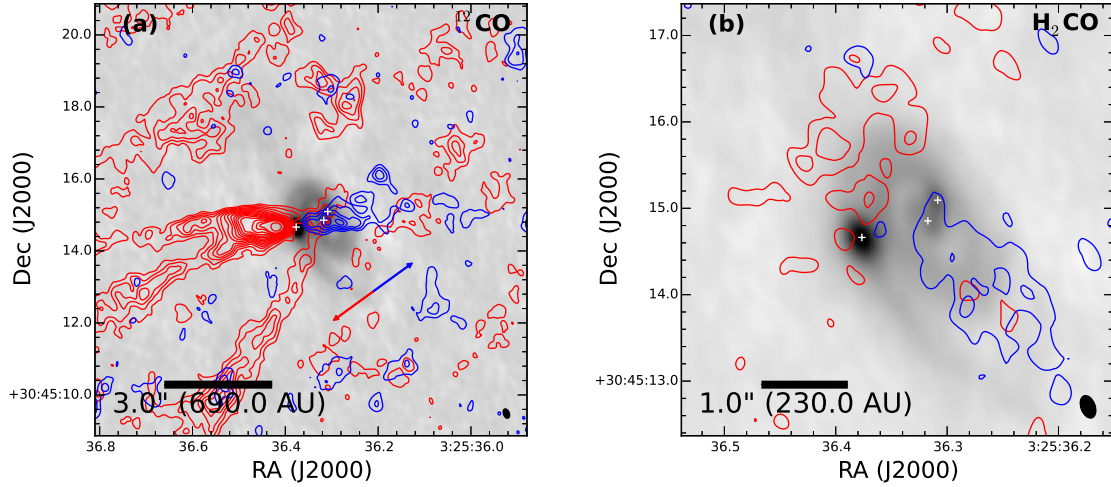
We next calculate a disk model using a self-consistently derived temperature profile based on the temperature dependent opacities. For simplicity we only consider the benchmark accretion rate $\dot{M} \approx 10^{-7} M_\odot \text{ yr}^{-1}$. We create a grid of allowed temperatures and identify the radial temperature profile that is consistent with the radiative equilibrium model of Equation 6. In Extended Data Figure 4, we show the derived surface density, temperature, Q and β profiles. Once again we find that the disk is marginally unstable at radii in the vicinity of the recently formed tertiary, IRS3B-c.

We note that although the surface density profile was derived using the optically thin approximation, our modeling indicates that the inner most regions may be marginally optically thick ($\tau \sim 2-5$). Thus, we are likely underestimating the total disk mass somewhat. This would only corroborate our conclusion that the disk has most likely experienced a recent episode of gravitational instability and fragmentation.

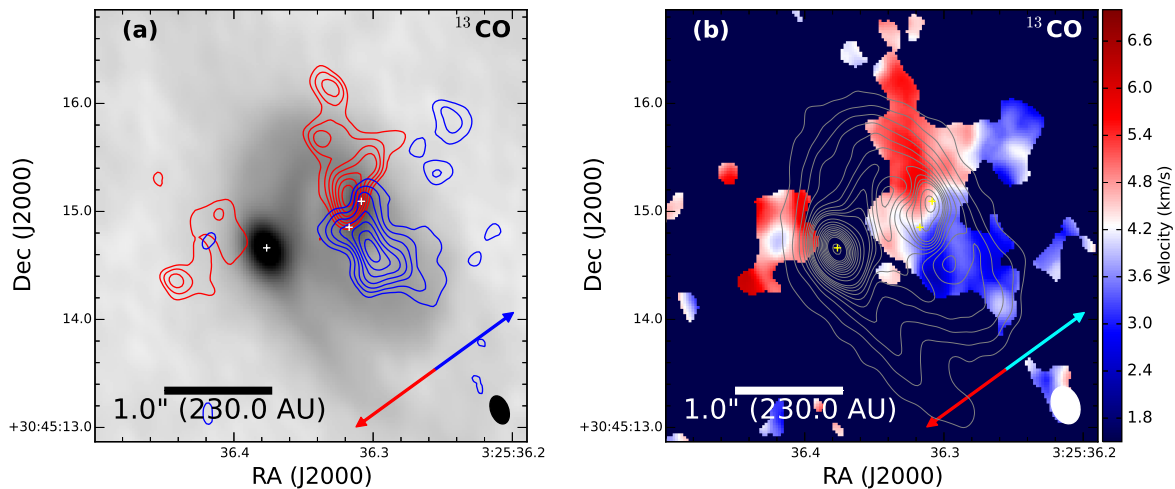
Finally, we can obtain a rough estimate of the mass scale of a fragment born in such a disk. The most unstable wavelength (i.e., physical scale) on which perturbations grow due to GI is $\lambda = 2\pi H$, where H is the disk scale height (see Equation 1 of the main text). The typical initial fragment mass is, to order of magnitude, $M_{\text{frag}} \approx \epsilon \Sigma \lambda^2$, where Σ is the surface density of the disk. Various studies have debated the appropriate order unity coefficient ϵ to properly account for contraction and spiral arm overdensities⁷³. For $\epsilon = 0.5$, the fragment mass ranges from $0.005 - 0.015 M_\odot$ between $150 - 320 \text{ AU}$ for our disk model with an accretion rate of $10^{-7} M_\odot \text{ yr}^{-1}$.

Data Availability Statement: The ALMA data utilized to generate Figures 1 & 2 in the main text and Extended Data Figures 1, 2, 3, & 5 are from ALMA project 2013.1.00031.S. These data will be available in the ALMA archive and from the corresponding author upon request. The data from which Figure 3 of the main text was generated and Extended Data Figure 4 are provided with the paper.

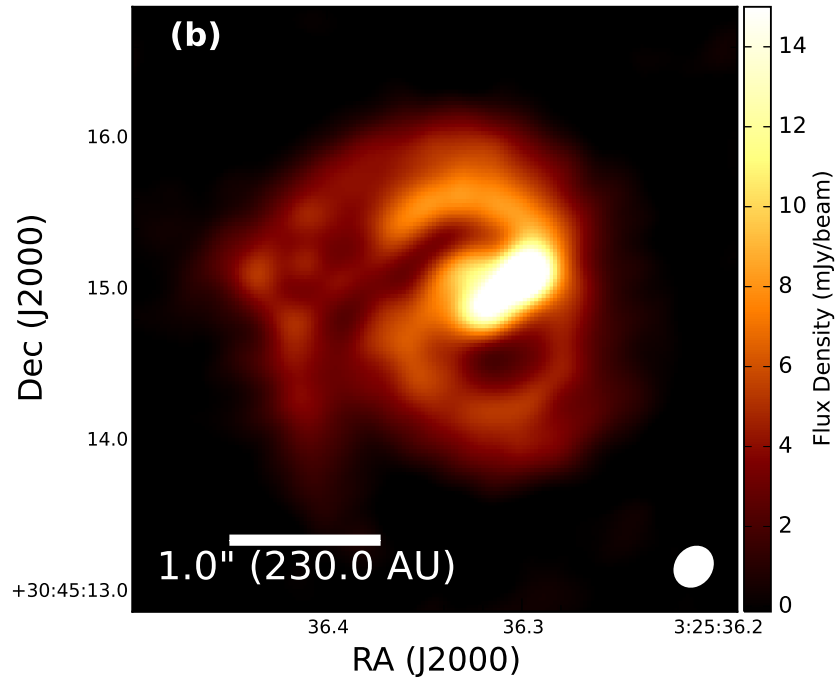
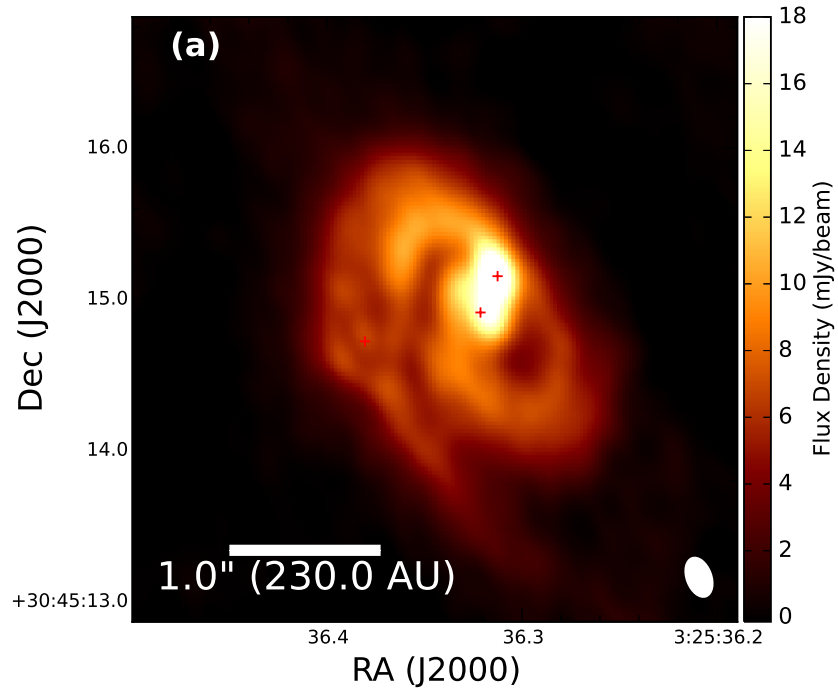
Code Availability: The data were reduced and processed using the Common Astronomy Software Applications (CASA) version 4.5.0. The code for generating the model found in Extended Data Figure 5 is publicly available⁵⁵ <http://dx.doi.org/10.5281/zenodo.13823>. The theoretical calculations presented in Figure 3 of the main text and in Extended Data Figure 4 are derived from solving algebraic equations and thus involve no significant code development.



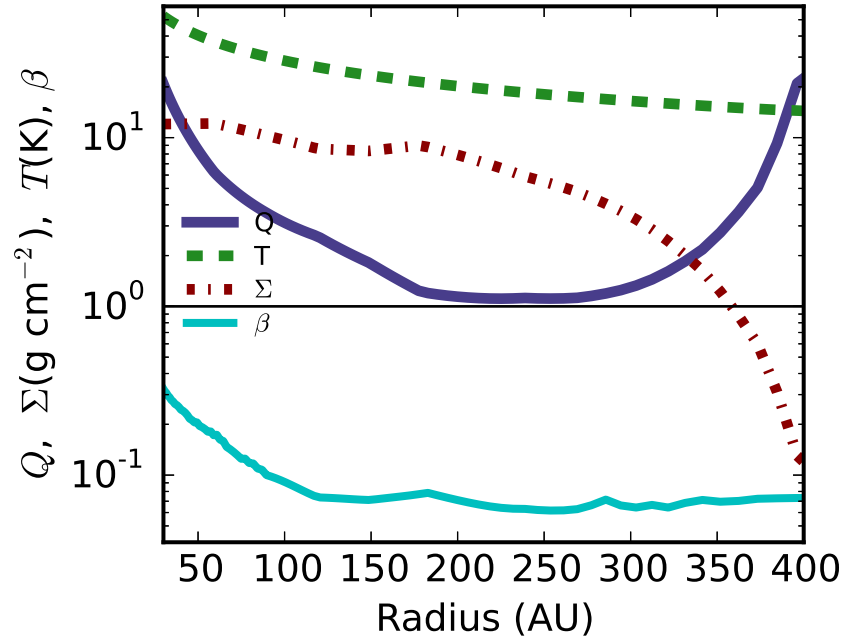
EXTENDED DATA FIGURE 1. Images of the ^{12}CO and H_2CO emission in the vicinity of L1448 IRS3B. Panel (a) shows the ^{12}CO red and blue-shifted contours overlaid on the 1.3 mm continuum and the same is shown in panel (b) but for H_2CO . The ^{12}CO emission in panel (a) most clearly shows a red-shifted outflow from the three protostars. There is a wide cavity that is traced back to IRS3B-a/b and a more collimated outflow is emitted from IRS3B-c, that is potentially generating the red-shifted arc that is within the wide outflow cavity. The blue-shifted side of the outflow is more diffuse and not well-recovered in our data, but appears associated with all three sources. The H_2CO emission in panel (b) has low intensity and traces a rotation gradient in the inner envelope and disk surrounding the sources. The source positions are marked with white crosses and the outflow direction is marked with the blue and red arrows in the ^{12}CO map. The contours in both plots start at 3σ and increase in 2σ intervals. The ^{12}CO emission was integrated over $-5.5 - 1.5 \text{ km s}^{-1}$ and $6.5 - 10.0 \text{ km s}^{-1}$ for the blue and red-shifted maps, respectively, and $\sigma_{\text{Blue}}=6.88 \text{ K km s}^{-1}$ and $\sigma_{\text{Red}}=5.02 \text{ K km s}^{-1}$. The H_2CO emission was integrated over $2.75 - 4.0 \text{ km s}^{-1}$ and $5.25 - 6.25 \text{ km s}^{-1}$ for the blue and red-shifted maps, respectively, and $\sigma_{\text{Blue}}=2.25 \text{ K km s}^{-1}$ and $\sigma_{\text{Red}}=2.05 \text{ K km s}^{-1}$.



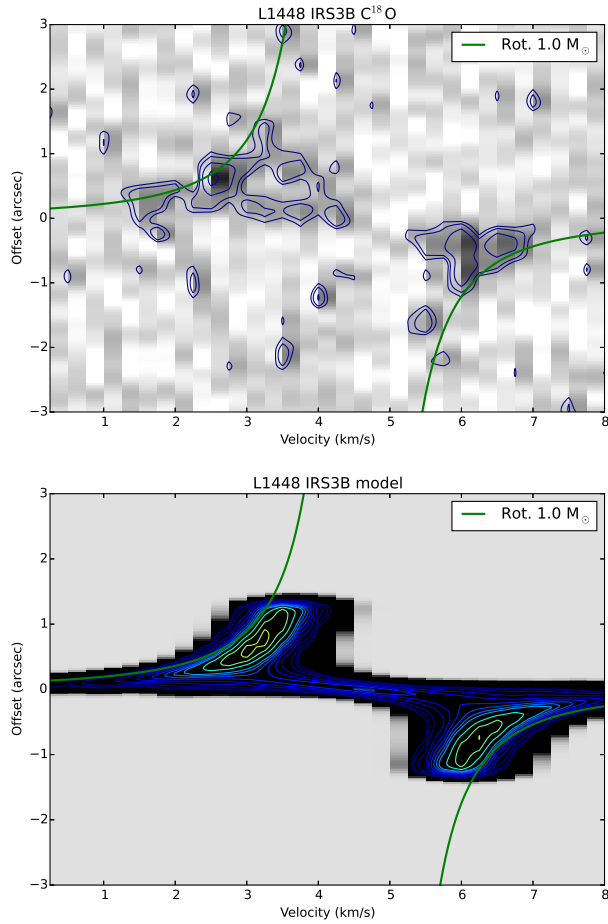
EXTENDED DATA FIGURE 2. Images of the ^{13}CO emission and its corresponding velocity maps from the disk around L1448 IRS3B, showing a rotation signature. (a) Red and blue-shifted ^{13}CO ($J = 2 \rightarrow 1$) emission overlaid on the ALMA 1.3 mm continuum image (grayscale) as red and blue contours. (b) Line-center velocity map of the ^{13}CO emission with 1.3 mm continuum contours overlaid in gray. The ^{13}CO traces higher-velocity emission near IRS3B-a and IRS3B-b and little of the extended disk due to spatial filtering. The source positions are marked with white or yellow crosses. The outflow direction¹⁴ is denoted by the blue and red arrows. The angular resolution of these data are given by the ellipse in the lower right corners, $0''.36 \times 0''.25$ ($83 \text{ AU} \times 58 \text{ AU}$). The contours in panel (a) start at 4σ and increase in 1σ intervals. The ^{13}CO emission was integrated over $1.25 - 4.0 \text{ km s}^{-1}$ and $5.5 - 7.0 \text{ km s}^{-1}$ for the blue and red-shifted maps, respectively. The noise levels for ^{13}CO are $\sigma_{Blue} = 4.99 \text{ K km s}^{-1}$ and $\sigma_{Red} = 3.2 \text{ K km s}^{-1}$. The continuum (gray) contours in panel (b) start at and increase by 10σ , at 100σ the levels increase in steps of 30σ , and at 400σ the levels increase by 100σ steps; $\sigma = 0.14 \text{ mJy beam}^{-1}$.



EXTENDED DATA FIGURE 3. ALMA 1.3 mm images with the brightest protostar (IRS3B-c) removed. The image with IRS3B-c removed, as observed is shown in panel (a). Panel (b) shows the image deprojected (rotated and corrected for system inclination), also with IRS3B-c removed.



EXTENDED DATA FIGURE 4. Plot of observed disk surface density (Σ), Temperature (T), Toomre Q , and dimensionless cooling β . Temperature is plotted as the green line, Toomre Q (blue line) is calculated self-consistently from the inferred surface density profile (red) using the temperature dependent opacities⁶⁶. The cyan line indicates the dimensionless cooling rate, β . The black line demarcates unity, where the disk is expected to be gravitationally unstable.



EXTENDED DATA FIGURE 5. Position-velocity diagrams of L1448 IRS3B and a model disk showing the rotation profile. A position-velocity (PV) cut is taken along the major axis of the disk (analogous to a long-slit spectrum), across the position of IRS3B-a and IRS3B-b and shown in the top panel. The solid line drawn in the panel is a Keplerian rotation curve for a $1.0 M_{\odot}$ central protostar, assumed to be the combined mass of IRS3B-a/b. A PV diagram for a thin disk model is shown in the bottom panel, demonstrating that a model with the same inclination angle shows a consistent PV diagram.

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