

# Water, methanol and dense gas tracers in the local ULIRG Arp 220: results from the new SEPIA Band 5 Science Verification campaign

M. Galametz,<sup>1</sup>★ Z.-Y. Zhang,<sup>1,2</sup> K. Immer,<sup>1</sup> E. Humphreys,<sup>1</sup> R. Aladro,<sup>3,4</sup>  
C. De Breuck,<sup>1</sup> A. Ginsburg,<sup>1</sup> S. C. Madden,<sup>5</sup> P. Møller<sup>1</sup> and V. Arumugam<sup>1</sup>

<sup>1</sup>European Southern Observatory, Karl-Schwarzschild-Str. 2, D-85748 Garching-bei-München, Germany

<sup>2</sup>Institute for Astronomy, University of Edinburgh, Royal Observatory, Blackford Hill, Edinburgh EH9 3HJ, UK

<sup>3</sup>European Southern Observatory, Avda. Alonso de Córdova 3107, Vitacura, Santiago, Chile

<sup>4</sup>Department of Earth and Space Sciences, Chalmers University of Technology, Onsala Observatory, SE-439 92 Onsala, Sweden

<sup>5</sup>Laboratoire AIM, CEA, Université Paris Diderot, IRFU/Service d'Astrophysique, Bat. 709, F-91191 Gif-sur-Yvette, France

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## ABSTRACT

We present a line survey of the ultraluminous infrared galaxy Arp 220, taken with the newly installed SEPIA (Swedish-European Southern Observatory PI receiver for APEX) Band 5 instrument on APEX (Atacama Pathfinder Experiment). We illustrate the capacity of SEPIA to detect the 183.3 GHz H<sub>2</sub>O 3<sub>1,3</sub>–2<sub>2,0</sub> line against the atmospheric H<sub>2</sub>O absorption feature. We confirm the previous detection of the HCN(2–1) line, and detect new transitions of standard dense gas tracers such as HNC(2–1), HCO<sup>+</sup>(2–1), CS(4–3), C<sup>34</sup>S(4–3) and HC<sub>3</sub>N(20–19). We also detect HCN(2–1)  $v_2 = 1$  and the 193.5 GHz methanol (4–3) group for the first time. The absence of time variations in the megamaser water line compared to previous observations seems to rule out an AGN nuclear origin for the line. It could, on the contrary, favour a thermal origin instead, but also possibly be a sign that the megamaser emission is associated with star-forming cores washed out in the beam. We finally discuss how the new transitions of HCN, HNC and HCO<sup>+</sup> refine our knowledge of the interstellar medium physical conditions in Arp 220.

**Key words:** ISM: molecules – galaxies: individual: Arp 220 – galaxies: ISM – galaxies: starburst.

## 1 INTRODUCTION

Arp 220 is one of the nearest ( $z=0.018$ ) ultraluminous infrared galaxies (ULIRGs;  $L_{\text{IR}} > 10^{12} L_{\odot}$ ). Many studies have analysed its massive molecular component and the physical conditions powering its intense infrared (IR) luminosity. Diffuse (e.g. low- $J$  CO) and dense (HCN, HNC, HCO<sup>+</sup>) gas tracers have been detected (Krips et al. 2008; Greve et al. 2009; Imanishi et al. 2010, hereafter K08, G09 and I10), providing us with estimates of the molecular gas masses and densities. Physical conditions of the interstellar medium can also be probed using maser (microwave amplification by stimulated emission of radiation) emission (Lo 2005). Even if the 22 GHz water megamaser emission has been detected in more than 150 extragalactic objects so far (it usually arises from active galactic nucleus (AGN) central engines in Seyfert 2 or LINER galaxies; e.g. Braatz, Wilson & Henkel 1996; Lo 2005), this line was not detected towards Arp 220. Instead, Arp 220 exhibits OH megamaser emission (Rovilos et al. 2003). Using the Institut de Radioastronomie Millimétrique (IRAM) 30 m telescope, Cernicharo, Pardo & Weiss (2006, hereafter C06) have detected the H<sub>2</sub>O(3<sub>1,3</sub>–2<sub>2,0</sub>) line at

183.3 GHz. The line, originating from a transition at a low energy level above ground state, is of sufficient luminosity to be classified as a megamaser. Given the absence of the 22 GHz water emission, they suggest that the 183 GHz emission is arising from the relatively spatially extended starburst (from about 10<sup>6</sup> star-forming cores) as opposed to denser molecular gas in the AGN nuclei.

The SEPIA (Swedish-European Southern Observatory PI receiver for APEX) Band 5 (159–211 GHz) instrument is a dual-polarization two-sideband receiver based on the Atacama Large Millimeter/submillimeter Array (ALMA) Band 5 receivers (Billade et al. 2012). The Science Verification (SV) campaign was carried out on the Atacama Pathfinder Experiment (APEX) antenna in 2015. The galaxy Arp 220 was observed as part of the SV campaign because of its importance as a local benchmark for high-redshift galaxies. The goals of the following analysis are to (i) compare the H<sub>2</sub>O(3<sub>1,3</sub>–2<sub>2,0</sub>) and HCN(2–1) lines observed with SEPIA Band 5 with previous observations in Arp 220 in order to test the calibration of the instrument and monitor the water line emission, (ii) target millimetre wavelength CH<sub>3</sub>OH maser emission and analyse new transitions of standard dense gas tracers such as HNC, CS, HC<sub>3</sub>N or CH<sub>3</sub>CN now accessible with SEPIA Band 5. We describe the observations and data reduction technique in Section 2 and present a detailed overview of the various

\*E-mail: [maud.galametz@eso.org](mailto:maud.galametz@eso.org)

lines detected in Section 3. We analyse the observations in Section 4.

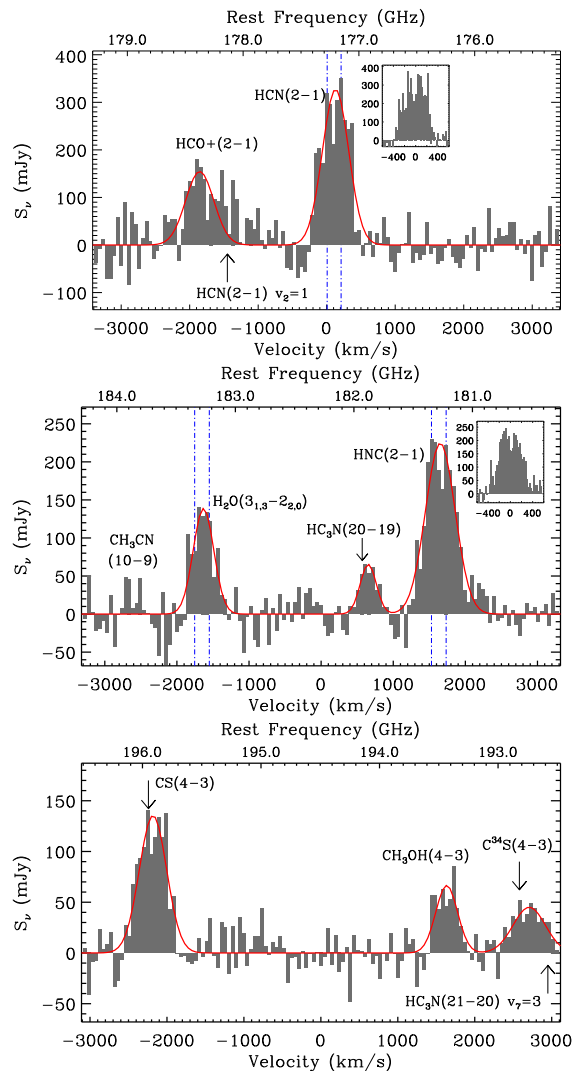
## 2 OBSERVATIONS AND DATA REDUCTION

SEPIA Band 5 covers the frequency range 159–211 GHz. The lower and upper sideband (LSB and USB) are separated by 12 GHz. Each sideband is recorded by 2 XFSTS (eXtended bandwidth Fast Fourier Transform Spectrometer) units of 2.5 GHz each, with a 1 GHz overlap. The spectral resolution of the data can reach up to  $0.065 \text{ km s}^{-1}$  and the beam size is  $\sim 35$  arcsec (equivalent to 12.7 kpc for Arp 220). We use a Jy/K factor of 34 to convert between antenna temperature  $T_A^*$  and flux density.<sup>1</sup> We refer to Billade et al. (2012) and Immer et al. (in preparation) for more details on the receiver. The performance of the instrument will be presented in a future paper. We conducted observations towards Arp 220 (Project ID: 095.F-9801; PI: Galametz) using the wobbler switch mode in 2015 July and September with a precipitable water vapour  $< 0.7$  mm. The observations were carried out for a total time of 6.2 h, with 2.8 h on source. Pointings were regularly checked on Carina and a calibration scan was taken every 10 min. Data were reduced using CLASS/GILDAS.<sup>2</sup> We removed a linear baseline from each individual spectrum before averaging the data. The rms reached for the water and HCN observations are 0.7 and 1.2 mK, respectively, at a  $50 \text{ km s}^{-1}$  resolution.

## 3 OVERVIEW OF THE DETECTED LINES

In this section, we present the lines detected by SEPIA Band 5 receiver towards Arp 220. We show the spectra in Fig. 1 and describe the lines in order of decreasing peak temperature. We provide the parameters of the Gaussian fits of the lines in Table 1. Note, however, that some of the lines have broad non-Gaussian profiles.

**HCN and HNC** – the hydrogen cyanide and isocyanide are very sensitive tracers of the dense molecular phase. The HCN luminosity is in particular used as a robust proxy for the star formation rate (e.g. Privon et al. 2015). Both HCN and HNC have been widely observed in Arp 220. HCN(2–1) is the strongest line in our observations. Its line width ( $452 \text{ km s}^{-1}$ ) is consistent with that of the HCN(3–2) line estimated by C06 but smaller than the line width of the HCN(2–1) already observed by K08. The two HCN(2–1) flux densities, however, match within the uncertainties (Table 1). The right-hand side of the HCN profile looks truncated. The width of the HNC(2–1) line is more consistent with the HCN line widths estimated in K08. Its peak temperature is  $\sim 1.5$  times lower than that of HCN(2–1) (see Section 4.3). Both HCN and HNC exhibit a double-peak profile. The HCN(2–1) line has a stronger peak at higher velocities ( $+170 \text{ km s}^{-1}$  shift compared to the systemic velocity). This asymmetry has already been found for HCN in the 3–2 and 4–3 transitions while the 1–0 line shows two symmetric peaks. On the contrary, the HNC(2–1) line shows a stronger peak at lower velocities ( $-50 \text{ km s}^{-1}$  shift compared to the systemic velocity) already observed in the 1–0 transition. Arp 220 possesses two nuclei, each of them being separated by about 1 arcsec and resolved in the radio continuum at 1 and 6 mm by Sakamoto et al. (1999) and Rodríguez-Rico et al. (2005). Our observations could indicate that a larger fraction of the HCN emission is arising from the redshifted eastern nucleus (as already suggested by I10) and that a larger fraction of the HNC emission is arising from the blueshifted western



**Figure 1.** SEPIA Band 5 observations towards Arp 220 (6.2 h). The plotted spectral resolution is  $50 \text{ km s}^{-1}$ . We observe strong water and methanol emission as well as many tracers of the dense gas (HCN, HNC,  $\text{HCO}^+$ , CS,  $\text{HC}_3\text{N}$ ,  $\text{C}^{34}\text{S}$ ). Insets show the HCN and HNC double-peak profiles at a  $25 \text{ km s}^{-1}$  resolution ( $x$ -axis: velocity relative to the centre of the line). We overlay the position of their respective peaks on the  $\text{H}_2\text{O}$ , HCN and HNC lines (blue).

nucleus. However, recent observations of HCN with ALMA have shown that a double-peak velocity profile is observed towards each individual nucleus (Scoville et al. 2015; Martín et al. 2016). The double-peak profiles are thus probably linked with self- and continuum absorption along the line of sight. Follow-up imaging with ALMA would help disentangle the two effects in the future, when Band 5 becomes available on ALMA.

**$\text{HCO}^+$**  – the formylm line traces gas  $\sim 100$ – $500$  times denser than the low- $J$  CO. Graciá-Carpio et al. (2006) suggested  $\text{HCO}^+$  as a more robust dense gas tracer in (U)LIRGs than HCN because of the enhanced HCN-to-CO ratio of these galaxies linked with IR pumping or chemical enhancement. Previous observations of  $\text{HCO}^+$  in Arp 220 in the 1–0, 3–2 and 4–3 transitions are reported in Graciá-Carpio et al. (2006), G09, I10 and Aladro et al. (2015). We detect  $\text{HCO}^+(2-1)$  ( $\nu_{\text{rest}} = 178.4 \text{ GHz}$ ) for the first time towards Arp 220 (Fig. 1, top). The  $\text{HCO}^+$  line is blended with the HCN(2–1) line in its  $\nu_2 = 1$  vibrational state ( $\nu_{\text{rest}} = 178.1 \text{ GHz}$ ;  $\text{HCN}_{\text{vib-to-HCN}(2-1)}$  line ratio  $< 0.5$ ). In order to properly fit the  $\text{HCO}^+$  profile, we

<sup>1</sup> <http://www.eso.org/sci/activities/apexsv/sepia/sepia-band-5.html>

<sup>2</sup> Gildas Continuum and Line Analysis Single-dish Software, <http://www.iram.fr/IRAMFR/GILDAS>

**Table 1.** Characteristics of the various fitted lines (Gaussian fitting).

Setup	Line	$\nu_{\text{rest}}^a$ (GHz)	Vel. offset <sup>b</sup> (km s <sup>-1</sup> )	FWHM (km s <sup>-1</sup> )	$T_{\text{peak}}^c$ (mK)	$\int S dv$ (Jy km s <sup>-1</sup> )	$L'$ (10 <sup>7</sup> K km s <sup>-1</sup> pc <sup>2</sup> )	Previous measurements <sup>d</sup> FWHM	$L'$
Tuning 1 <sup>e</sup>	HCN(2–1)	177.26	+99 ± 13	452 ± 25	9.6 ± 0.4	157 ± 11	80 ± 6	530 ± 20	88 ± 26
	HCO <sup>+</sup> (2–1)	178.37	0 <sup>g</sup>	484 ± 25 <sup>f</sup>	4.6 ± 0.3	79 ± 7	40 ± 3	–	–
Tuning 2 <sup>f</sup>	HNC(2–1)	181.32	+80 ± 11	484 ± 25	6.6 ± 0.3	116 ± 8	57 ± 4	–	–
	HC <sub>3</sub> N(20–19)	181.94	+84 ± 29	262 ± 54	1.9 ± 0.2	18 ± 4	9 ± 2	–	–
	H <sub>2</sub> O(3 <sub>1,3</sub> –2 <sub>2,0</sub> )	183.31	+67 ± 14	332 ± 31	4.1 ± 0.3	50 ± 6	24 ± 3	310	25
	C <sup>34</sup> S(4–3)	192.82	+128 ± 42	495 ± 92	1.3 ± 0.1	23 ± 5	10 ± 2	–	–
	CH <sub>3</sub> OH(4–3)	193.45	+25 ± 27	333 ± 42	2.0 ± 0.2	23 ± 5	10 ± 2	–	–
	CS(4–3)	195.95	+60 ± 12	418 ± 25	4.0 ± 0.2	60 ± 5	25 ± 2	–	–

Notes. <sup>a</sup> $z = 0.018$ .

<sup>b</sup>Velocity offset of the fitted peak with the systemic velocity.

<sup>c</sup>In  $T_{\text{A}}^*$  scale.

<sup>d</sup>Same units. References: K08 and C06.

<sup>e</sup>Frequency ranges observed = [172.2–176.2 GHz] (LSB) and [184.2–188.2 GHz] (USB).

<sup>f</sup>Frequency ranges observed = [177.1–181.1 GHz] (LSB) and [189.1–193.1 GHz] (USB).

<sup>g</sup>The HCO<sup>+</sup>(2–1) line is blended with the HCN(2–1)  $\nu_2 = 1$  line. The central velocity and width of the HCO<sup>+</sup> line are fixed for the fitting.

thus use a Gaussian positioned at the expected frequency of the line and assume a similar width as the well-constrained HNC line (i.e. 484 km s<sup>-1</sup>). We find a peak temperature of 4.6 mK (0.16 Jy), twice lower than that of HCN.

CS – up to the 4–3 transition, the carbon monosulfide molecule traces gas densities of 10<sup>5</sup>–10<sup>6</sup> cm<sup>-3</sup> while higher transitions trace densities > 10<sup>6</sup> cm<sup>-3</sup> (Bayet et al. 2009; Shirley 2015). In contrast with HCN or HNC, CS seems to be less affected by radiative excitation from X-ray-dominated regions (XDRs) or by IR pumping (Meijerink, Spaans & Israel 2007). Previous observations of CS(2–1), (3–2), (5–4) and (7–6) towards Arp 220 are reported in G09 and Aladro et al. (2015). For the first time, the CS(4–3) line ( $\nu_{\text{rest}} = 195.9$  GHz) is prominently detected in Arp 220 (Fig. 1, bottom). Its width is three times larger than the average line width of CS(4–3) lines observed in local ( $D < 10$  Mpc) galaxies (Bayet et al. 2009). The CS(4–3) profile resembles the asymmetric broad profiles already observed in other transitions, with a redshifted stronger component (G09). The asymmetry could, as for HCN, be the sign of self-absorption: Scoville et al. (2015) also detected a double-peaked profile in the CS line in both nuclei with ALMA. Zhang et al. (2014) have shown that a strong correlation (that extends over eight orders of magnitude) exists between the CS luminosity and the IR luminosity, making the CS luminosity an excellent proxy of the star formation rate. In Arp 220, the  $L'_{\text{CS}(7-6)}$ -to- $L'_{\text{CS}(4-3)}$  ratio is equal to 0.76. Using the calibration of Zhang et al. (2014), this leads to an  $L_{\text{IR}}$  estimate of  $1.6 \times 10^{12} L_{\odot}$ , consistent with the  $L_{\text{IR}}$  determined by Rangwala et al. (2011) using *Herschel*/Spectral and Photometric Imaging Receiver (SPIRE) data (i.e.  $1.8 \times 10^{12} L_{\odot}$ ). The relation linking  $L_{\text{CS}(4-3)}$  and  $L_{\text{IR}}$  can now be calibrated with SEPIA Band 5.

H<sub>2</sub>O – submillimetre water line emission is considered to be an important gas cooling mechanism in warm molecular clouds and the H<sub>2</sub>O line ( $\nu_{\text{rest}} = 183.3$  GHz) only requires moderate densities ( $n(\text{H}_2) = 10^4$ – $10^5$  cm<sup>-3</sup>) and low temperatures ( $T \sim 40$ K) to be excited (Cernicharo et al. 1994). We confirm the detection of H<sub>2</sub>O(3<sub>3,1</sub>–2<sub>2,0</sub>) emission in Arp 220 (Fig. 1, middle), with a peak flux of 0.14 Jy and a width at half-power of 332 ± 31 km s<sup>-1</sup>. These values are consistent with the peak flux (0.17 Jy) and width at half-power (310 km s<sup>-1</sup>) obtained with IRAM data by C06 (see section 4.1 for further analysis).

CH<sub>3</sub>OH – extragalactic methanol maser emission has been detected towards many objects (Henkel et al. 1987; Sjouwerman et al. 2010).

In particular, millimetre methanol transitions are known to be sensitive (but sometimes degenerate) tracers of temperature and density. Millimetre CH<sub>3</sub>OH emission has been detected in Arp 220 by Martín et al. (2011), particularly the 5–4 group (241.8 GHz). Salter et al. (2008) also report an absorption feature from the 6.7 GHz methanol line. Our SEPIA observations give us an access to the 193.5 GHz methanol (4–3) group (Fig. 1, bottom). The peak flux is 68 mJy. The water and methanol lines have similar widths (equal to 333 km s<sup>-1</sup> for CH<sub>3</sub>OH). We note that this is at the lower end of the line width range (280–590 km s<sup>-1</sup>) estimated from thermal molecules in Arp 220 by G09.

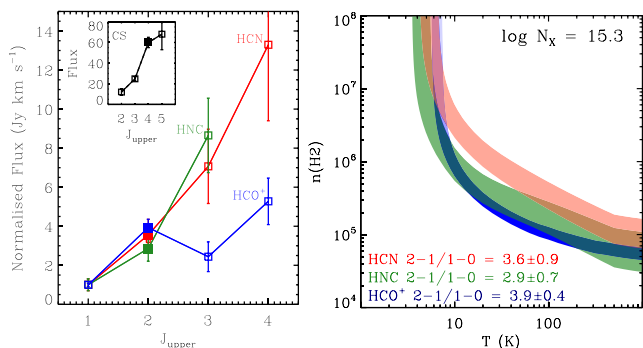
HC<sub>3</sub>N – molecules from the cyanopolyne family like HC<sub>3</sub>N are standard tracers of the dense gas in Galactic star-forming regions but are also detected in extragalactic environments (Aladro et al. 2011; Lindberg et al. 2011) in many transitions (from 9–8 to 28–27). Arp 220 is one of the very few ‘HC<sub>3</sub>N-luminous’ galaxies (i.e. that the HC<sub>3</sub>N(10–9) line intensity is at least 15 per cent of the HCN(1–0) line intensity; Aalto et al. 2002; Martín et al. 2011). The HC<sub>3</sub>N(20–19) line is narrow and detected at 16 per cent of the HNC(2–1) luminosity. Its profile looks symmetric but its peak is shifted. This is consistent with the asymmetry observed in the 10–9 transition by Aalto et al. (2002) and could suggest that most of the emission emerges from the eastern nucleus.

C<sup>34</sup>S – we observe a strong and broad C<sup>34</sup>S(4–3) line at 192.8 GHz. The line is unfortunately at the edge of the velocity range observed. An HC<sub>3</sub>N vibrationally excited line resides in a similar frequency range (192.6 GHz; see Fig. 1, bottom). Vibrationally excited emission of this species was already reported in Martín et al. (2011). The line should not contribute significantly due to its high energy level.

CH<sub>3</sub>CN – in Arp 220, the acetonitrile line was detected in both the ground and vibrationally excited states by Martín et al. (2011). Our observations cover the CH<sub>3</sub>CN(10–9) transition, for which we do not have a clear detection ( $2\sigma$ ; Fig. 1, middle). It is then difficult to say if the line is double-peak as the HCN line.

#### 4 ANALYSIS

This section presents a first analysis of the data focused on the water, methanol, HCN, HNC and HCO<sup>+</sup> lines. A more detailed analysis of the present and complementary data sets will be presented in a following paper.



**Figure 2.** Left: rotational ladders of HCN, HNC and HCO<sup>+</sup> (fluxes are normalized to that of the 1–0 transition). Inset: CS rotational ladder (not normalized). Filled squares are from this work. Right: RADEX modelling of the HCN, HNC and HCO<sup>+</sup> 2–1-to-1–0 ratios (units of Jy km s<sup>-1</sup>) in a kinetic temperature (in K) versus H<sub>2</sub> gas density (in cm<sup>-3</sup>) diagram. The bands indicate the parameter space covered by each ratio, including uncertainties. For the model presented here, the column density of each molecular species ( $N_X$ ) is fixed to the  $N(\text{HCN})$  estimated by Rangwala et al. (2011).

#### 4.1 Nature of the 183 GHz water emission

The 183 GHz H<sub>2</sub>O(3<sub>3,1</sub>–2<sub>2,0</sub>) line was observed by C06. Its intense luminosity suggested a megamaser (i.e. >10<sup>6</sup> times the luminosity of Galactic analogues) origin. We re-observe the H<sub>2</sub>O line to monitor its possible time variation if it is excited by AGN activity. We derive the isotropic line luminosity using

$$\frac{L}{L_{\odot}} = 10^{-3} \frac{1}{1+z} \frac{\nu_{\text{rest}}}{[\text{GHz}]} \frac{D^2}{[\text{Mpc}^2]} \frac{\int S dv}{[\text{Jy km s}^{-1}]} \quad (1)$$

with  $z$  being the redshift of the source,  $\nu_{\text{obs}}$  the rest frequency of the line,  $D$  the distance to the source and  $\int S dv$  the integrated flux density. We find an isotropic luminosity of  $4.6 \times 10^4 L_{\odot}$ . This is million times more luminous than the typical  $10^{-3} L_{\odot}$  Galactic 22 GHz maser emission (Genzel & Downes 1977). In Galactic targets, maser emission is typically stronger at 22 GHz compared to 183 GHz. Our detection is thus also more than a million times more luminous than the typical Galactic 183 GHz maser emission. It is also two to three orders of magnitude higher than 22 GHz masers detected in active galaxies (Braatz et al. 1996, among others). The Seyfert 2 galaxy NGC 3079 was the first extragalactic object where 183 GHz water emission was observed (Humphreys et al. 2005). The line was narrow (<50 km s<sup>-1</sup>), and interpreted as arising from a portion of the AGN central engine. In contrast, the 183 GHz water line in Arp 220 is much broader than Galactic 183 GHz analogues (Waters et al. 1980) or than in NGC 3079. It has a profile similar to the HCN or HNC lines [the full width at half-maximum (FWHM) is 1.4 times narrower for H<sub>2</sub>O] that could suggest a thermal origin. The single-dish observations do not allow us to disentangle between an extended/thermal and a nuclear/maser origin. AGN water megamasers at 22 GHz typically vary by factors of a few. Yet, we do not find variation between the 183 GHz water emission measured in 2005 by C06 and with SEPIA in 2015: the line has a similar width and luminosity. This lack of variability probably argues against the emission originating from an AGN nucleus in Arp 220 and provides weight to the proposal by C06 that the water megamaser emission in Arp 220 is associated with star formation and arises from individual cores (estimated at 10<sup>6</sup> by C06). The variabilities in these individual cores could then be washed out in a single-dish spectrum. OH megamasers are, in the same way, not commonly reported to be highly variable in single-

dish studies (variability ≤30 per cent; e.g. Darling & Giovanelli 2002). More SEPIA observations are needed to monitor possible variations in the water line. We finally note that González-Alfonso et al. (2012) suggested that the extreme H<sub>2</sub>O abundances in Arp 220 could also be the result of a hot core chemistry (i.e. characterized by the evaporation of icy grain mantles; see Charnley 1995, among others).

#### 4.2 Nature of the 193.5 GHz methanol emission

Extragalactic methanol maser emission was only classified as megamaser in NGC 1068 (84.5 GHz; Wang et al. 2014) and Arp 220 (36.2 and 37.7 GHz; Chen et al. 2015). We find an isotropic line luminosity of  $2.3 \pm 0.3 \times 10^4 L_{\odot}$  for the 193.5 GHz methanol line, here again suggesting a megamaser origin. The luminosity is nine orders of magnitude higher than typical Galactic or Large Magellanic Cloud methanol masers (Breen et al. 2010; Ellingsen et al. 2010). The line is much broader than typical extragalactic methanol masers or than the megamaser emission detected in Arp 220 by Chen et al. (2015), which could also favour a thermal origin. We note that the luminosity of the CH<sub>3</sub>OH(2<sub>k</sub>–1<sub>k</sub>) line ( $7.5 \times 10^7$  K km s<sup>-1</sup> pc<sup>2</sup>; Aladro et al. 2015) is lower than the luminosity of the 4–3 transition. In local thermodynamic equilibrium (LTE) conditions, the CH<sub>3</sub>OH(4–3) is expected to be more difficult to excite than the lower transitions if it was resulting from a pure thermal excitation. We cannot, however, rule out a thermal origin in the case of highly non-LTE excitation conditions. Getting more CH<sub>3</sub>OH lines at other frequencies would be necessary to understand the origin of the emission.

#### 4.3 Line ratios and excitation conditions

Fig. 2 shows the HCN, HNC, HCO<sup>+</sup> and CS spectral line energy distributions (SLEDs). Fluxes are gathered from this work, Aalto et al. (2009) and G09 (we use the rms weighted averages of the reported line luminosities). Arp 220 is one of the few objects where the HNC/HCN(1–0) and (3–2) line ratios are close or above unity (Huettemeister et al. 1995; G09). In contrast, we find an HNC/HCN(2–1) luminosity ratio well below unity (~0.7) that suggests a more standard HNC chemistry. A proper HNC(2–1) survey would help assess how the HNC/HCN(2–1) ratio in Arp 220 compares with other nearby galaxies. We use the non-LTE radiative transfer code RADEX (van der Tak et al. 2007) to generate diagnostic plots of the HCN, HNC and HCO<sup>+</sup> 2–1/1–0 ratios for various kinetic temperatures and H<sub>2</sub> number densities. We simply assume here that the cosmic microwave background 2.73 K emission dominates the background emission in the mm regime. Fig. 2 shows the  $T$ – $n(\text{H}_2)$  parameter spaces covered by the 2–1/1–0 ratios if we assume the column density derived by Rangwala et al. (2011) from HCN high- $J$  absorption lines, i.e.  $2 \times 10^{15}$  cm<sup>-2</sup>. We find that the domains covered by the three species barely overlap. Increasing the background blackbody temperature shifts all solutions to lower  $T$  and  $n(\text{H}_2)$  but does not make the domains converge. A solution domain is reached when we assume higher column densities (10<sup>16</sup> cm<sup>-2</sup>). It requires lower  $n(\text{H}_2)$  (~10<sup>5</sup> cm<sup>-3</sup>) than those derived in Rangwala et al. (2011), but similarly high temperatures (>300 K). Another solution domain is reached for  $n(\text{H}_2)$  ~10<sup>7</sup> cm<sup>-3</sup> but the temperatures (<10 K) requested are too low to be realistic (much lower than the temperature of the dust in Arp 220). G09 and I10 have shown that using HCO<sup>+</sup> line ratios results in  $n(\text{H}_2)$  estimates one to two orders of magnitude lower than those derived from HCN. In contrast, our moderate HCN/HCO<sup>+</sup>(2–1) line ratio does not lead to such a strong discrepancy. We note however that

the estimated  $\text{HCO}^+(2-1)$  flux is high (Fig. 2, left). Because the  $\text{HCN}(2-1) v_2 = 1$  line partly contributes to the flux, the  $\text{HCO}^+(2-1)$  flux estimate is uncertain and should probably be considered as an upper limit.

The high  $\text{HCN}(2-1/1-0)$  renders the ratios difficult to reconcile with our simple model. In fact, many additional physical processes are at work in Arp 220. First, its HCN emission is thought to be enhanced via IR pumping. Rangwala et al. (2011) suggested, for instance, that the HCN high- $J$  absorption lines might be populated via an IR pumping of the vibrational states. Vibrationally excited HCN is indeed observed towards Arp 220 (Aalto et al. 2015; Martín et al. 2016), and our detection of the  $\text{HCN}(2-1) v_2 = 1$  line is an additional sign that IR pumping processes are at work in Arp 220. These processes should also affect the HNC line, as suggested by Aalto et al. (2007). Secondly, our simple model neglects the effects of a hot dust background that could be non-negligible (even on the line ratios) in extreme environments such as Arp 220 (Tunnard & Greve 2016). Finally, the double-peak profiles of the HCN and HNC lines are strong signs of self- and continuum absorption whose effects on the various transitions of the lines still need to be properly quantified (Papadopoulos, Isaak & van der Werf 2010; Rangwala et al. 2015; Martín et al. 2016). These processes will be investigated and integrated in a more complex modelling in a following paper.

## 5 CONCLUSIONS

This Letter presents new observations towards Arp 220 using the SEPIA Band 5 instrument on APEX. We demonstrate the capacity of SEPIA to observe the 183 GHz water and 193.5 GHz methanol lines in this extreme extragalactic environment. Their luminosities classify them as megamaser but the absence of variation in the water line could rule out an AGN nucleus origin. SEPIA is an ideal instrument to monitor variations of maser lines. Our observations complement the low- $J$  SLEDs of standard molecular species, with lines such as  $\text{HCN}(2-1)$ ,  $\text{HNC}(2-1)$  and  $\text{HCO}^+(2-1)$  or  $\text{CS}(4-3)$  lines. Arp 220 exhibits a more standard  $\text{HNC}/\text{HCN}(2-1)$  ( $< 1$ ) and moderate  $\text{HCN}/\text{HCO}^+(2-1)$  line ratio compared to other transitions. The HCN and HNC lines exhibit double-peak profiles probably linked with self- and continuum absorption. New transitions of  $\text{HC}_3\text{N}$  and  $\text{CH}_3\text{CN}$  and a vibrational excited HCN line are also detected with the instrument.

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