

LOFAR, VLA, AND *CHANDRA* OBSERVATIONS OF THE TOOTHBRUSH GALAXY CLUSTER

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## ABSTRACT

We present deep LOFAR observations between 120 and 181 MHz of the “Toothbrush” (RX J0603.3+4214), a cluster that contains one of the brightest radio relic sources known. Our LOFAR observations exploit a new and novel calibration scheme to probe 10 times deeper than any previous study in this relatively unexplored part of the spectrum. The LOFAR observations, when combined with VLA, GMRT, and *Chandra* X-ray data, provide new information about the nature of cluster merger shocks and their role in re-accelerating relativistic particles. We derive a spectral index of  $\alpha = -0.8 \pm 0.1$  at the northern edge of the main radio relic, steepening toward the south to  $\alpha \approx -2$ . The spectral index of the radio halo is remarkably uniform ( $\alpha = -1.16$ , with an intrinsic scatter of  $\leq 0.04$ ). The observed radio relic spectral index gives a Mach number of  $\mathcal{M} = 2.8^{+0.5}_{-0.3}$ , assuming diffusive shock acceleration. However, the gas density jump at the northern edge of the large radio relic implies a much weaker shock ( $\mathcal{M} \approx 1.2$ , with an upper limit of  $\mathcal{M} \approx 1.5$ ). The discrepancy between the Mach numbers calculated from the radio and X-rays can be explained if either (i) the relic traces a complex shock surface along the line of sight, or (ii) if the radio relic emission is produced by a re-accelerated population of fossil particles from a radio galaxy. Our results highlight the need for additional theoretical work and numerical simulations of particle acceleration and re-acceleration at cluster merger shocks.

**Key words:** galaxies: clusters: individual (RX J0603.3+4214) – galaxies: clusters: intracluster medium – large-scale structure of universe – radiation mechanisms: non-thermal – X-rays: galaxies: clusters

## 1. INTRODUCTION

Merging clusters of galaxies are excellent laboratories to study cosmic ray (CR) acceleration in dilute cosmic plasmas, investigate the self-interaction properties of dark matter, and probe the physics of shocks and cold fronts in the intracluster medium (ICM). Diffuse radio emission has been observed in the majority of massive merging clusters, indicating the

presence of relativistic particles (i.e., cosmic rays) and magnetic fields. Given the large extent of these sources ( $\sim 1$  Mpc) and the limited lifetime of the synchrotron emitting electrons, these relativistic particles must be created or (re)-accelerated in situ (e.g., Jaffe 1977).

The diffuse cluster radio sources in merging clusters are referred to as radio halos and relics, depending on their location, morphology, and polarization properties. During the last decade, considerable progress has been made in our understanding of these sources (for reviews see Feretti et al. 2012; Brunetti & Jones 2014).

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Radio halos are centrally located in merging galaxy clusters and have typical sizes of about one Mpc. They have smoother brightness distributions than radio relics and the radio emission roughly follows the X-ray emission from the thermal ICM. Halos are unpolarized above the few percent level, although a few cases of polarized emission from halos have been reported. However, the nature of the polarized emission is still being debated (e.g., Govoni et al. 2005; Pizzo et al. 2011).

For radio halos, a correlation has been found between cluster X-ray luminosity ( $L_X$ ) and radio halo power (e.g., Cassano 2010; Cassano et al. 2013). A similar correlation with radio power is found (e.g., Basu 2012; Cassano et al. 2013) for the integrated Sunyaev–Zel’dovich (SZ) effect signal, given by  $Y_{SZ}$ .  $Y_{SZ}$  traces the integrated pressure along the line of sight. Both  $L_X$  and  $Y_{SZ}$  are used as observational proxies of cluster mass.

Two main physical models to explain the presence of radio halos have been proposed. In the secondary (or hadronic) model, the CR electrons are secondary products of collisions between thermal ions and relativistic protons in the ICM (e.g., Dennison 1980; Blasi & Colafrancesco 1999; Dolag & Enßlin 2000; Miniati et al. 2001; Pfrommer & Enßlin 2004; Pfrommer et al. 2008; Keshet & Loeb 2010; Enßlin et al. 2011; Zandanel et al. 2014). In the turbulent re-acceleration model, electrons are re-accelerated by merger induced magnetohydrodynamical turbulence (e.g., Schlickeiser et al. 1987; Brunetti et al. 2001; Petrosian 2001; Pinzke et al. 2015). Secondary models are challenged by the large energy content of cosmic ray protons needed to explain radio halos with very steep spectra (Brunetti 2004; Brunetti et al. 2008) and by the non-detection of  $\gamma$ -rays (e.g., Jeltema & Profumo 2011; Brunetti et al. 2012; Ackermann et al. 2014).

Radio relics are elongated and often arc-like sources found in the outskirts of galaxy clusters. There is a considerable amount of evidence that radio relics trace shock fronts where particles are being (re-)accelerated, as was proposed initially in Enßlin et al. (1998). The key observational facts that support the shock-radio relic connection are (i) the high polarization of some relics, with the apparent magnetic field lines being parallel to the major axis of the relics (e.g., van Weeren et al. 2010; Bonafede et al. 2012; Kale et al. 2012; de Gasperin et al. 2014), indicating that the ICM and associated magnetic fields are compressed, (ii) the presence of spectral index gradients, indicating electron cooling in the post-shock region of an outward traveling shock wave (e.g., Giacintucci et al. 2008; van Weeren et al. 2010, 2011; Stroe et al. 2013; Hindson et al. 2014), (iii) a change from power-law radio spectra at the outer edge of relics toward curved spectra in the post-shock region, indicating a site of acceleration and electron cooling (van Weeren et al. 2012a; Stroe et al. 2013), and (iv) the presence of ICM density and/or temperature jumps at the location of relics (e.g., Finoguenov et al. 2010; Macario et al. 2011; Akamatsu & Kawahara 2013; Bourdin et al. 2013; Ogorean & Brüggén 2013; Ogorean et al. 2014a). X-ray observations indicate that the shock Mach numbers ( $\mathcal{M}$ ) are low, typically  $\lesssim 3$ .

Although the above observations provide support for the basic scenario where particles are accelerated by outward traveling shock waves, there are still several open questions and observational puzzles. One of these puzzles concerns the low Mach numbers. The efficiency of acceleration via the diffusive shock acceleration (DSA) mechanism (e.g., Drury 1983) at these weak shocks in clusters is thought to be very low. A possible solution to this problem is that the shock re-

accelerates a population of already mildly relativistic electrons, instead of directly accelerating them from the thermal pool (e.g., Markevitch et al. 2005; Kang & Ryu 2011; Kang et al. 2012). There is some observational support for this shock re-acceleration model (van Weeren et al. 2013; Bonafede et al. 2014; Shimwell et al. 2015) and this would help to explain the existence of bright radio relics (Pinzke et al. 2013).

Particle-in-cell (PIC) simulations can provide crucial information about electron and proton acceleration at shocks. Recent PIC simulations indicate that instead of DSA, shock drift acceleration (SDA) is a more important mechanism at the initial stages of the acceleration process (Caprioli & Spitkovsky 2014; Guo et al. 2014a, 2014b). In these simulations, SDA results in efficient acceleration by low Mach number shocks in typical ICM conditions and thus potentially solves the low-acceleration efficiency problem of DSA.

Another observational puzzle is that some clusters host shocks that do not show corresponding radio relics (Russell et al. 2011; Shimwell et al. 2014). The opposite situation, i.e., relics without clear shocks, is also observed (e.g., Ogorean et al. 2014b). However, it is important to note that the X-ray detection of shocks in cluster outskirts, in particular those with low Mach numbers ( $\mathcal{M} \lesssim 2$ ), is difficult. Given that radio relics seem to trace low Mach number shocks, it is to be expected that in some cases these shocks escape detection in the X-rays, simply because of observational limitations. Some shocks coincide with the “edges” of radio halos (e.g., Markevitch et al. 2005; Shimwell et al. 2014), suggesting a relation between shocks and halos.

### 1.1. The “Toothbrush” Galaxy Cluster

The “Toothbrush” galaxy cluster (RX J0603.3+4214) hosts one of the largest and brightest radio relics, and therefore serves as an excellent target for a low-frequency radio study. The cluster was discovered by van Weeren et al. (2012a) and is located at  $z = 0.225$  (W. A. Dawson et al. 2015, in preparation). Besides the bright toothbrush-shaped radio relic, the cluster hosts at least two additional fainter relics and an elongated radio halo. The integrated Toothbrush radio relic spectrum has a power-law shape<sup>27</sup> between 74 MHz and 4.9 GHz with  $\alpha = -1.10 \pm 0.02$ . A clear north–south spectral index gradient is also observed for the relic. In addition, the radio spectra at the relic’s outer edge are straight and relatively flat, while the spectrum curves in regions with a steeper spectral index. From a radio color–color analysis (Katz-Stone et al. 1993) an injection spectral index of  $-0.65 \pm 0.05$  was found, corresponding to a Mach number of 3.3–4.6 (under the assumption of DSA in the linear test particle regime). The radio observations thus suggest that the Toothbrush traces a moderately strong, northwards-moving shock, that accelerates particles which later lose energy in the post-shock region.

Numerical simulations suggest that a triple merger event causes the peculiar shape of the putative shock that produces the Toothbrush relic (Brüggén et al. 2012). This merger scenario consists of a main north–south merger event, together with a third less massive substructure moving in from the southwest. Although *XMM-Newton* observations revealed that the main merger event takes place along a north–south axis, no evidence was found for a third smaller substructure (Ogorean et al. 2013). Several candidate shocks

<sup>27</sup>  $F_\nu \propto \nu^\alpha$ , where  $\alpha$  is the spectral index.

were identified in the cluster from X-ray surface brightness edges. In the north, a weak edge was found, but the location of the edge was offset from the expected shock location based on the relic position. W. A. Dawson et al. (2015, in preparation) find additional substructure in the cluster which indicates that the merger is likely more complex than a simple north–south event.

Deep H $\alpha$  imaging of the Toothbrush cluster and CIZA J2242.8+5301 was presented by Stroe et al. (2014). No boost in the number of H $\alpha$  emitters was found in the Toothbrush cluster, in stark contrast with the merging cluster CIZA J2242.8+5301.

Here we present LOw Frequency ARray (LOFAR) observations of the Toothbrush galaxy cluster. LOFAR is a powerful new radio telescope operating between 10 and 240 MHz (van Haarlem et al. 2013) which allows studies of diffuse cluster radio emission with high-angular resolution and sensitivity. To produce high-quality LOFAR images we developed a new calibration scheme which is presented in a separate paper. This calibration scheme will serve as a reference for future deep observations and surveys carried out with LOFAR. We combine the LOFAR data with Karl G. Jansky Very Large Array (VLA) 1–2 GHz L-band observations and previously published Giant Metrewave Radio Telescope (GMRT) data to study the radio spectral index. We also present new deep *Chandra* ACIS-I observations of the cluster. This allows us to search for surface brightness edges associated with shocks and cold fronts, and to map the ICM temperature distribution.

The layout of this paper is as follows. We start with an overview of the observations and data reduction in Section 2. We present the radio and X-ray images, together with the spectral index and temperatures maps in Section 3. The results of our search for shocks and cold fronts are described in Section 4. This is followed by a discussion and conclusions in Sections 5 and 6.

In this paper, we adopt a  $\Lambda$ CDM cosmology with  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_m = 0.3$ , and  $\Omega_\Lambda = 0.7$ . With the adopted cosmology, 1'' corresponds to a physical scale of 3.614 kpc at  $z = 0.225$ , the redshift of the Toothbrush cluster.

## 2. OBSERVATIONS

### 2.1. LOFAR HBA Observations

The Toothbrush cluster was observed on 2013 February 24, for 10 hr with the LOFAR High Band Antenna (HBA) stations. Complete frequency coverage between 112 and 181 MHz was obtained on the Toothbrush field. The nearby calibrator 3C147 was simultaneously observed using a second station beam.

For the observations 13 remote and 42 (split) core stations were used (see van Haarlem et al. 2013 for a description of the stations), giving baselines that range between 68 m and 80 km. The large number of short baselines is important to image diffuse extended emission.

Advanced calibration and processing techniques are needed to obtain deep high-fidelity images at low radio frequencies. We employ the facet calibration scheme to correct for direction dependent effects (DDEs) which are caused by the ionosphere and imperfect knowledge of the station beam shapes. Below we briefly summarize the

**Table 1**  
VLA Observations

Array configuration	D
Observation date	2013 Jan 28
Frequency coverage (GHz)	1–2
On source time (hr)	3.2
Correlations	full stokes
Channel width (MHz)	1
Integration time (s)	5
Resolution (arcsec)	$31 \times 24$
Image noise ( $\mu\text{Jy beam}^{-1}$ )	35

calibration of the LOFAR data. For more details about the calibration and observations the reader is referred to van Weeren et al (2016).

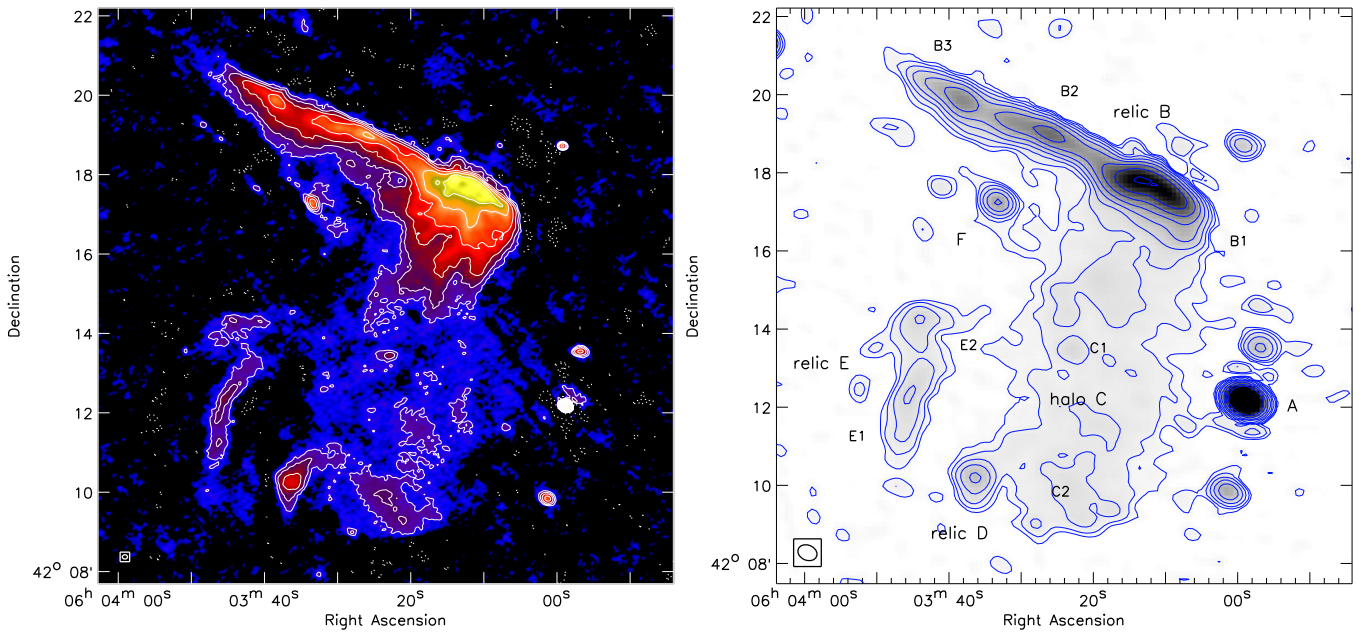
The first part of the data reduction consists of the removal of radio frequency interference (RFI), subtraction of bright off-axis sources appearing through the station beam sidelobes, and transfer of the flux-scale and clock corrections from the calibrator source 3C147 to the target field observation. After applying these corrections, we calibrate these data against a target field model derived from GMRT 150 MHz observations (van Weeren et al. 2012a). We decided to discard the data below 120 MHz since these data are generally more noisy and harder to calibrate. In the second part of the calibration, we obtain DDE corrections toward about 70 directions, so-called “facets,” and apply these corrections during the imaging. With this scheme we can overcome the limits of traditional self-calibration which cannot correct for DDEs. This facet calibration scheme allows us to obtain deep ( $\sim 0.1 \text{ mJy beam}^{-1}$ ) and high-resolution ( $\sim 5''$ ) images of the cluster in the HBA band.

### 2.2. VLA Observations

VLA L-band observations were taken in D-array on 2013 January 28. The data were recorded with 16 spectral windows each having 64 channels, covering the 1–2 GHz band. An overview of the observations is given in Table 1.

The data were reduced and calibrated with CASA (McMullin et al. 2007). As a first step the data were Hanning smoothed. We removed RFI using the AOFlogger and applied the elevation dependent gain tables and antenna offset positions. We calibrated the antenna delays, bandpass, and polarization leakage terms using the primary calibrator 3C147. The cross-hand delays and polarization angles were set using 3C138. Gain calibration was obtained for these calibrator sources as well. The gain solutions were then transferred to the target sources.

To refine the gain calibration several rounds of self-calibration were carried out. All imaging was done with W-projection (Cornwell et al. 2005, 2008), MS-MFS (Rau & Cornwell 2011), and multi-scale clean (Cornwell 2008). We used three Taylor terms ( $n_{\text{terms}}=3$ ) to take the frequency dependence of the sky into account. For the self-calibration we used Briggs (1995) weighting, with a `robust` factor of 0.0. Clean masks were also employed. The final images were corrected for the primary beam attenuation.



**Figure 1.** Left: LOFAR HBA 120–181 MHz image. The beam size is  $8''.0 \times 6''.5$  as indicated in the bottom left corner. Contour levels are drawn at  $[1, 2, 4, 8, \dots] \times 4\sigma_{\text{rms}}$ , with  $\sigma_{\text{rms}} = 93 \mu\text{Jy beam}^{-1}$ . Negative  $-3\sigma_{\text{rms}}$  contours are shown with dotted lines. Right: D-array VLA 1–2 GHz image with various sources labelled, following van Weeren et al. (2012a). Some of the sources are further subdivided into numbered components, e.g., relic E is divided into E1 and E2. The beam size is  $31'' \times 24''$ . Contour levels are drawn at the same levels as in the left panel, but with  $\sigma_{\text{rms}} = 35 \mu\text{Jy beam}^{-1}$ .

**Table 2**  
Properties of Diffuse Radio Sources in RX J0603.3+4214

	$S_{1.5\text{GHz}}^a$ (mJy)	$S_{150\text{MHz}}^a$ (Jy)	$P_{1.4\text{GHz}}$ ( $10^{24} \text{ W Hz}^{-1}$ )	LLS <sup>b</sup> (Mpc)	$\alpha_{150}^{1500}$
B	$317 \pm 31$	$3.9 \pm 0.4$	65	1.9	$-1.09 \pm 0.06$
C	$46 \pm 5$	$0.55 \pm 0.06$	9.1	$\sim 1.5$	$-1.08 \pm 0.06$
D	$6.5 \pm 0.7$	$0.086 \pm 0.009$	1.4	0.25	$-1.11 \pm 0.06$
E	$14.3 \pm 0.2$	$0.15 \pm 0.02$	2.5	1.1	$-1.03 \pm 0.06$

#### Notes.

<sup>a</sup> Integrated flux densities were measured from the LOFAR and VLA images that were used to make Figure 7. The regions where the fluxes were extracted are indicated in Figure 18.

<sup>b</sup> Largest linear size.

### 2.3. Chandra Observations

*Chandra* ACIS-I observations totaling 237 ks were taken in late 2013 (ObsID: 15171, 15172, 15323). The data were calibrated with the *chav* package<sup>28</sup>, following the processing described in Vikhlinin et al. (2005), applying the most recent calibration files.<sup>29</sup> The calibration and data reduction consists of filtering of counts with a recomputed ASCA grades 1, 5, or 7 and those from bad pixels, a correction for position-dependent charge transfer inefficiency, and the application of gain maps to calibrate photon energies. Periods of high background were also filtered by examining the count rate in the 6–12 keV band. The final exposure time available for the *Chandra* data after filtering was 236 ks. Standard blank sky background files were used for background subtraction. We used a pixel binning factor of 2. For more details about the data reduction the reader is referred to Vikhlinin et al. (2005).

<sup>28</sup> <http://hea-www.harvard.edu/~alexey/CHAV/>

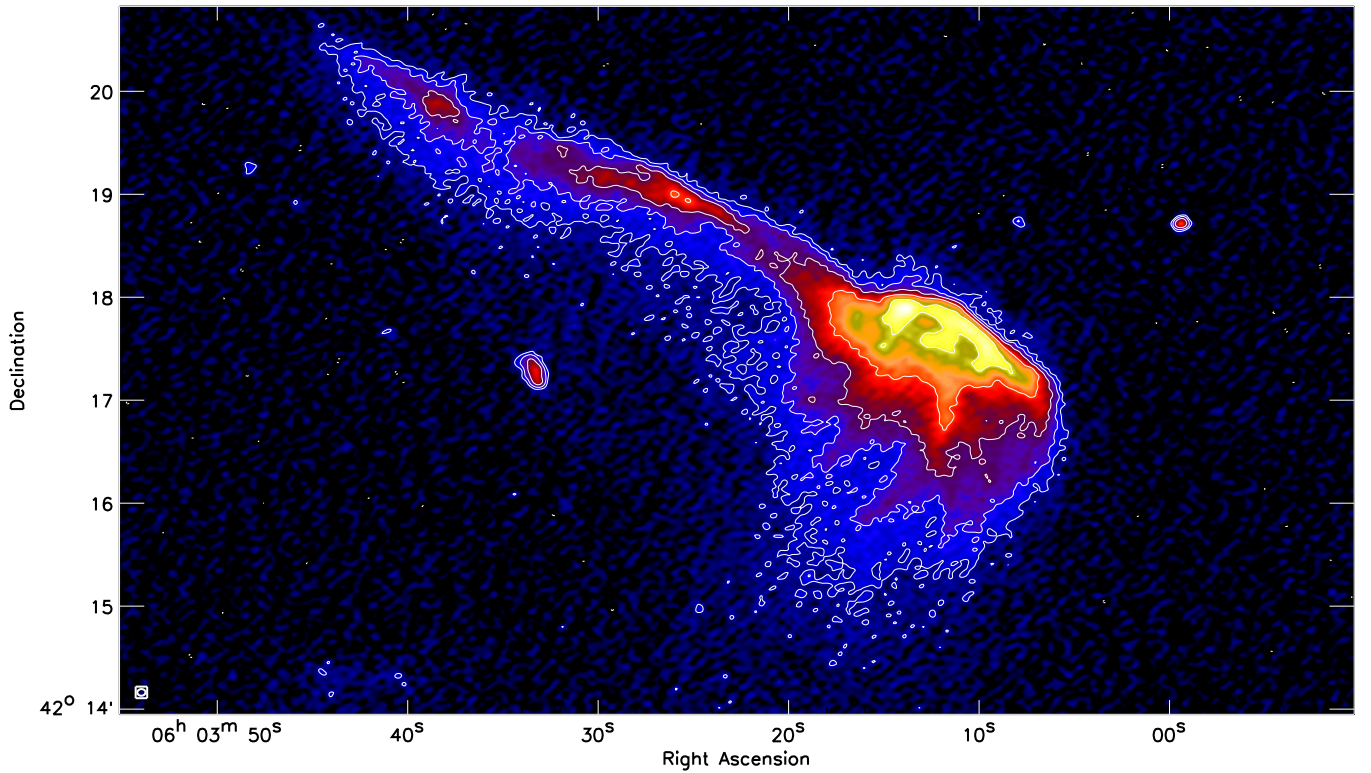
<sup>29</sup> We used CIAO v4.6 and CALDB v4.6.5.

## 3. RESULTS

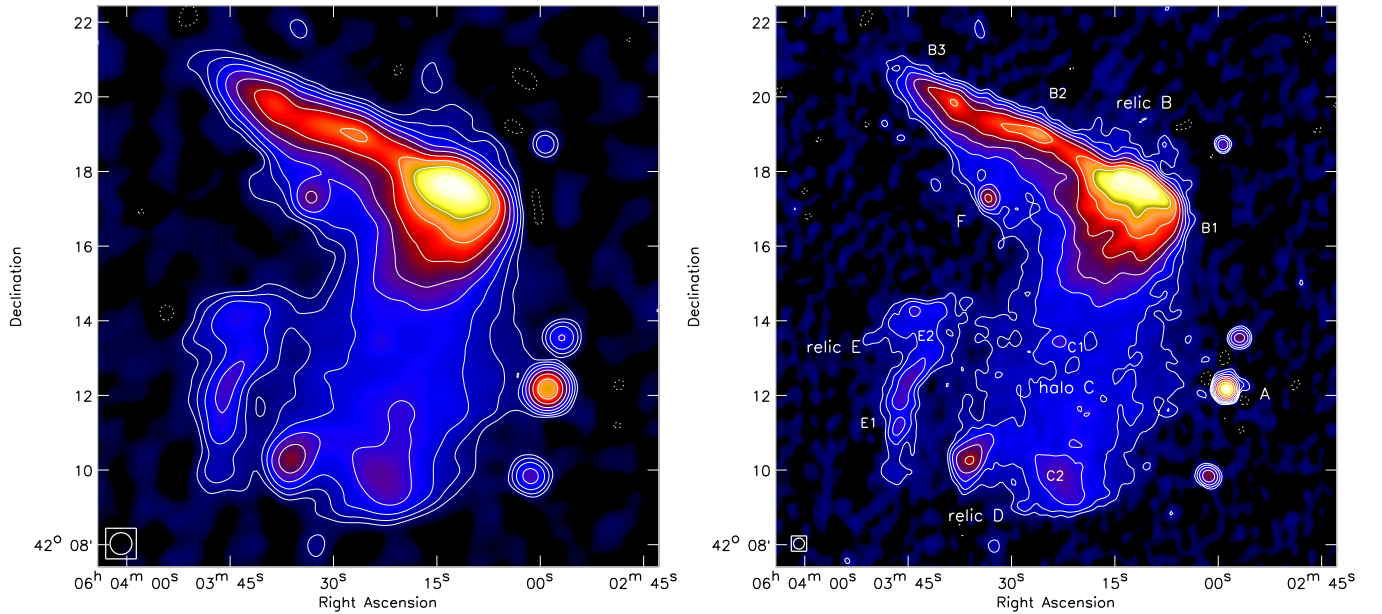
### 3.1. LOFAR 150 MHz Radio Continuum Images

Our deep  $8''.0 \times 6''.5$  resolution LOFAR image of the cluster, with *robust* =  $-0.25$  weighting, is shown in Figure 1 (left panel). The VLA 1–2 GHz D-array image, with a resolution of  $31'' \times 24''$ , is shown in the right panel of Figure 1. To facilitate the discussion, we labeled the diffuse and compact sources on the VLA image, following the scheme from van Weeren et al. (2012a). The LOFAR image (Figure 1) reveals both the previously known radio relics (sources B, D, E) and the extended radio halo (source C). The precise boundary, between where the relic emission from B1 ends and the emission from radio halo C starts, is difficult to determine. We define the boundary of the halo and relic to be located at the point where the spectral index approaches a value of about  $-1.1$ , the average spectral index of the radio halo, see Section 3.2 and Figure 18. An overview of the properties of the diffuse radio sources in the cluster is given in Table 2.

In terms of depth and resolution (assuming a spectral index of about  $-1$ ), the LOFAR image is comparable to our deep



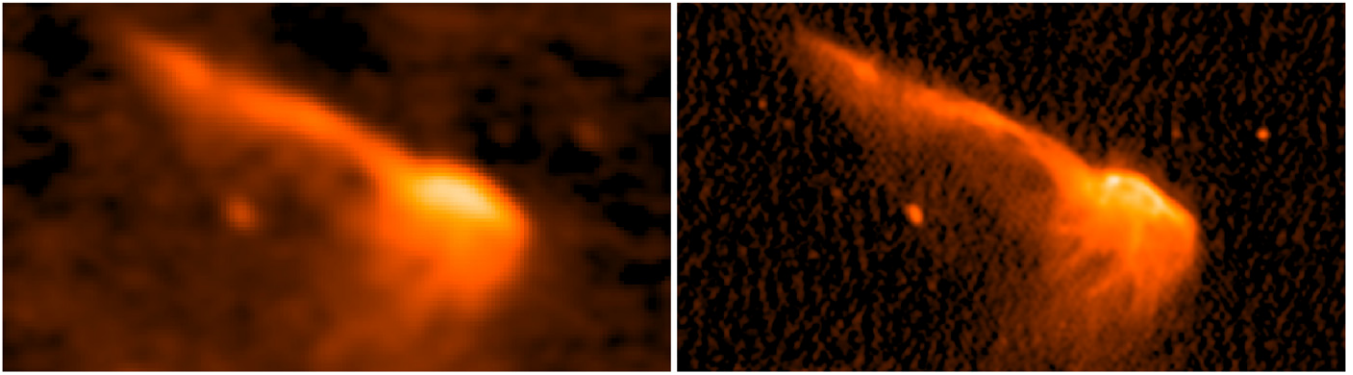
**Figure 2.** LOFAR HBA 120–181 MHz images made with `robust=-1.5` (i.e., close to uniform) weighting of the visibilities. The beam size is  $4''.8 \times 3''.9$  and indicated in the bottom left corner. Contour levels are drawn at  $[1, 2, 4, 8, \dots] \times 4\sigma_{\text{rms}}$ , with  $\sigma_{\text{rms}} = 121 \mu\text{Jy beam}^{-1}$ .



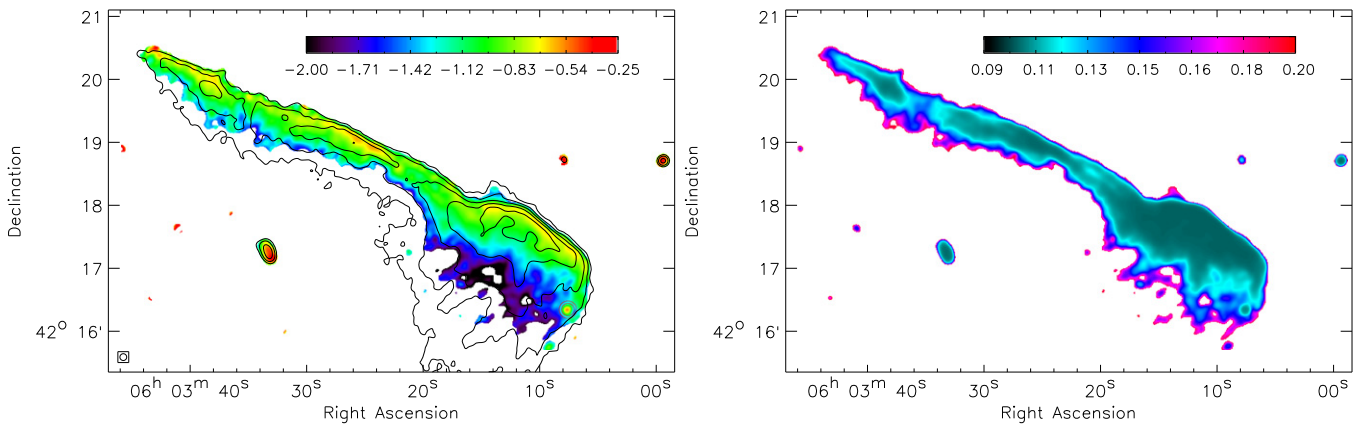
**Figure 3.** LOFAR HBA 120–181 MHz images with uniform weighting and uv-tapers of  $40''$  (left panel) and  $20''$  (right panel; sources are labelled as in Figure 1). The resulting beams are close to circular with sizes of  $35''.2$  and  $17''.7$  as indicated in the bottom left corner. Contour levels are drawn at  $[1, 2, 4, 8, \dots] \times 4\sigma_{\text{rms}}$ , with  $\sigma_{\text{rms}} = 358$  and  $212 \mu\text{Jy beam}^{-1}$ . Negative  $-3\sigma_{\text{rms}}$  contours are shown with dotted lines.

GMRT 610 MHz image (van Weeren et al. 2012a). A difference with respect to the GMRT 610 MHz observations is that the radio halo is clearly detected in the LOFAR image, while it is missing in the GMRT image. This difference is caused by the much denser inner uv-coverage of LOFAR, allowing excellent mapping of extended low-surface brightness

emission. Another difference with respect to the GMRT 610 MHz observations is that relic B is wider (i.e., extends further to the south, in particular component B1). This is expected given the steep north–south spectral index gradient. For a comparison with our GMRT 150 MHz image from van Weeren et al. (2012a), see Figure 4.



**Figure 4.** A comparison between the GMRT 150 MHz image (from van Weeren et al. 2012a), with a resolution of  $26 \times 22''$ , and the 120–181 MHz LOFAR image (from Figure 2), with a resolution of  $4.8 \times 3.9''$ . The colors in both images were scaled manually to achieve similar image surface brightness levels. Except for the fine-scale structures, the images show very good agreement.



**Figure 5.** Left: spectral index map between 151 and 610 MHz at  $6.5''$  resolution. Black contours are drawn at levels of  $[1, 2, 4, \dots] \times 1.25 \text{ mJy beam}^{-1}$  and are from the 151 MHz image. Pixels with values below  $4.5\sigma_{\text{rms}}$  were blanked. The circled source at the western end of the relic is discussed in Section 5.3. Right: corresponding spectral index error map. The errors are based on the individual rms noise levels in the two maps and using an absolute flux calibration uncertainty of 10% at 150 and 610 MHz.

We made an image with close to uniform weighting (Figure 2) to better enhance small-scale details in relic B. Several filamentary “streams” extend from the north of B1 to the south. These “streams” were also seen in the GMRT 610 MHz image and seem to “feed” the radio halo (C1) below B1. The northern boundary of relic B remains relatively sharp at  $4''$ – $5''$  resolution.

To better bring out the radio halo, we re-imaged the dataset with uniform weighting and tapered with Gaussians in the  $uv$ -plane corresponding to resolutions of  $20''$  and  $40''$ . These two images are shown in Figure 3. At these resolutions, relic B extends as far south as source F and faint emission connects the halo with relic E. The southern part of the halo, C2, has an enhanced surface brightness compared to the northern part C1.

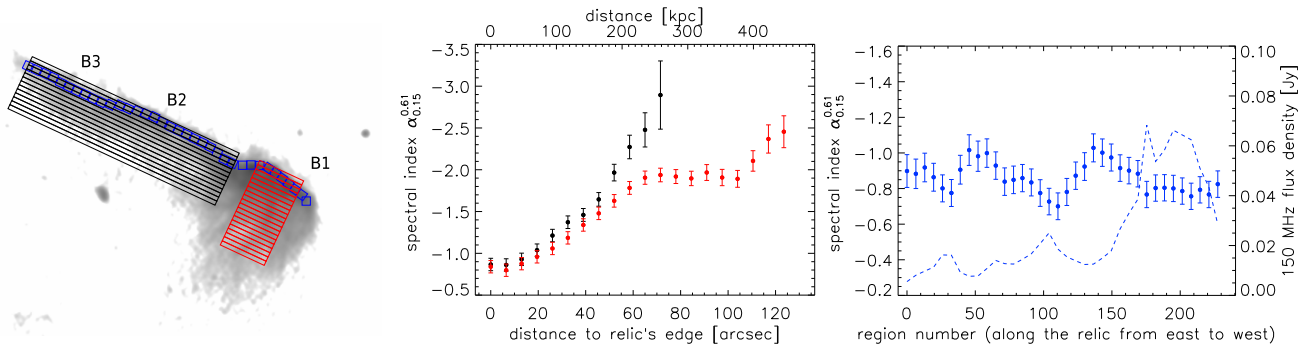
### 3.2. Spectral Indices

The high resolution LOFAR observations allow us, for the first time, to map the low-frequency 150–610 MHz spectral index at  $<10''$  resolution. To make the spectral index maps, we employed uniform weighting, used common inner  $uv$ -range cuts, and convolved the final images to the same resolution. The  $6.5''$  resolution 150–610 MHz spectral index map is shown in Figure 5. We blanked pixels with values less than  $4.5\sigma_{\text{rms}}$  in both the 150 and 610 MHz maps. We also computed spectral profiles across the relic, see Figure 6.

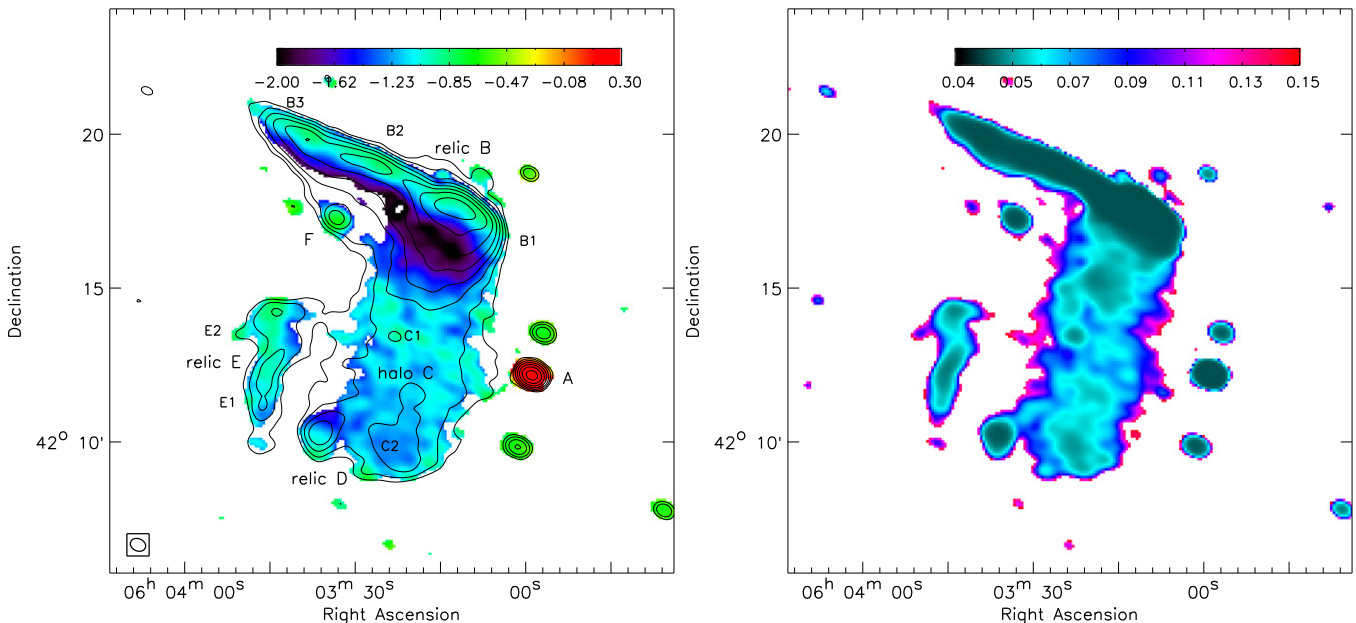
The spectral index map and profiles of relic B display a clear north–south spectral index gradient, with steepening from  $\alpha = -0.8$  to  $\alpha \lesssim -2.5$ . These numbers are consistent with our previous spectral index estimates ( $\alpha \approx -0.6$  to  $-0.7$ ) considering the difference of 10% in the flux-scale between the GMRT and LOFAR 150 MHz maps. With the new results presented here, a spectral index as flat as  $-0.6$  is however disfavored. In Figure 6 we also display the variation of the spectral index along the northern edge of the relic from east to west (a distance of 1.9 Mpc). The spectral index varies between  $-0.7 \pm 0.1$  and  $-1.0 \pm 0.1$ , with the brighter parts corresponding to flatter spectra.

The spectral index also steepens along the “streams” of emission that originate from B1. However, at their southern ends we find that the spectral indices are somewhat flatter than the surrounding emission, by about 0.3–0.5 units (see Figure 5). The streams might be the result of stretched magnetic field lines, allowing electrons to diffuse more efficiently along these field lines. Also stronger magnetic fields, combined with a diffusion time smaller than the electron cooling time, would “illuminate” lower energy electrons giving a flatter spectral index.

We also created a “low-resolution” ( $31'' \times 24''$ ) spectral index map by combining the LOFAR with the VLA 1–2 GHz D-array data (Figure 7). Pixels with values below  $3.5\sigma_{\text{rms}}$  were



**Figure 6.** Left: regions where the spectral indices were extracted (displayed in the middle and right panel) on top of a HBA 150 MHz image (used to make Figure 5). Middle: the spectral index between 150 and 610 MHz across relic B1, from north to south, is shown with red points. The black points cover the spectral index for the B2 and B3 part of the relic. The regions have a width of  $6''$ . Right: the blue points trace the spectral index along the northern edge of relic B, with the distance increasing from east to west. The blue regions have sizes of  $12''$ . The dashed blue line traces the LOFAR 150 MHz flux density along the relic, corresponding to the spectral indices shown in the same color.



**Figure 7.** Left: spectral index map between 0.151 and 1.5 GHz at  $31'' \times 24''$  resolution. Sources are labeled as in Figure 1). Black contours are drawn at levels of  $[1, 2, 4, \dots] \times 2 \text{ mJy beam}^{-1}$  and are from the 0.151 GHz image. Pixels with values below  $3.5\sigma_{\text{rms}}$  were blanked. Right: corresponding spectral index error map. The errors are based on the individual rms noise levels in the two maps and using an absolute flux calibration uncertainty of 10% at 150 MHz and 2.5% at 1.5 GHz.

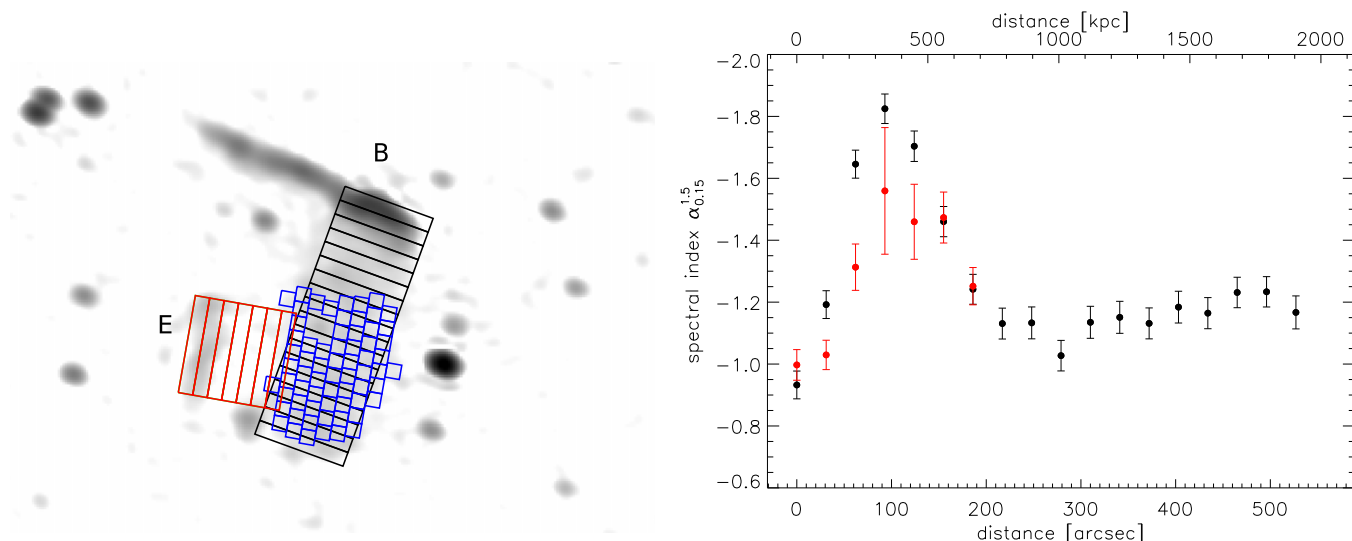
blanked. In this “low-resolution” spectral index map, we detect a spectral index gradient across source D, with  $\alpha$  steepening from  $-0.9$  to  $-1.4$  from south to north. We again extracted spectral profiles in different regions (Figure 8). We find a similar north–south trend across relic B1 as before, although in this case  $\alpha$  only steepens to about  $-1.8$  because of the limited spatial resolution and neighboring regions with flatter spectrum to the north and south. Similarly, the spectral index at the northern edge of relic B1 steepens to  $-0.9$  due to mixing of emission from the steeper spectrum regions just to the south. For the brightest part of relic E, the spectral index is  $-1.0 \pm 0.1$ . The spectrum then steepens toward the west to  $\alpha = -1.6 \pm 0.2$ , at the interface between relic E and radio halo C. This is a similar trend as seen for relic B1.

After steepening across relic B1 from north to south, the spectrum flattens and remains relatively uniform across radio halo C. Going from C1 toward C2, the spectral index only steepens slightly from  $-1.1$  to  $-1.2$  (Figure 8). Across almost the entire radio halo, the spectral index uncertainties are smaller

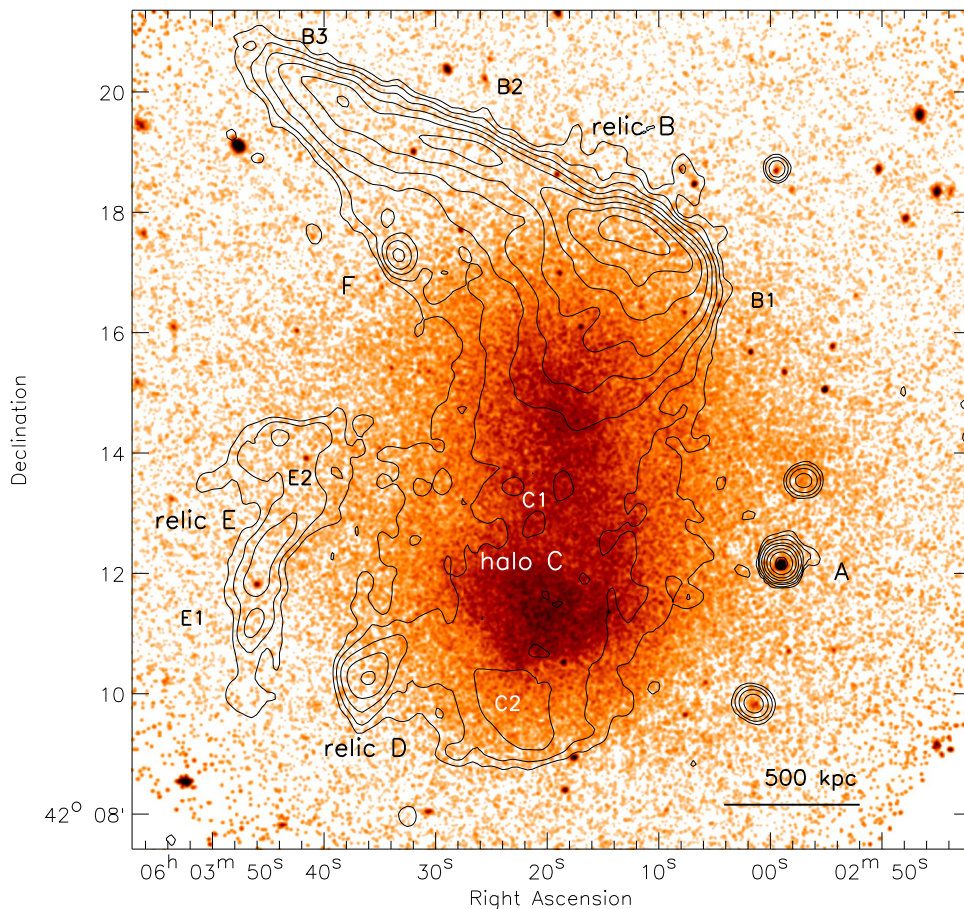
than 0.1. This makes it the best radio halo spectral map constructed so far (compare for example to the maps in, e.g., Feretti et al. 2004; Orrú et al. 2007; Vacca et al. 2014). In this case, the large spectral baseline, together with a deep low-frequency map, reduces the spectral index errors by a factor  $\gtrsim 2$ . Additionally, it increases the number of independent regions where the spectral index can be measured accurately (i.e., a spectral index uncertainty  $< 0.1$ ) by about an order of magnitude. In Section 5.5, we discuss the radio halo spectral index in more detail.

### 3.3. Chandra Image

The *Chandra* 0.5–2.0 keV image of the cluster is displayed in Figure 9. It shows extended north–south X-ray emission consisting of two main concentrations which were also clearly seen in previously published *XMM-Newton* images (Ogrea et al. 2013). The southern subcluster is better defined and brighter than the northern one. The southern subcluster displays a bullet-like shape with a surface brightness edge on



**Figure 8.** Left: regions where spectral indices were extracted are shown on top of the VLA 1–2 GHz image. The spectral indices for the red and black regions are shown in the left right panel. The spectra for the blue regions are shown in Figure 16. We excluded the regions affected by a compact source located at R.A.  $06^{\text{h}}03^{\text{m}}23^{\text{s}}$ , Decl.  $42^{\circ}13'26''$ . Right: the spectral index between 0.15 and 1.5 GHz across the E radio relic is shown in red (from east to west). The spectral index across relic B1 and the radio halo, from north to south, is shown in black. The regions have a width of  $31''$ .

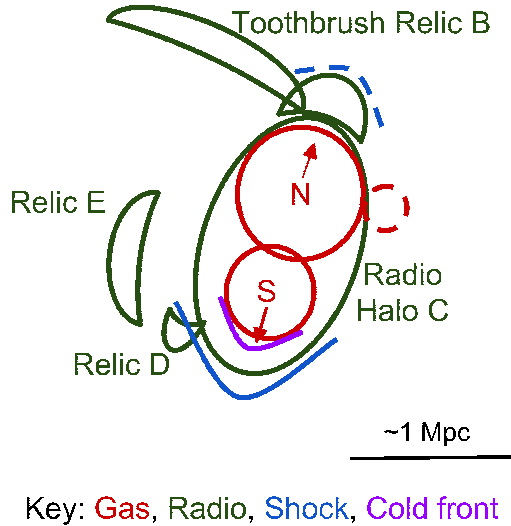


**Figure 9.** *Chandra* ACIS-I 0.5–2.0 keV image (bin = 2) of the Toothbrush cluster. The image is smoothed with a Gaussian with a FWHM of  $3''$ . Radio contours from the LOFAR image with a resolution of  $17''/7$  (right panel Figure 3) are overlaid. Radio sources are labelled as in Figure 1. The X-ray image shows two main concentrations, aligned north–south. The southern structure is better defined and has a bullet-like shape (see also Figure 12).

its southern side. Further to the south of this “bullet,” the *Chandra* image reveals a bow-shaped surface brightness edge. The edge coincides with the southern edge of the radio halo. We discuss the nature of these edges in Section 4. To

the west of the two main subclusters, about  $1'$  north of source A, we detect a smaller substructure (see also Figure 12). A schematic overview of the cluster is presented in Figure 10.





**Figure 10.** Schematic overview of the radio and X-ray components detected in the cluster. The dashed lines shows features whose identification is less certain.

### 3.4. X-Ray Temperature Map

To create a temperature map of the cluster, we divided the cluster into individual regions whose edges follow the surface brightness contours using `contbin` (Sanders 2006). This method is particularly useful for clusters where the surface brightness distribution has edges associated with density discontinuities. We required that each spectrum contains at least 5000 counts per region in the 0.5–7.0 keV band. Compact sources were detected in the 0.5–7.0 keV band with the CIAO task `wavdetect` using scales of 1, 2, 4, 8 pixels and cutting at the  $3\sigma$  level. We manually inspected the output and added a couple more sources that were identified by eye. The resulting spectra were then fitted with `XSPEC` (v12.8.2, Arnaud 1996). The regions covered by the compact sources were excluded from the spectra.

ICM spectra were fitted with an absorbed thermal emission model (`phabs` \* `APEC`). The metallicity was fixed to a value of  $0.3Z_{\odot}$  using the abundance table of Anders & Grevesse (1989). For the redshift we took the default value of  $z = 0.225$  (van Weeren et al. 2012a). The total Galactic H I column density was fixed to a value of  $N_{\text{H}} = 0.214 \times 10^{22} \text{ cm}^{-2}$  (the weighted average  $N_{\text{H}}$  from the Leiden/Argentine/Bonn (LAB) survey, Kalberla et al. 2005). The resulting temperature map is shown in Figure 11. The uncertainties on the temperatures are shown in Figure 17.

For most of the cluster regions, we measure temperatures of 7–10 keV. The main deviation from that is a region of hot gas located to the south of the southern subcluster’s core. The temperature in this region is 14–15 keV, with the  $1\sigma$  statistical uncertainties in the temperature of about  $-2$  keV and  $+3$  keV. The presence of this hot gas was also hinted at by the *XMM-Newton* observations (Ogrea et al. 2013). We extracted the spectra for the cores of the two main subclusters; these two regions are indicated in Figure 12. Using the same fitting parameters as for the temperature map, we find temperatures of  $8.43^{+0.26}_{-0.25}$  and  $9.00^{+0.28}_{-0.28}$  keV for the northern (*coreN*) and southern (*coreS*) subclusters, respectively.

## 4. CHARACTERIZATION OF SURFACE BRIGHTNESS EDGES

A visual inspection of the *Chandra* image reveals two edges in the surface brightness distribution in the southern subcluster. One of the surface brightness edges is located at the southern boundary of the radio halo, while the other one seems to be a contact discontinuity between the southern core and the hotter ICM in the cluster outskirts, see Figure 9. The locations of these discontinuities are indicated in Figures 10 and 12.

We used `PROFFIT` (Eckert et al. 2011) to extract and fit the surface brightness profiles. The areas covered by the compact sources were excluded from the fitting (see Section 3.4). We modeled the surface brightness profiles across the edges assuming underlying broken power-law density models, plus a constant ( $S_{\text{back}}$ ) to optionally include a background level. The emissivity is assumed to be proportional to the density squared and this model is then integrated along the line of sight assuming spherical symmetry. The sectors for which the fitting was performed are indicated in Figure 12 (see also Table 3). For the inner “edge” of the southern subcluster’s core we considered two different sectors (*coreS\_east*, *coreS\_west*) as the edge has a peculiar right-angled shape with a small apex. For the binning of photons for the partial annuli, we required a signal-to-noise ratio (S/N) of 5 per bin.

The broken power-law model, which can be used to represent a density jump associated with a shock or cold front, is given by

$$n(r) = \begin{cases} (n_2/n_1)n_0 \left(\frac{r}{r_{\text{edge}}}\right)^{a_2}, & r \leq r_{\text{edge}} \\ n_0 \left(\frac{r}{r_{\text{edge}}}\right)^{a_1}, & r > r_{\text{edge}}. \end{cases} \quad (1)$$

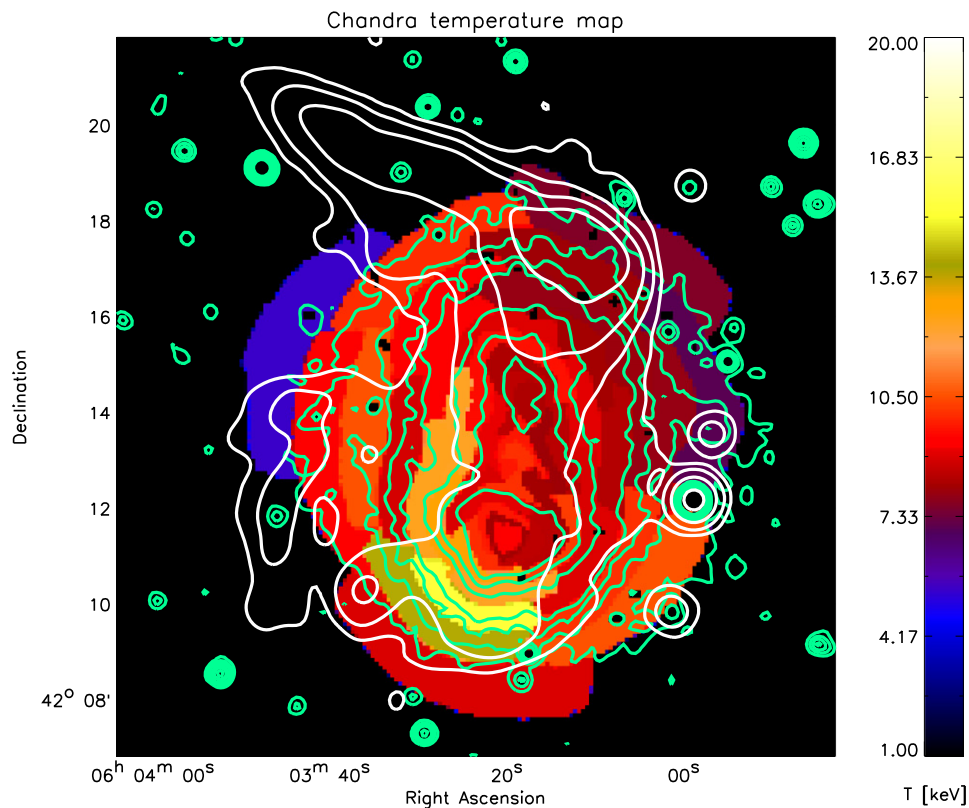
The subscripts 1 and 2 denote the upstream and downstream regions, respectively. The parameter  $n_0$  is a normalization factor,  $a_1$  and  $a_2$  define the slopes of the power-laws, and  $r_{\text{edge}}$  the position of the jump. The parameter  $n_2/n_1 \equiv C$  represents the electron density jump. In the case of a shock,  $n_2/n_1$  is related to the Mach ( $\mathcal{M}$ ) number via the Rankine–Hugoniot relation

$$\mathcal{M} = \left[ \frac{2C}{\gamma + 1 - C(\gamma - 1)} \right]^{1/2}, \quad (2)$$

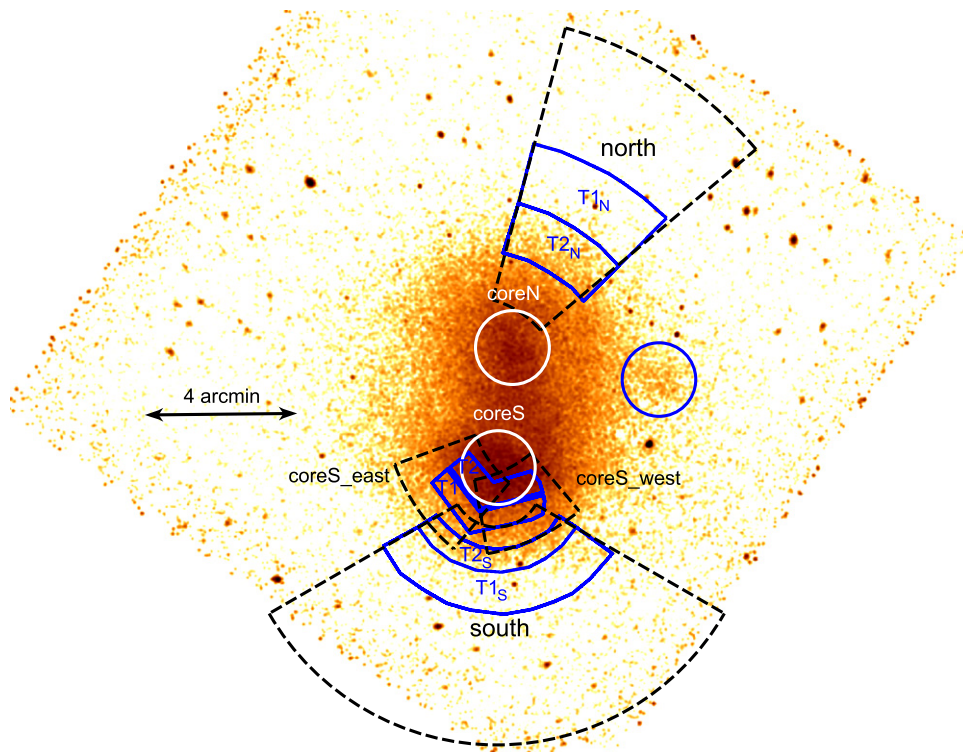
where  $\gamma$  is the adiabatic index of the gas. In this work, we assume  $\gamma = 5/3$ .

The modeling was done assuming a fixed background level of zero ( $S_{\text{back}} = 0$ ) because the *Chandra* instrumental and sky background components are already subtracted.

To determine the nature of the density jumps (i.e., shocks or cold fronts), we also need to measure temperatures on both sides of  $r_{\text{edge}}$ . For that we defined regions on each side of  $r_{\text{edge}}$  and extracted the spectra from the *Chandra* data. The locations of these regions are indicated in Figure 12. We fit the temperatures in the same way as described in Section 3.4. What we fit is thus the projected temperature, which is an approximation of the real 3D temperature. We note that the regions where the temperatures are extracted are quite large. We did not attempt to correct for possible temperature gradients or variations within these bins.



**Figure 11.** X-ray temperature map of the cluster. Each region was required to contain at least 5000 net counts in the 0.5–7.0 keV band. X-ray contours (green) and radio contours (white, Figure 3 left panel) are overlaid. Black colored regions are blanked, either due to compact sources or because of an insufficient number of ICM counts.



**Figure 12.** Smoothed (4 pixels) *Chandra* 0.5–2.0 keV image showing the sectors (dashed) in which we extracted and fitted the surface brightness profile, see Table 3. The regions where the temperatures ( $T_1$  and  $T_2$ ) are extracted are shown in blue. The surface brightness edges are located at the boundary between the  $T_1$  and  $T_2$  regions. The blue circle represents the area where we extracted the X-ray luminosity of the western substructure. The white circles indicate the areas for which we extracted the temperatures for the cores of the northern (*coreN*) and southern (*coreS*) subclusters.

**Table 3**  
Surface Brightness Profiles

	South	coreS_east	coreS_west	North
Sector center (R. A., J2000)	90°835417	90°785958	90°862000	90°853138
Sector center (Decl., J2000)	42°188797	42°223506	42°239089	42°225425
Opening angle	210°–330°	200°–230°	280°–310°	40°–75°
Fitting range	1'2–7'0	3'0–5'3	3'2–5'2	2'6–10'0
$n_2/n_1$	1.57 <sup>+0.09</sup> <sub>-0.08</sub>	1.35 <sup>+0.05</sup> <sub>-0.05</sub>	1.46 <sup>+0.05</sup> <sub>-0.05</sub>	1.26 <sup>+0.14</sup> <sub>-0.10</sub>
$r_{\text{edge}}$	2'45 <sup>+0.02</sup> <sub>-0.02</sub>	4'09 <sup>+0.03</sup> <sub>-0.02</sub>	4'12 <sup>+0.03</sup> <sub>-0.02</sub>	5'31 <sup>+0.17</sup> <sub>-0.19</sub>
$\chi^2/\text{d.o.f.}$	68.4/76	62.0/48	56.5/41	68.7/76

For a cold front, the pressure is the same on both sides of  $r_{\text{edge}}$ . For a shock and assuming  $\gamma = 5/3$ , the Rankine–Hugoniot jump conditions relate the upstream ( $T_1$ ) and downstream temperature ( $T_2$ ) to the Mach number (e.g., Finoguenov et al. 2010)

$$\mathcal{M} = \frac{\left(8 \frac{T_2}{T_1} - 7\right) + \left[\left(8 \frac{T_2}{T_1} - 7\right)^2 + 15\right]^{1/2}}{5}. \quad (3)$$

#### 4.1. Southern Sector

For the southern discontinuity (at the edge of the radio halo), we find that the broken power-law model provides a good fit to the data, with  $C = 1.57^{+0.09}_{-0.08}$ , see also Table 3.<sup>30</sup> The profile and resulting best-fitting model are shown in the right panel of Figure 13.

For the temperatures on both sides of the surface brightness edge, we measure  $T_{1S} = 7.8^{+1.6}_{-1.4}$  keV and  $T_{2S} = 14.5^{+2.9}_{-1.8}$  keV. Using Equation (3), we find  $\mathcal{M} = 1.8^{+0.5}_{-0.3}$ . Taking the density jump and Equation (2) we obtain  $\mathcal{M} = 1.39^{+0.06}_{-0.06}$ . Thus the results from the temperature and density jump indicate that the southern brightness edge traces a shock.

We also divided the southern sector into three equal parts, with opening angles of 40° to search for variations of the Mach number along the shock front. From east to west we obtain  $\mathcal{M} = 1.16^{+0.11}_{-0.07}$  (210°–250°),  $\mathcal{M} = 1.57^{+0.16}_{-0.13}$  (250°–290°), and  $\mathcal{M} = 1.43^{+0.09}_{-0.08}$  (290°–330°), based on the density jump determined from the surface brightness profile. This result indicates that the eastern part of the southern shock is weaker.

#### 4.2. Southern Core

For the southern subcluster’s eastern sector (*coreS\_east*), we obtained  $C = 1.37^{+0.05}_{-0.05}$ . For the western sector (*coreS\_west*), we obtained  $C = 1.46^{+0.05}_{-0.05}$ , which indicates very similar density jumps (see Figure 14 for the profiles). The reduced  $\chi^2$  of 1.3–1.4 (Table 3) means the fits are reasonable but not perfect, which is not surprising given the assumption of spherical symmetry and the power-law approximation to the inner density slope.

For the temperatures on both sides of the edge, we find  $T_1 = 11.4^{+1.2}_{-0.9}$  keV and  $T_2 = 9.0^{+0.4}_{-0.4}$  keV. The subscripts 1 and 2

<sup>30</sup> We repeated the fit, but not subtracting the sky background and leaving  $S_{\text{back}}$  as a free parameter. The value of  $C = 1.65^{+0.12}_{-0.11}$  we obtained in this way is consistent with our earlier result.

refer to the regions indicated in Figure 12. Here, we combined the eastern and western regions for the temperature measurements, motivated by the very similar density jumps measured for the eastern and western sectors (*coreS\_east* and *coreS\_west*). Taking the temperature and density jumps, we find that the pressures on both sides of  $r_{\text{edge}}$  are consistent within their errors as indicated by the ratio  $T_1/T_2$  and  $n_2/n_1$ . We therefore classify this surface brightness edge as a cold front. This combination of a merger related cold front and a shock ahead of it (Section 4.1), is very similar to what is seen in the Bullet Cluster (Markevitch et al. 2002).

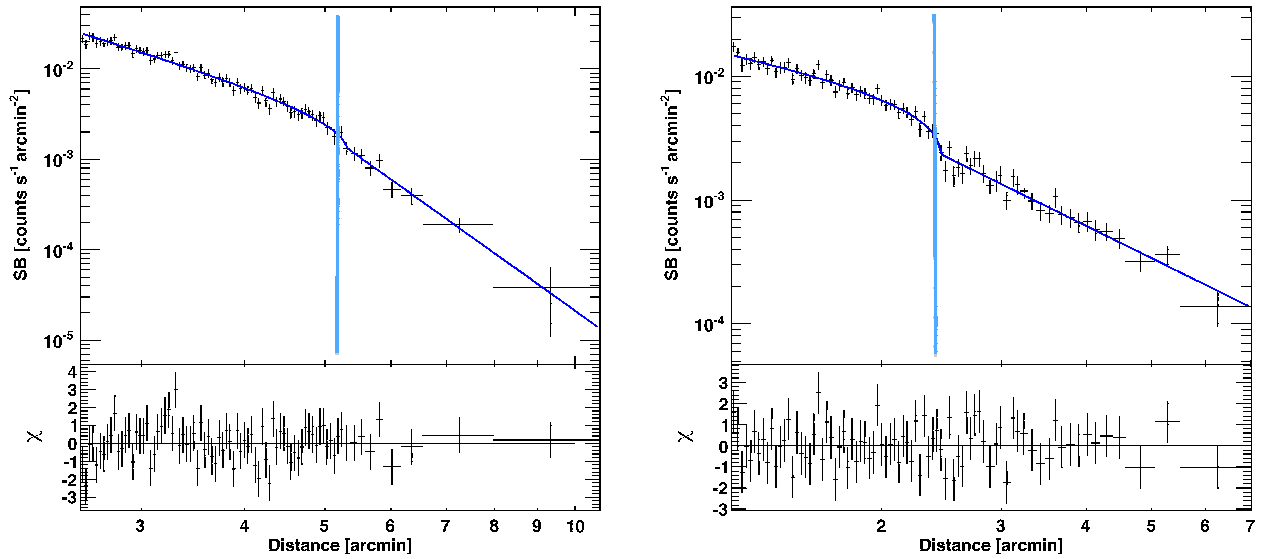
The distance between a shock and the blunt body producing the shock (i.e., the remnant subcluster’s core, Vikhlinin et al. 2001) depends on the Mach number. Applying this to the Toothbrush cluster, the distance ratio between shock and cold front location ( $\approx 2$ ) is consistent with our obtained Mach number (Equation 11, Farris & Russell 1994). This confirms the association between the southern subcluster’s cold front and the shock ahead of it. We can also use the Mach angle ( $\mu$ ) to compute the Mach number of the shock, via the relation  $\sin(\mu) = \frac{1}{\mathcal{M}}$ . From Figure 9, we estimate a Mach angle  $\mu \approx 40^\circ$  which corresponds to  $\mathcal{M} \approx 1.5$ , again consistent with the previous results.

#### 4.3. Northern Sector

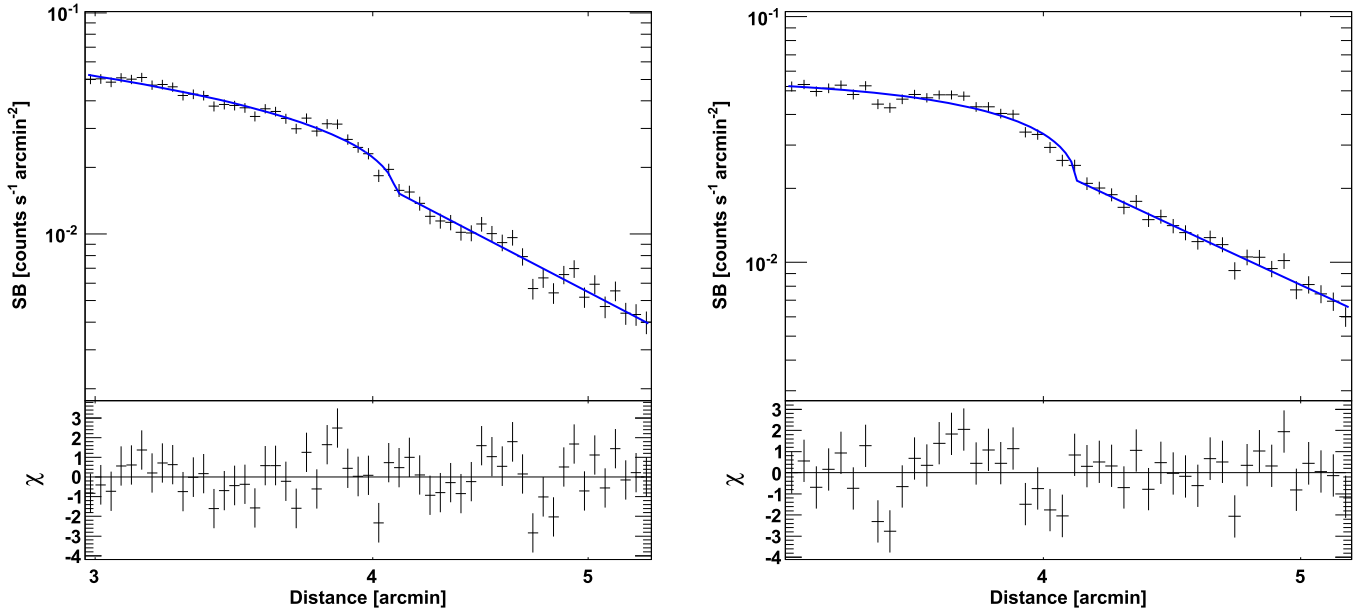
In addition to the three sectors in the southern part of the cluster, we also extracted a profile across the northern radio relic (B1) to search for a possible edge associated with the front of the relic, see Figure 12. The sector shape is defined by the curvature and shape of the radio relic. The resulting profile is shown in left panel of Figure 13. Interestingly, the profile indicates a change in the slope of the surface brightness profile around the location of the radio relic. We fitted a broken power-law density model to this profile and obtained  $C = 1.26^{+0.14}_{-0.10}$ , with the model providing a good fit to the data. The position of the jump ( $r_{\text{edge}}$ ) is also in good agreement with the location of the front of the radio relic. However, at the 90% confidence limit we cannot exclude the possibility of having a kink in the profile (i.e.,  $C = 1$ ). Better statistics will be needed to determine if a jump is really required.

For the northern partial annuli across relic B1, we measure  $T_{1N} = 8.3^{+3.2}_{-2.4}$  keV and  $T_{2N} = 8.2^{+0.7}_{-0.9}$  keV. Clearly, the error on the putative pre-shock temperature is too large to confirm or rule out a shock, which is expected given the low Mach number and the low count rate in the pre-shock region. Taking the  $1\sigma$  lower and upper limit for the pre and post-shock temperatures, respectively, we obtain an upper limit on the Mach number of 1.5.

Interestingly, both *Chandra* and *XMM-Newton* data show that any shock that traces the Toothbrush relic (the B1-part) must have a very low Mach number, otherwise the jump should have been detected unambiguously. To determine an upper limit on the Mach number, we refitted the *Chandra* profile keeping the location of the shock fixed to the value reported in Table 3 (which coincides with the front of relic B1). From this we determined an upper limit of  $C = 1.7$  at the 90% confidence level. Thus we find that the putative shock must have a Mach number of  $\mathcal{M} \leq 1.5$ . This number is consistent with the upper limit we found from the temperature change across the relic.



**Figure 13.** Left: X-ray surface brightness profile across the bright part (B1) of the northern Toothbrush relic. Right: X-ray surface brightness profile in the southern sector. The blue lines show the best-fitting broken power-law density model (Equation (1)) and the vertical line the outer boundary of the radio relic or halo emission.



**Figure 14.** Left: X-ray surface brightness profiles for the right-angled edges around the southern cluster's core. The profile for the eastern sector is shown in the left panel and the western one in the right panel, see Figure 12. The blue lines show the best-fitting broken power-law density model (Equation (1)).

#### 4.4. Comparison with XMM-Newton Results

We compare the above results with previously published XMM-Newton data by Ogrea et al. (2013). XMM-Newton also found a shock in the southern part of the cluster. The Mach number measured by Chandra via the density jump is consistent with the  $\mathcal{M} = 1.5 \pm 0.1$  reported by Ogrea et al., but the actual position of the shock does not agree with the position found by Chandra. After a closer investigation, we conclude that small-scale substructure in the downstream region of the ICM appeared as an edge in XMM's surface brightness profile. Together with XMM's lower spatial resolution this caused the fitting procedure to pick up the wrong shock location. We also note that the sector in which the surface brightness profile is fitted is different from the one used

by Ogrea et al. (2013). Since the edge is not clearly visible in the XMM image by eye there was no clear guidance on what sector to pick in this case. However, if we take the position of the shock found by Chandra as an initial guess and re-fit the XMM data, we obtain consistent results for  $r_{\text{edge}}$ .

Turning to the northern subcluster, a less significant edge was reported to be located about  $1'$  to the north of relic B1 in the XMM image (see Figures 8–11 from Ogrea et al. (2013)). At that location there is no jump visible in the Chandra profile. Given the rather low-significance ( $< 2\sigma$ ) of this jump its status remains unclear. We also note that the actual significance of this jump is lower than the fitting results suggest due to the look-elsewhere effect.

Ogrea et al. (2013) did not find a shock at the front of relic B1, although a slightly different sector was used to extract the

surface brightness profile. We refitted the XMM profile using the sector shown in Figure 12 with a broken power-law model. The initial guesses for the fitting were taken from the best fit *Chandra* model (Figure 13). In this case, the fitted XMM profile was consistent with the results from *Chandra* and also revealed a change in the slope of the surface brightness profile at the front of relic B1.

We conclude that both the *XMM-Newton* and *Chandra* data do not provide substantial evidence for a “relic shock offset problem.” Instead, a change in the slope of the surface brightness profile is observed at the expected location of a shock that would trace relic B1. The best-fitting broken power-law model indicates a small jump in the profile, corresponding to  $C = 1.26^{+0.14}_{-0.10}$ , at the location where the slope changes. However, the current *Chandra* data are not deep enough to fully exclude the possibility that this feature is just a change in the slope of the profile and not a jump.

## 5. DISCUSSION

### 5.1. Radio Spectral Index and Shock Mach Number for Relic B1

For a shock front, DSA (in the linear test particle regime) relates the shock Mach number to the injection spectral index  $\alpha_{\text{inj}}$  (e.g., Blandford & Eichler 1987)

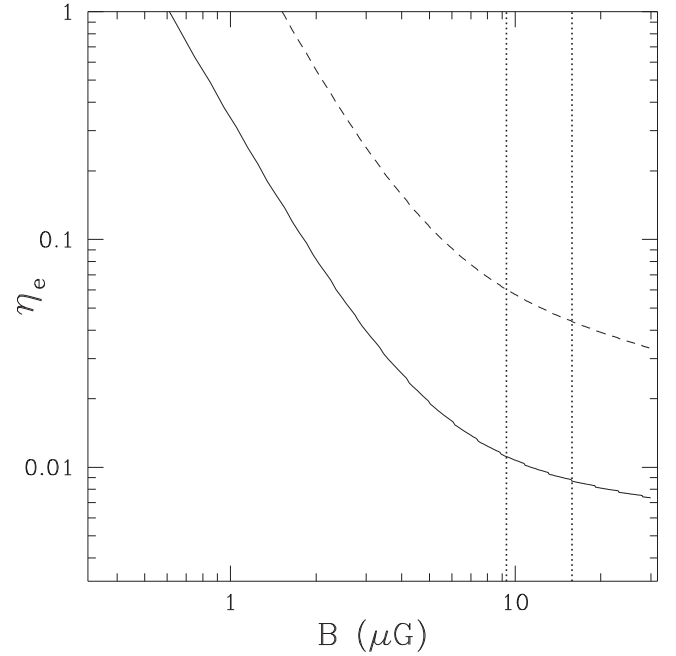
$$\alpha_{\text{inj}} = \frac{1}{2} - \frac{\mathcal{M}^2 + 1}{\mathcal{M}^2 - 1}. \quad (4)$$

The injection spectral index is related to the energy spectrum of the electrons via  $\alpha_{\text{inj}} = -(\delta_{\text{inj}} - 1)/2$ , with  $dN/dE \propto E^{-\delta_{\text{inj}}}$ . The volume integrated spectral index,  $\alpha_{\text{int}}$ , is steeper by 0.5 units compared to  $\alpha_{\text{inj}}$  for a simple planar shock model where the electron cooling time is much shorter than the lifetime of the shock (Ginzburg & Syrovatskii 1969). Recently, it has been shown that this approximation ( $\alpha_{\text{inj}} = \alpha_{\text{int}} + 0.5$ ) does not hold for spherically expanding shocks with properties comparable to those found in clusters (Kang 2015a, 2015b). Therefore, if possible, it is preferred to take the spectral index from spatially resolved maps to obtain  $\alpha_{\text{inj}}$  and avoid using  $\alpha_{\text{int}} + 0.5$  in combination with Equation (4).

For the Toothbrush relic (source component B1, Figure 1) with  $\alpha_{\text{inj}} = -0.8 \pm 0.1$ , we obtain  $\mathcal{M} = 2.8^{+0.5}_{-0.3}$ , which is not consistent with the limit of 1.5 set on the Mach number by the density jump. Working the other way around, i.e., starting with a Mach number of 1.5 (to take the most extreme value), we calculate  $\alpha_{\text{inj}} = -2.1$ , which is in strong conflict with the radio spectral index maps.

One could argue that the limit on the Mach number obtained from the density jump is not reliable, given the assumption of spherical symmetry and unknown radius of curvature. However, the Mach number derived via the temperature jump should be less sensitive to these effects. Given that we obtained a limit of about  $\mathcal{M} \leq 1.5$ , from the temperature jump we conclude that it is unlikely that there could be a strong well-defined  $\mathcal{M} = 2.8$  shock hiding in the data.

Recent PIC simulations indicate that particles are initially accelerated via SDA (also a Fermi I process) and not DSA (Guo et al. 2014a, 2014b) at cluster shocks. However, these simulations suggest that the resulting radio spectral index will be similar to that created by DSA.



**Figure 15.** The fraction of the kinetic energy flux at the shock ( $\eta_e$ ) that needs to be channelled into the supra-thermal and relativistic electron components that is required to explain the observed flux density of relic B1. The solid line is for  $\delta_{\text{inj}} = 2.2$  and the dashed line for  $\delta_{\text{inj}} = 2.6$ . For any reasonable value of the magnetic field strength ( $B \lesssim 10 \mu\text{G}$ ) we find the extreme requirement that more than several percent of the kinetic energy flux crossing the shock should be converted into the acceleration of relativistic electrons. The two vertical dotted lines indicate the magnetic field value for plasma  $\beta$  values of 3 and 1.

We therefore conclude that using the radio spectral index and Equation (4) leads to an overestimation of the X-ray-predicted Mach number for the Toothbrush relic, considering in particular that we took the most “extreme scenario,” i.e., the steepest possible injection spectral index and the largest Mach number allowed by the X-ray surface brightness profile.

### 5.2. Acceleration Efficiency

If the Mach number is indeed low, as suggested by the X-ray observations, another important consideration is the electron acceleration efficiency required to produce such a bright radio relic (Vazza & Brügggen 2014). The kinetic energy flux that becomes available at the shock ( $\Delta F_{\text{KE}}$ ) is:

$$\Delta F_{\text{KE}} = 0.5 \rho_1 v_s^3 \left(1 - \frac{1}{C^2}\right), \quad (5)$$

where  $v_s$  is the shock speed and  $\rho_1$  is the upstream density. Part of this energy can be converted into electron acceleration and then into synchrotron and IC emission.

In Figure 15, we show the efficiency of electron acceleration that is necessary at the shock to explain the observed synchrotron luminosity (Brunetti & Jones 2014) of the “bright” part of the relic, component B1. The efficiency,  $\eta_e$ , is defined as the fraction of the kinetic energy flux at the shock that is channelled into the supra-thermal and relativistic electron component, i.e.,:

$$0.5 \eta_e \rho_1 v_s^3 \left(1 - \frac{1}{C^2}\right) = \epsilon_e v_2 \quad (6)$$

where  $\epsilon_e$  and  $v_2$  are the energy density of accelerated electrons and the downstream velocity of the flow, respectively.<sup>31</sup>

We assume a power-law distribution of the accelerated electrons in momentum  $Q(p) \propto p^{-\delta_{\text{inj}}}$  with two relevant values of the injection spectral index, one consistent with the synchrotron injection spectrum measured at the northern edge of the relic ( $\delta_{\text{inj}} = 2.6$ ) and another one that is consistent with the integrated spectrum of the relic,  $\alpha_{\text{int}} = 1.1$  (i.e.,  $\alpha_{\text{inj}} = -0.6$ , resulting in  $\delta_{\text{inj}} = 2.2$ ; van Weeren et al. 2012a). The required efficiency  $\eta_e$  decreases as the strength of the magnetic field in the relic region (downstream region) is increased, as a progressively larger fraction of the shock energy flux is radiated into the radio band via synchrotron emission; large values of the magnetic fields,  $B \gtrsim 10 \mu\text{G}$ , are however excluded under the assumption that thermal pressure is larger than the magnetic field pressure downstream.

From Figure 15 we come to the extreme requirement that more than several percent of the kinetic energy flux crossing the shock should be converted into acceleration of relativistic electrons (e.g., Kang et al. 2012). This argues against standard DSA and provides support for re-acceleration for which this requirement does not apply.

### 5.3. Shock Re-acceleration

In the re-acceleration scenario, the post-shock radio spectral slope is also given by Equation (4), unless the initial fossil spectrum has a flatter spectrum than what would be produced by Equation (4). Therefore, we could explain the Toothbrush spectrum if we start with a fossil population with slope of  $-0.8$  (i.e.,  $\delta = 2.6$ ). This slope is then preserved during re-acceleration because the shock is too weak to change the slope of the particle distribution. We note that a spectral index of  $-0.8$  is very typical for radio galaxies. At low enough frequencies, these particles are not yet affected by spectral ageing (and resulting change in slope). The peculiar brightness distribution of the Toothbrush relic would then reflect the distribution of fossil plasma electrons in this region. While in this case the slope of the spectrum is preserved, the normalization will increase by about a factor of about

$$\frac{3C}{C + 2 - \delta(C - 1)} \quad (7)$$

at the shock (Markevitch et al. 2005).

We can estimate this number to check how bright this fossil plasma would be in comparison with the relic. For  $\delta = 2.6$  (a spectral index of  $-0.8$ ) and a density compression factor of  $C = 1.26$ , the normalization should increase the radio surface brightness by about a factor of 1.5 at the shock. A higher compression factor of  $C = 1.7$  (i.e.,  $\mathcal{M} = 1.5$ ) would result in an increase in surface brightness at the shock of about 2.7.

Besides re-acceleration, there should also be a boost of the fossil electron spectrum due to the magnetic field compression. This will increase the normalization by a factor of about

$$C^{(\delta+2)/3}. \quad (8)$$

With  $\delta = 2.6$  and  $C = 1.7$ , this gives a boost in the surface brightness of a factor of 2.3. It is not clear though whether this boost can be added to the one given by Equation (7) (Markevitch et al. 2005). However, in both cases the boost is

rather small. Therefore, if there is any pre-shock fossil plasma, it should only be a factor of  $\lesssim 5$  times fainter in comparison with the relic. There is a hint of some faint emission just to the north of B1, but this emission has a surface brightness that is a factor of  $\sim 10^2$  fainter than the northern part of relic B1. In addition, for the western part of the relic (B2 and B3) we do not detect pre-shock fossil plasma, while we have sufficient sensitivity to do so if it were present. This apparent lack of pre-shock fossil plasma can be explained if we assume that the fossil plasma was confined into an active galactic nucleus (AGN) lobe or radio tail that was stripped, accelerated, and transported downstream of the shock. In that case, the relic directly traces the remnants of a radio galaxy whose plasma has been re-accelerated. This scenario has also been invoked for the relic in the Bullet Cluster (Shimwell et al. 2015) and PLCKG287.0+32.9 (Bonafede et al. 2014). The morphology of the relic in the Bullet Cluster, and its location with respect to the X-ray emission from the ICM, is very similar to the Toothbrush cluster. However, no discrepancy between the radio and X-ray-derived Mach numbers was found by Shimwell et al. (2015).

A candidate that could have supplied this fossil plasma is the AGN located at R.A. 06<sup>h</sup>03<sup>m</sup>07<sup>s</sup>.5 Decl. +42°16'21".4, at the southwestern end of B1 (Figure 5). This inverted spectrum radio source ( $S_{610\text{MHz}} = 2.0 \pm 0.2 \text{ mJy}$ ) is associated with a cluster member ( $z = 0.210$ ; W. A. Dawson et al. 2015, in preparation). The source is unresolved at the  $\approx 4''$  resolution of the GMRT 610 MHz observations (van Weeren et al. 2012a), but it is located within the overall emission associated with B1. While the source is rather faint compared to the relic, it may have been brighter in the past. However, given the spatial extent of the relic the proximity could also be circumstantial.

A problem with the above scenario of an AGN tail that was stripped and re-accelerated, is that one would expect to observe a spectral index steepening along the relic due the energy losses along the tail. Re-acceleration at a weak shock would preserve this spectral index trend. The spectral index along the northern edge of relic B does not show clear evidence of spectral steepening from west to east (Figure 6). However, it should be noted that the above sketched re-acceleration scenario assumed a power-law fossil electron distribution. If the fossil electrons originated from a radio galaxy, the assumption of a power-law distribution is not valid, as the distribution is expected to have a cutoff due to energy losses. In this case, a weak shock can change the observed spectral index (e.g., Enßlin & Gopal-Krishna 2001). Simulations will be needed to determine if this can explain the observed spectral index for the Toothbrush relic (e.g., Kang & Ryu 2015).

### 5.4. A Mixture of Mach Numbers

Shocks tracing radio relics should consist of a mixture of Mach numbers along the line of sight, as hydrodynamical simulations indicate (e.g., Skillman et al. 2011, 2013; Vazza et al. 2012). If the acceleration efficiency is a nonlinear function of the Mach number (e.g., Hoeft & Brüggem 2007), the synchrotron-weighted Mach number would naturally be higher than the X-ray weighted Mach number. In addition, if the X-ray Mach number is measured via the density jump, then the projection of different shock surfaces and shock substructures would smear out the density jump, increasing the discrepancy between the radio- and the X-ray-derived Mach numbers.

<sup>31</sup> Note however that the efficiency in many cases is defined without the term  $(1 - \frac{1}{C})$  (e.g., Ryu et al. 2003; Kang et al. 2007).

Recent simulations of particle acceleration at cluster shocks by Hong et al. (2015) provide some support for this explanation. The Mach numbers derived from mock radio and X-ray maps were found to differ significantly. The X-ray observations were found to underestimate the volume weighed shock Mach number. The derived mock radio Mach numbers gave a fairer representation of the volume averaged Mach number of shocks associated with radio relics. Therefore, (part of) the difference between the radio- and X-ray-derived Mach numbers in the Toothbrush cluster could also be explained by a mixture of Mach numbers along the line of sight. It should be noted though that these simulations did not take into account the fact that pre-existing CR electrons could be re-accelerated at weak shocks, which could influence the results. One would also expect spectral index variations along the length of relics. For the Toothbrush relic, we do observe some spectral index variations along the northern edge of the relic (Figure 6). It is not clear though if these variations are consistent with these simulations, because Hong et al. (2015) only reported on the integrated spectral indices for radio relics. Future deep X-ray and radio observations are critical to determine whether discrepancies between radio- and X-ray-derived Mach numbers of shocks tracing relics are common.

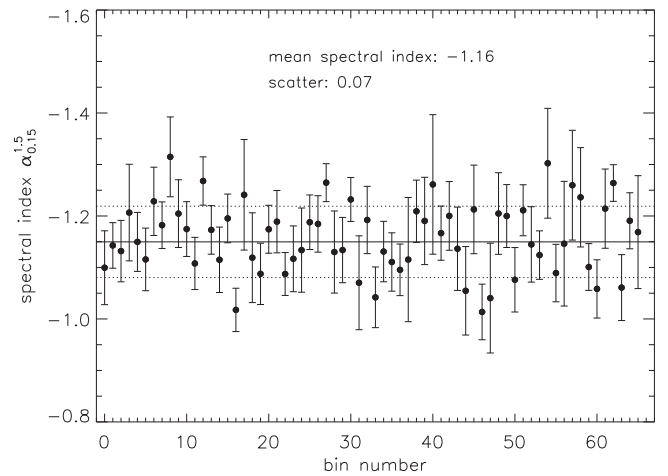
### 5.5. Radio Halo Spectral Index Variations

To study the spectral index across the radio halo we extracted spectra in several regions, where the halo is well detected by both LOFAR and the VLA, as indicated in Figure 8. We avoided the extreme southern part of the radio halo due to the presence of several compact sources. The regions where the spectra were extracted correspond to boxes with physical sizes of about 110 kpc on a side.

The spectral indices for the remaining 66 regions are displayed in Figure 16. From the 66 spectra we find that the spectral index across the radio halo over a  $0.8 \text{ Mpc}^2$  region stays rather constant. Following Cassano et al. (2013), to properly take the measurement uncertainties into account, we compute a raw scatter from the mean spectral index of 0.07. This implies that the intrinsic spectral index scatter is  $\leq 0.04$ , given that measurement errors from map noise alone would result in a spread of  $0.06 \sqrt{0.06^2 + 0.04^2} \approx 0.07$ . The largest intrinsic variations are  $\Delta\alpha \sim 0.1$ . This statement applies to the part of the radio halo that is detected at high S/N at both 1.5 and 0.15 GHz, but we note that this already covers regions that in “deep” old-VLA or GMRT observations would have been to be close to the detection limit.

The uniform spectral index is quite remarkable. In the case of Abell 2256, inhomogeneous turbulence was invoked to explain the low frequency spectral steepening (van Weeren et al. 2012b). In addition, Vacca et al. (2014), Orrú et al. (2007), Feretti et al. (2004) find evidence for spectral variations across radio halos that are larger than for the Toothbrush. However, it should be mentioned that the spectral index maps presented there have considerably larger spectral index uncertainties and less dense uv-coverage.

For the Toothbrush cluster the turbulence seems to be homogeneous. In the context of turbulent re-acceleration models, it would imply that the combination of turbulent energy flux and turbulent damping on relativistic electrons is essentially constant when averaged on 110 kpc cells and integrated along the line of sight. It also implies that the turbulent energy does not change significantly across the



**Figure 16.** Spectral index distribution across the radio halo between 0.15 and 1.5 GHz in regions measuring  $31 \text{ arcsec}^2$  (see also Figure 7). The regions where the spectral indices are extracted, are indicated in the left panel of Figure 8 in blue. The solid horizontal line indicates the average spectral index. The dotted lines represent the measured scatter of the spectral indices. Systematic uncertainties in the flux-scale were not included in the error bars as they do not affect the measured scatter. The region (bin) numbering was randomized.

overall region  $0.8 \text{ Mpc}^2$  region, that is the region just crossed by the two DM cores. In principle these constraints are very important for the physics of turbulence and acceleration in the ICM, however detailed simulations including physics of electron acceleration and turbulence are necessary to draw meaningful conclusions.

A region with a somewhat steeper spectrum,  $\alpha \simeq -1.4$  to  $-1.5$ , is located at the eastern part of the halo, toward relic E (Figure 8). This steep-spectrum region may indicate a decline of the magnetic field or a significant steepening of the particle spectrum (Brunetti et al. 2001). However, it could also be that this steepening is due to electron energy losses downstream of relic E.

The spectral index variations we find for the halo are much smaller than those reported by Orrú et al. (2007) for Abell 2744. For Abell 2744, the spectral index values varied between  $-0.7 \pm 0.1$  and  $-1.5 \pm 0.2$ . For Abell 2744, Orrú et al. (2007) also reported a correlation of spectral index with the ICM temperature, where the regions with the flattest spectra corresponded to regions with the highest temperature. We do not find any evidence for this correlation in RX J0603.3+4214. A similar result was reported by Vacca et al. (2014) for Abell 520. Either the physical conditions related to particle acceleration in Abell 2744 are different, or these variations are caused by the lower S/N of the Abell 2744 observations. We leave a detailed investigation of the physical implications of the uniform spectral index and temperature distribution for future work.

### 5.6. Merger Scenario

One of the puzzling aspects of the Toothbrush cluster is the very long linear extension of relic B to the east. Although the complex shape of the relic could reflect the distribution of fossil plasma in the re-acceleration scenario, the merger still needs to produce a shock that re-accelerates this fossil plasma along the visible extent of the relic. Simulations of Brügggen et al. (2012) suggest that this extension could be caused by a merger

between two main subclusters and a third lower mass system. In the *Chandra* image, we find a smaller substructure to the west of the cluster (Figures 10 and 12). Assuming it is at the same redshift as the main cluster, we compute an X-ray luminosity of  $L_{X,0.5-2.0\text{ keV}} \approx 5 \times 10^{42} \text{ erg s}^{-1}$ . Using the  $L_X - M$  scaling relation from Vikhlinin et al. (2009) this would correspond to a mass of  $\sim 6 \times 10^{13} M_\odot$  which is similar to the mass used by Brüggén et al. (2012) of  $3.5 \times 10^{13} M_\odot$ . The current location of the substructure is, however, very different from that suggested by the simulations of Brüggén et al. (2012). New simulations, taking the position of this subcluster into account, as well as weak lensing and spectroscopic measurements that provide estimates of the subcluster masses (W. A. Dawson et al. 2015, in preparation; Jee et al. 2015), should shed more light on this issue.

### 5.7. Radio Halo, Relic, Shocks and Turbulence Connection

The southern edge of the radio halo, C1, coincides with the location of the shock front. A similar situation is observed in the Bullet cluster (Shimwell et al. 2014) and in Abell 520 (Markevitch et al. 2005; Markevitch 2010; Vacca et al. 2014), the Coma Cluster (Planck Collaboration et al. 2013; Uchida et al. 2015), and Abell 754 (Macario et al. 2011). Therefore, this seems to be a common phenomenon. In van Weeren et al. (2012a), we speculated whether the southern part of the radio halo (C1) could actually be a fainter relic, given the somewhat enhanced surface brightness. The rather uniform spectral index across the radio halo argues against such an interpretation.

Relic D is located at the eastern end of the southern shock. The spectral index gradient across this small relic, with spectral steepening toward the north (Section 3.2), is precisely what one would expect for a southwards traveling shock. Since the relic is quite small, a largest linear size (LLS) of 0.25 Mpc, and has a complex morphology (van Weeren et al. 2012a), we speculate that this source also traces a remnant AGN lobe where particles are re-accelerated. We could not identify the source of this fossil plasma, but there are several cluster galaxies located in its vicinity.

The radio halo seems to trace the region where the X-ray gas is highly disturbed (Figure 9), as the core of the southern subcluster flew through another cluster (that is now located north of it). Baryons in the ICM behave collisionally, and in major mergers after the core passage they are strongly stirred and partially decouple from the collisionless DM cores (e.g., Roettiger et al. 1997; Clowe et al. 2006). In this situation, fluid/kinetic instabilities and turbulence are generated which play a role in dissipating part of the kinetic energy of the gas into gas heating, particle acceleration, and magnetic field amplification.

An interesting feature is the connection between the northern relic and the halo. A strong spectral steepening is measured downstream of the relic, while the spectrum returns to being progressively flatter in the halo (Figures 7 and 8). This spectral behavior is suggestive of the interplay between different mechanisms. The most appealing hypothesis is that the flattening is due to the re-acceleration by turbulence of “aged” electrons downstream of the relic (Markevitch et al. 2005; Markevitch 2012).

On one hand, we can assume that the large scale turbulence that is responsible for the radio halo is generated by the shock passage downstream, and that it takes time to decay into the smaller scales that are relevant for particle acceleration

(Brunetti & Jones 2014). Thus, spectral flattening would occur in the region where the decay time of large-scale gas motions,  $\sim L/\delta V$  ( $\delta V$  and  $L$  being the velocity and scales of the large eddies), equals the time spent by the passage of the shock,  $\sim \text{distance}/v_2$ . This sets a condition  $\delta V \sim 700 \text{ km s}^{-1} \times L/(100 \text{ kpc})$  (assuming  $\mathcal{M} \sim 1.5$  and upstream temperature from the observations) that is qualitatively consistent with the requirements of the models of radio halos based on particle re-acceleration by compressive MHD turbulence (see Brunetti & Lazarian 2007).

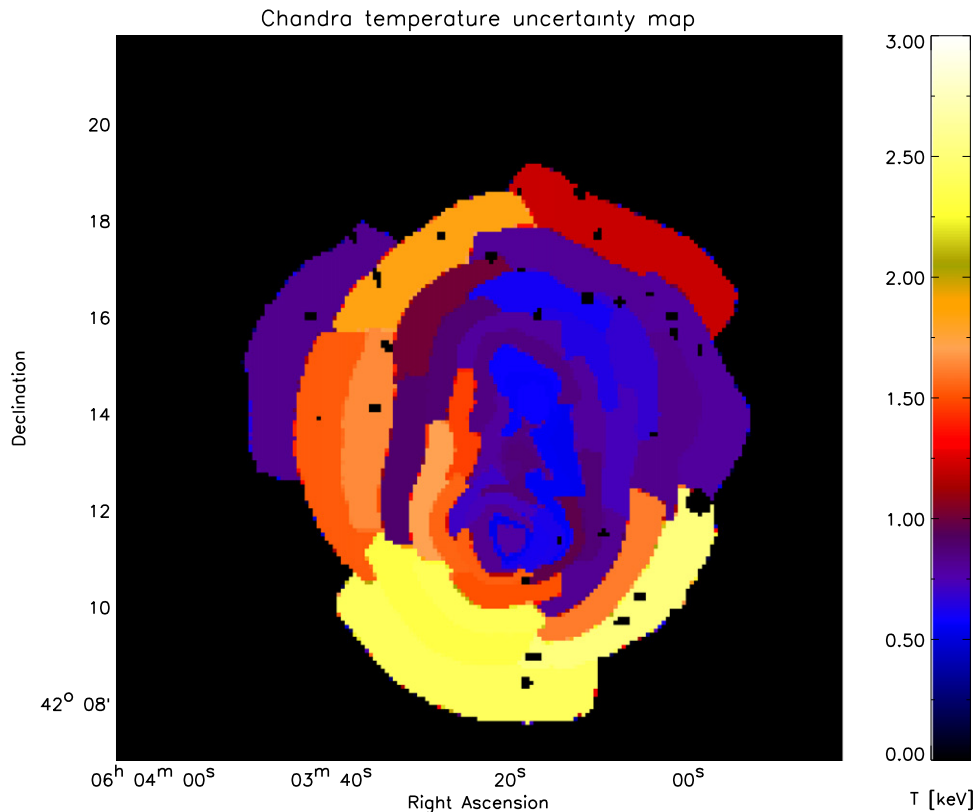
On the other hand, it is also possible that the weak shock associated with the northern relic is in fact not generating enough turbulence and that turbulence is only present in the region where we observe spectral flattening. The absence of strong spectral steepening in the southern region of the halo, where the southern shock is detected, might favor this second hypothesis. In this case the southern shock should still be traveling in a turbulent region, advecting turbulent plasma in the region associated with the passage of the two DM cores, while the northern shock should be traveling ahead of the DM core that is moving northward.

## 6. CONCLUSIONS

In this paper, we have presented results from deep LOFAR and *Chandra* observations of the Toothbrush cluster. The LOFAR data have allowed us to create high-quality spectral index maps, combining with GMRT and VLA data, and probe diffuse emission with unprecedented sensitivity in the 120–181 MHz range. To obtain these deep LOFAR images, we developed a new calibration scheme, which we name “facet calibration,” that corrects for the time-varying complex station beams and direction dependent ionospheric effects. Below, we summarize our results:

1. We obtained the first deep LOFAR images of diffuse cluster radio emission, reaching a noise level of  $93 \mu\text{Jy beam}^{-1}$ . These images allowed the study of the diffuse emission down to  $5''$ -scales for the first time below 200 MHz.
2. We detect both the radio halo and the relics. For the main northern Toothbrush relic, we observe a strong spectral index gradient. The spectral index is  $-0.8 \pm 0.1$  at the northern edge of the relic and steepens to  $\approx -2$  at the boundary of the relic and the halo. The spectral index, along the northern edge of the Toothbrush relic, varies (over a distance of 1.9 Mpc) between  $-0.7 \pm 0.1$  and  $-1.0 \pm 0.1$ , with the brighter parts corresponding to flatter spectra.
3. The radio halo has an integrated spectral index of  $\alpha = -1.08 \pm 0.06$ . The radio spectral index across the halo is remarkably uniform. For the main part of the radio halo, covering a region of  $0.8 \text{ Mpc}^2$ , we find that the intrinsic spectral index scatter is  $\leq 0.04$ .
4. With the *Chandra* observations we find a merger-related shock and cold front in the southern part of the cluster. The shock is located at the southern edge of the radio halo, and the shock/halo connection is similar to the Bullet Cluster. The southern shock has a Mach number of  $\mathcal{M} = 1.39_{-0.06}^{+0.06}$  measured via the density jump and  $\mathcal{M} = 1.8_{-0.3}^{+0.5}$  via the temperature jump. We also find evidence for variations of the Mach number along the





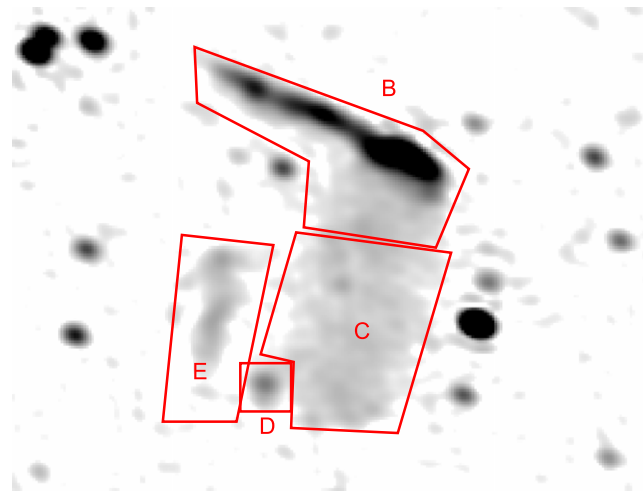
**Figure 17.** Temperature uncertainty map corresponding to Figure 11. The average of the upper and lower  $1\sigma$  temperature errors is shown.

shock front, with the Mach number ranging from  $\mathcal{M} = 1.16_{-0.07}^{+0.11}$  to  $\mathcal{M} = 1.57_{-0.13}^{+0.16}$ .

5. We observe a change in the slope of the X-ray surface brightness profile at the location of the Toothbrush relic. This change of slope could indicate the presence of a very weak,  $\mathcal{M} \sim 1.2$ , shock. The absence of a stronger shock seems to be at odds with the radio-derived Mach number of  $\mathcal{M} = 2.8_{-0.3}^{+0.5}$  predicted from DSA. Thus, the observationally derived radio and X-ray shock Mach numbers differ significantly for this radio relic. In addition, for DSA to explain the observed radio flux density of the relic, an unrealistic fraction of the shock energy is required to produce the non-thermal electron population.
6. The apparent conflict between the X-ray derived Mach number and the one from the radio spectral index can be reconciled with a scenario where a weak shock re-accelerates a population of fossil electrons with an energy slope of  $\delta \simeq 2.4$ . The complex brightness distribution of the relic suggests that the source of these fossil electrons was likely an old radio galaxy. Alternatively, the discrepancy between the radio- and X-ray-derived Mach number might be explained if the relic traces multiple shock surfaces along the line of sight, coupled with the nonlinear scaling of radio brightness with Mach number.
7. In the *Chandra* image, we detected a third, smaller, substructure. However, it is not clear if this substructure is responsible for the peculiar linear eastern extension of the main Toothbrush relic. New simulations of the merger event are required to shed more light on this issue.

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**Figure 18.** VLA 1–2 GHz images depicting the regions where the integrated flux densities were measured. The image has a resolution of  $31'' \times 24''$ .

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*Facilities:* GMRT, CXO, VLA, LOFAR.

#### APPENDIX TEMPERATURE UNCERTAINTY MAP AND INTEGRATED FLUX DENSITIES

Figure 17 displays  $1\sigma$  uncertainties in the temperature map shown in Figure 11. The regions where we measured the integrated flux densities at 0.15 and 1.5 GHz are depicted in Figure 18. The integrated flux densities are reported in Table 2.

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