

Transfer reactions in inverse kinematics: An experimental approach for fission investigations

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Inelastic and multinucleon transfer reactions between a ^{238}U beam, accelerated at 6.14 MeV/u, and a ^{12}C target were used for the production of neutron-rich, fissioning systems from U to Cm. A Si telescope, devoted to the detection of the targetlike nuclei, provided a characterization of the fissioning systems in atomic and mass numbers, as well as in excitation energy. Cross sections and angular and excitation-energy distributions were measured for the inelastic and transfer channels. Possible excitations of the targetlike nuclei were experimentally investigated for the first time, by means of γ -ray measurements. The decays from the first excited states of ^{12}C , ^{11}B , and ^{10}Be were observed with probabilities of 0.12–0.14, while no evidence for the population of higher-lying states was found. Moreover, the fission probabilities of ^{238}U , ^{239}Np and $^{240,241,242}\text{Pu}$ and ^{244}Cm were determined as a function of the excitation energy.

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I. INTRODUCTION

Multinucleon transfer reactions have been widely used for the investigation of the fission process, generally in direct-kinematics experiments with projectiles lighter than He and actinide targets. Good examples are the measurements of transfer-induced fission probabilities, which were used in the past as an experimental observable for the study of actinide fission barriers [1].

In the framework of the surrogate-reaction technique, which is discussed in Sec. I, transfer-induced fission probabilities allow for the estimation of neutron-induced fission cross sections when direct neutron-irradiation measurements are not feasible [2–5]. These measurements are important for nuclear energy applications, such as the development of new generation nuclear reactors or the recycling of radioactive waste.

In addition, the use of alternative reactions for fission investigations allows to extend the number of fissioning systems accessible to the experimental research and the investigation of fundamental properties of the fission process [6–9].

Valuable results on fission probabilities were obtained by transfer-induced fission involving heavier projectiles, which are consistent with measurements using lighter projectiles

and enlarged the accessible excitation-energy range [10]. Furthermore, experiments of this type have permitted the investigation of some aspects of fission dynamics. In particular, high excitation energies and high angular momenta [11,12] of the fissioning systems can be explored by increasing the charge and the mass transferred. Significant survival probabilities against fission have been observed, which give hints on the description of the fission times and the deexcitation of the compound nucleus [13].

The interplay between fission and survival probabilities of the produced nuclear species brings particular interest to transfer-induced fission experiments where the fission of heavy or superheavy elements may be investigated. Actually, multinucleon transfer is expected to allow the production of more neutron-rich nuclei with longer half lives [14].

Our experimental approach uses inelastic and multinucleon transfer reactions between a ^{238}U beam and a ^{12}C target for the production of the fissioning systems of interest, pushing the transfer-induced fission method towards heavier transfer reactions. In this way, a single experiment gives access to a higher variety of neutron-rich actinides and allows to explore different excitation energy regimes, depending on the transferred nucleons. In addition, the use of a heavy beam and a light target define an inverse-kinematics scenario, in which the fission fragments are emitted in the forward direction with relatively high kinetic energies. Using a magnetic spectrometer, the accurate isotopic identification of the heavy and light fission fragments is possible. Results regarding isotopic fission fragment yields from an earlier experimental campaign can be found in Refs. [15,16].

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The present work aims to investigate the potential of $^{238}\text{U} + ^{12}\text{C}$ transfer reactions for fission investigations, by providing a detailed characterization of the different transfer channels and a discussion of the experimental data in the framework of the surrogate-reaction technique. The challenges associated with the detection of the relatively heavy targetlike nuclei at high intensities and the use of inverse kinematics will be discussed. The latter, while improving the quality of the identification of the fission fragments [6], results in a degradation of the excitation-energy resolution. However, it also ensures cleaner experimental conditions than direct kinematics, where reactions on the backing of the actinide targets usually complicate the analysis of the experimental data and, in particular, the measurement of fission probabilities [3].

The experimental setup and data analysis are discussed in Secs. II and III. Section IV is dedicated to the characterization of $^{238}\text{U} + ^{12}\text{C}$ channels, providing cross sections and angular and total excitation-energy distributions. The excitation of the targetlike transfer partners is discussed in Sec. V, in the light of γ -ray measurements. These results are especially important for surrogate-reaction fission experiments, as the excitation of targetlike nuclei influences the excitation energy of the fissioning system. Fission probabilities are given in Sec. VI for ^{238}U , ^{239}Np , $^{240,241,242}\text{Pu}$, and ^{244}Cm , as a function of the total excitation energy. The latter is a short-lived ($T_{1/2} = 18$ yr) minor actinide, for which transmutation in accelerator-driven systems is nowadays a subject of recognized interest [17,18]. The present work enlarges the known excitation-energy range for the investigation of the fission probabilities of this particular nucleus. Finally, a summary and conclusions are presented in Sec. VII.

A. The surrogate-reaction technique

Following the Bohr hypothesis [19], the surrogate-reaction method considers the formation and decay of the compound nucleus independent of each other. Therefore, in order to obtain information on its decay process, the compound nucleus of interest can be produced via a surrogate reaction, which is experimentally accessible. In fission investigations, inelastic and transfer reactions between nuclei in the actinide region and light nuclei are typically used.

The majority of surrogate applications invoke approximations, such as the Weisskopf-Ewing limit, where the decay branching ratios for the compound nucleus, formed with certain excitation energy, angular momentum, and parity, are only a function of the excitation energy; or assume that both the desired and the surrogate reactions populate similar angular-momentum distributions in the compound nucleus [20,21]. Under these conditions, the neutron-induced fission cross section σ_{nf} can be estimated as the product of the fission probability measured in the surrogate reaction P_f , and the compound-nucleus formation cross section σ_n^{CN} , which is calculated via an optical potential, i.e.,

$$\sigma_{nf}(E_n) = \sigma_n^{CN}(E_n)P_f(E_x). \quad (1)$$

In this expression, E_n represents the kinetic energy of the incident neutron and E_x is the excitation energy of the

compound nucleus. Both quantities are related through the neutron separation energy and the mass number of the compound nucleus, which are denoted as S_n^{CN} and A^{CN} in the equation below:

$$E_x = E_n \frac{A^{CN} - 1}{A^{CN}} + S_n^{CN}. \quad (2)$$

Although a good agreement has been observed between surrogate results and measured neutron-induced fission cross sections [3–5], their theoretical description has been the subject of intense investigation [22], as the method is, in principle, restricted to specific conditions. The Weisskopf-Ewing approximation is only justified for high excitation energies, where the decay of the compound nucleus is dominated by statistical level densities, which have no dependence on the spin and parity [20]. In addition, different angular-momentum distributions of the produced compound nucleus are usually expected from transfer and neutron-capture reactions [20,21,23]. This conundrum becomes even more unclear due to the limited amount of available information on the angular momentum of compound nuclei populated in transfer reactions, from both theoretical and experimental sides. Selecting the recoil angles of the surrogate reactions has been a method to show the influence of the angular momentum induced in the reaction on the fission-decay channel [24].

It is also relevant to remark that the surrogate-reaction technique has been usually investigated in direct kinematics, using $^{1,2,3}\text{H}$, or $^{3,4}\text{He}$ beams and actinide targets. The information concerning heavier transfer reactions, for which higher angular momenta are expected, is limited to previous measurements of $^{232}\text{Th}(^{12}\text{C}, ^8\text{Be})^{236}\text{U}$ and $^{236}\text{U}(^{12}\text{C}, ^8\text{Be})^{240}\text{Pu}$ [10], where a good agreement with direct neutron-induced fission cross sections was obtained. A more recent $^{12}\text{C}(^{238}\text{U}, ^{240}\text{Pu})^{10}\text{Be}$ experiment at GANIL, in inverse kinematics [25], suffers from insufficient resolution in both the isotopic identification of the reaction channel and the determination of the excitation energy.

An additional issue in the context of inelastic- or transfer-induced surrogate fission experiments lies in the determination of the actual excitation energy of the fissioning system. In standard surrogate-reaction measurements, the available excitation energy is usually attributed to the heavy transfer partner, i.e., the fissioning system. This is a fair consideration for reactions involving light nuclei where the first excited states are unbound or situated at high excitation energies, such as (α, α') [26], or where the breakup of the light transfer partner can be disregarded by a geometrical adjustment of the experimental setup [1,3]. However, in the perspective of a generalization of the method to heavier transfer reactions, its comprehensive application is more intricate. Besides the direct impact on the determination of the excitation energy at which fission occurs, the excited states of the light transfer partners eventually populated may decay by nucleon emission, leading to occasional misidentification of the transfer channels. In the present work, the light transfer partners present few bound states that decay through γ -ray emission. Their observation by means of γ -ray spectroscopy allows to investigate this challenging issue of the surrogate-reaction technique.

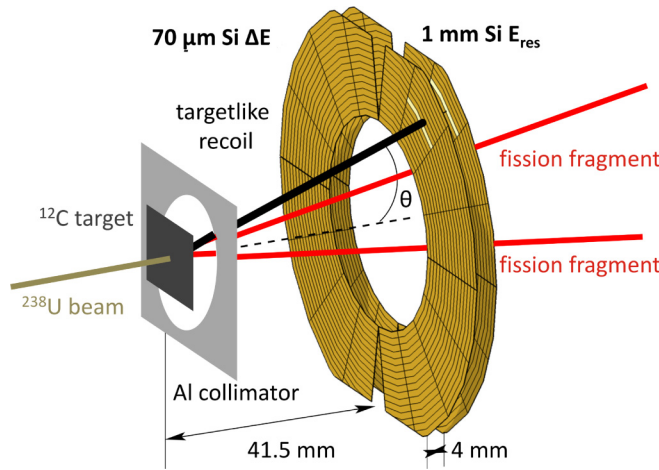


FIG. 1. (Color online) Schematic layout of the SPIDER telescope for the detection of targetlike nuclei. The beam and fission fragments passed through the inner hole of the telescope.

II. EXPERIMENTAL SETUP

A $^{238}\text{U}^{31+}$ beam with an average intensity of 10^9 pps was accelerated in the CSS1 cyclotron of GANIL up to 6.14 MeV/u. It impinged on a $100\text{-}\mu\text{g}/\text{cm}^2$ -thick ^{12}C target. As a result, inelastic scattering, fusion, and transfer reactions between the beam and the target were observed, producing a wide variety of excited actinides with a certain probability of decaying by fission. The incident energy in the center of mass, $E_{\text{c.m.}} = 70$ MeV, was approximately 10% above the Coulomb barrier. In this scenario, as has been shown in Ref. [27], fusion reactions, leading to the formation of ^{250}Cf , dominate by far the total reaction cross section, to which multinucleon transfer channels contribute with approximately 10%. The relatively low beam energy suppresses the opening of additional reaction channels.

The large-acceptance VAMOS spectrometer [28] was used for the identification of the fission fragments, described in detail in Refs. [15,29,30]. In the present work, it was used as a fission-event detector.

The detection of the targetlike nuclei was performed in a Si telescope named SPIDER, which is shown in Fig. 1. It was located 41.5 mm behind the target and covered polar angles between 30° , corresponding to the grazing angle [31], and 47° . The central hole of the detector ensured the noninterception of beamlike nuclei and fission fragments. In the inverse-kinematics conditions of the experiment, the former deviated only few degrees from the beam direction, while the latter were confined in a cone of about 25° . In order to avoid that ^{238}U nuclei of the beam halo impinged on SPIDER, a 0.5-mm-thick Al collimator, with a radius of 4 mm, was placed 3.5 mm behind the target.

The results from a previous experiment showed an increase of the current in the SPIDER detector, which reached several μA , as a consequence of the high counting rates, up to 40 kHz, of high-energy elastically scattered ^{12}C target nuclei. Therefore, the depletion of SPIDER was not permanently complete and its response dropped as a function of time. In addition, this earlier experience showed that the high counting

rates increased the temperature of the detector, deteriorating the energy resolution. These limitations were overcome in the present work by using new preamplifier concepts [32] and a cooling system, based on the circulation of liquid silicone at -35°C , which avoided the degradation of the energy resolution. Direct measurements of the temperature of the Si detectors indicated a value of $-1 \pm 1^\circ\text{C}$ during the whole experiment. Moreover, a magnetic field of approximately 750 G was used in the target region in order to prevent the arrival of δ electrons, produced in the interaction of the highly charged ^{238}U beam with the target, to SPIDER.

As shown in Fig. 1, SPIDER is composed of two double-sided Si detectors, which are 70 and $1042\text{ }\mu\text{m}$ thick. They were used to measure the energy loss ΔE and the residual energy E_{res} of the targetlike nuclei. The front and back sides of each Si detector are respectively segmented into 16 rings of 1.5 mm and 16 sectors, each covering an azimuthal, angular range of 11° . Ring (sector) sides are coated with 0.1- μm -thick Al (0.3- μm -thick Au) dead layers. The angles of the targetlike nuclei with respect to the beam direction were measured with an uncertainty below 1° thanks to the annular segmentation of the telescope.

In addition, three clovers of the EXOGAM array of Ge detectors [33] surrounded the target region. They were placed at backward angles, between 120° and 150° , the distances between the clovers and the target being 140.5, 160, and 158.1 mm. They were used in this work to investigate γ -ray emissions from the targetlike nuclei.

Finally, two main acquisition triggers were used. They corresponded to the detection of a targetlike nucleus in SPIDER, in coincidence and anticoincidence with the detection of a fission fragment in VAMOS. A reduction factor of 600 was applied to the latter in order to reduce the amount of $^{12}\text{C}(^{238}\text{U}, ^{238}\text{U})^{12}\text{C}$ elastic events treated by the data-acquisition system.

III. DATA ANALYSIS

The calibration of the SPIDER telescope was performed by means of $^{12}\text{C}(^{238}\text{U}, ^{238}\text{U})^{12}\text{C}$ elastic events. A beam energy of 6.11 MeV/u, which takes into account the slowing down in the first half of the target, was used for the kinematic calculations.

The rings and sectors of both ΔE and E_{res} detectors were individually calibrated. Up to 16 calibration points were obtained for each sector, using the coincidences with the associated rings. For each ring, only one elastic point was available and we made also use of the pedestals.

These calibrations were found to be strongly influenced by geometric parameters, such as the position of the beam or the distance between the target and the telescope, which were adjusted using a detailed simulation of elastic events in our experimental setup. This analysis showed that an accuracy of approximately 2 mm was achieved for the beam centering during the whole experiment. The beam size was of the same order, with full width at half maximum (FWHM) of approximately 2.4 and 3.5 mm in the horizontal and vertical directions, respectively.

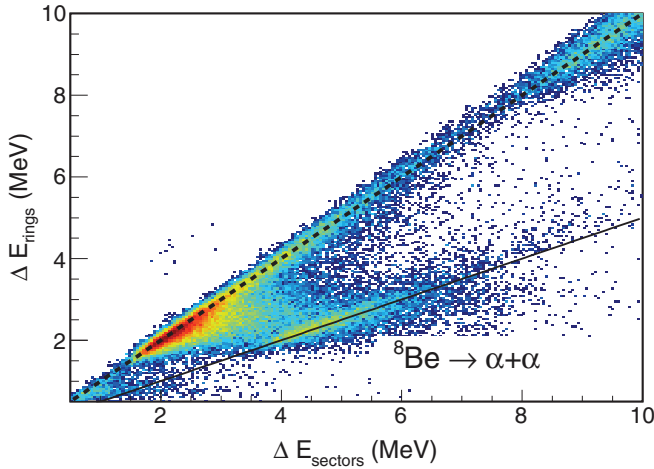


FIG. 2. (Color online) Correlation between ring and sector energy-loss measurements in the ΔE detector. Nuclei crossing a single ring of the ΔE detector are concentrated around the dashed line, for which $\Delta E_{\text{rings}} = \Delta E_{\text{sectors}}$. The solid line corresponds to $\Delta E_{\text{rings}} = 0.5\Delta E_{\text{sectors}}$. Events in this region belong to the $^{12}\text{C}(^{238}\text{U}, ^{242}\text{Pu})^8\text{Be}$ channel.

A double energy measurement was obtained in both ΔE and E_{res} detectors, as the same information was provided by ring and sector signals. For each event, the ring and sector signals were treated individually and the ones with the maximum deposited energy were selected. Geometrical considerations and the correlation between ring and sector energy measurements were used for the suppression of spurious events. Furthermore, the correlations between ring and sector energy measurements also determined the treatment of each event. Due to the thickness of the E_{res} detector, some nuclei crossed several rings before being stopped. These events were identified by using ring-sector correlations, as the energy deposited in individual rings was smaller than that measured in the sectors. An add-back procedure between adjacent rings was applied in order to recover the total energy deposited from the ring signals.

The total kinetic energy of the targetlike nuclei was finally obtained from SPIDER measurements as $E = \Delta E + E_{\text{res}} + \Delta E_{\text{Al+Au}}$. The term $\Delta E_{\text{Al+Au}}$ represents the energy loss in the Au and Al layers coating the detector sides and accounts for approximately 0.2–0.5 MeV. A precise estimation of $\Delta E_{\text{Al+Au}}$ was performed on an event-by-event basis, using the empirical formula given in Ref. [34], in which the different parameters were adjusted by means of LISE++ simulations [35].

In addition, the ring and sector energy measurements of the ΔE detector allowed to identify and to reconstruct multiplicity-2 events arising from $^{12}\text{C}(^{238}\text{U}, ^{242}\text{Pu})^8\text{Be}$ reactions and the subsequent decay of ^8Be into two strongly correlated α particles, with similar kinetic energies and a small angle between them. For those cases where the two α particles hit different rings and the same sector, a clear signature was found in the ring-sector energy correlation of the ΔE detector, which is illustrated in Fig. 2. The kinetic energy and the angle of ^8Be nuclei were reconstructed from the kinetic energy, momentum, and angle of the two α particles. They are denoted

by E_{α_i} , p_{α_i} , and θ_{α_i} in the equations below, where the angles are given with respect to the beam axis,

$$\begin{aligned} E_{^8\text{Be}} &= E_{\alpha_1} + E_{\alpha_2} + 2M_{\alpha} \\ &\quad - \sqrt{(E_{\alpha_1} + E_{\alpha_2} + 2M_{\alpha})^2 - (\vec{p}_{\alpha_1} + \vec{p}_{\alpha_2})^2}, \\ \cos(\theta_{^8\text{Be}}) &= \frac{p_{\alpha_1} \cos(\theta_{\alpha_1}) + p_{\alpha_2} \cos(\theta_{\alpha_2})}{\sqrt{(\vec{p}_{\alpha_1} + \vec{p}_{\alpha_2})^2}}. \end{aligned} \quad (3)$$

The $\Delta E - E$ correlation was used to identify targetlike nuclei according to their mass and atomic number, distinguishing the different transfer channels. The ring signals were used for this purpose because of their better energy resolution. The atomic and mass numbers of the complementary actinide partners were obtained assuming a binary reaction.

The measurement of the angle of the targetlike nucleus with respect to the beam direction completed the reconstruction of the reaction kinematics. The trajectory of the beam was considered to be perpendicular to the target and well centered with respect to the SPIDER telescope. Then, the total excitation energy in the exit channel E_x was determined by applying energy and momentum conservation laws,

$$\begin{aligned} E_x &= Q_{gg} + E_1 - E_4 - \sqrt{M_3^2 + p_3^2} + M_3, \\ p_3^2 &= p_1^2 + p_4^2 - 2p_1 p_4 \cos(\theta_4). \end{aligned} \quad (4)$$

In these expressions, Q_{gg} is the ground-state to ground-state Q value of the reaction, E_i represents the kinetic energy, and the subindexes $i = 1, 3$, and 4 refer to the beam, the beamlike, and the targetlike nucleus, respectively. The angle of the targetlike nucleus with respect to the beam direction is θ_4 . As shown in Fig. 8, an excitation-energy resolution of 2.7 MeV (FWHM) was achieved. The angular resolution, to which inverse-kinematics experiments are especially sensitive, was the main limiting factor.

The analysis of the experimental data revealed that the total excitation-energy distributions were affected by events where the targetlike nuclei hit the inside of the Al collimator behind the target, which is represented in Fig. 1. Nuclei scattered towards smaller polar angles in the collimator could reach the SPIDER telescope, biasing our reconstruction of the reaction kinematics and leading to an overestimation of the total excitation energy. For example, if ^{12}C nuclei, elastically scattered from the target, were scattered 2° with respect to their original trajectories, total excitation energies of 6 MeV would be obtained. Figure 3 provides evidence of this effect. The reconstructed total excitation-energy distributions are given for ^{12}C nuclei detected at different azimuthal angular ranges, namely $180\text{--}270^\circ$ and $270\text{--}360^\circ$. The difference between the two distributions is caused by an inaccurate positioning of the Al collimator with respect to the beam axis, so targetlike nuclei emitted at azimuthal angles of $270\text{--}360^\circ$ had a higher probability of being scattered in the inside of the collimator. As a result, an important contribution from elastic events, wrongly reconstructed, was mixed with real inelastic events. The shadow of the collimator observed in the outer rings of SPIDER during the experiment gives further support to this hypothesis.

Consequently, only a reduced range of azimuthal angles was used for the data analysis in the present work, for

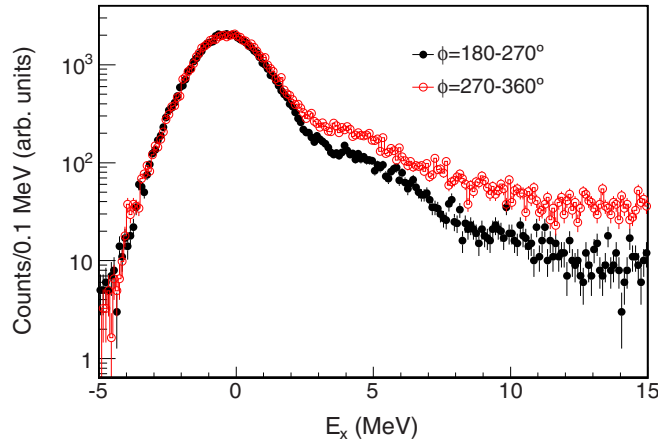


FIG. 3. (Color online) Total excitation-energy distributions reconstructed for $^{12}\text{C}(^{238}\text{U}, ^{238}\text{U})^{12}\text{C}$ events. Empty (red) and solid (black) symbols correspond to different azimuthal angular ranges. The first distribution contains a higher proportion of wrongly reconstructed events due to scattering in the Al collimator.

which no shadow of the Al collimator was observed in the detector. Background contributions to the excitation energy spectra were reduced in this way by approximately a factor of 4. An additional condition, which excluded the polar angles above 40° , was applied for the determination of the total excitation-energy distributions and fission probabilities, discussed in Secs. IV C and VI, in order to suppress residual background events.

The information provided by SPIDER was also used for the Doppler correction of the γ -ray energies performed with EXOGAM clovers. In this analysis, the targetlike nuclei, moving with typical velocities of 5 cm/ns, were considered as the γ -ray emitters.

The efficiency of the γ -ray detection was evaluated in two steps. First, the dependence on the γ -ray energy was determined from measurements performed during our experiment with a ^{152}Eu calibration source and $^{35}\text{Cl}(n, \gamma)$ data available

from an earlier work [36,37], in which γ -ray energies of up to 9 MeV were reached. Absolute efficiency values were then obtained by applying a scaling factor, based on a ^{60}Co calibration source located at the target position. A 20-Hz pulser was used in order to evaluate the dead-time contribution to the ^{60}Co measurements, which was taken into account afterwards. Absolute efficiency values of 0.014 and 0.006 were obtained at energies of 1.3 and 4 MeV, respectively. An approximate error of 6% was estimated for the efficiency values obtained by this procedure.

IV. CHARACTERIZATION OF $^{238}\text{U} + ^{12}\text{C}$ INELASTIC AND TRANSFER REACTIONS

A. Identification of the exit channels

The identification of targetlike nuclei provided by the SPIDER telescope is displayed in Fig. 4, where the energy loss in the ΔE detector is shown as a function of the total kinetic energy E . In this representation, the effective thickness crossed by targetlike nuclei emitted at different angles θ is accounted for by the factor $\cos(\theta)$. Therefore, the quantity $\Delta E \cos(\theta)$ depends only on the atomic number, the mass, and the total kinetic energy. Results are given for two trigger conditions, which select the anticoincidence and the coincidence with the detection of a fission fragment in VAMOS, as mentioned above. The latter requires a minimum excitation energy of approximately 6 MeV, corresponding to the height of the fission barrier. Hence, different relative populations of the exit channels are observed in the two cases. The anticoincidences with the detection of a fission fragment are clearly dominated by elastic events, which are suppressed when a coincidence with a fission fragment is required.

The atomic and mass numbers assigned to the targetlike nuclei were cross-checked with simulations of the energy loss in the SPIDER telescope. Graphical cuts on the identification of Fig. 4 were used for the selection of the different reaction channels. Assuming a linear dependence between $\Delta E \cos(\theta)$ and E^{-1} , the ratio $\frac{\Delta E \cos(\theta)}{C_0 E^{-1} + C_1}$ is presented in Fig. 5, in order

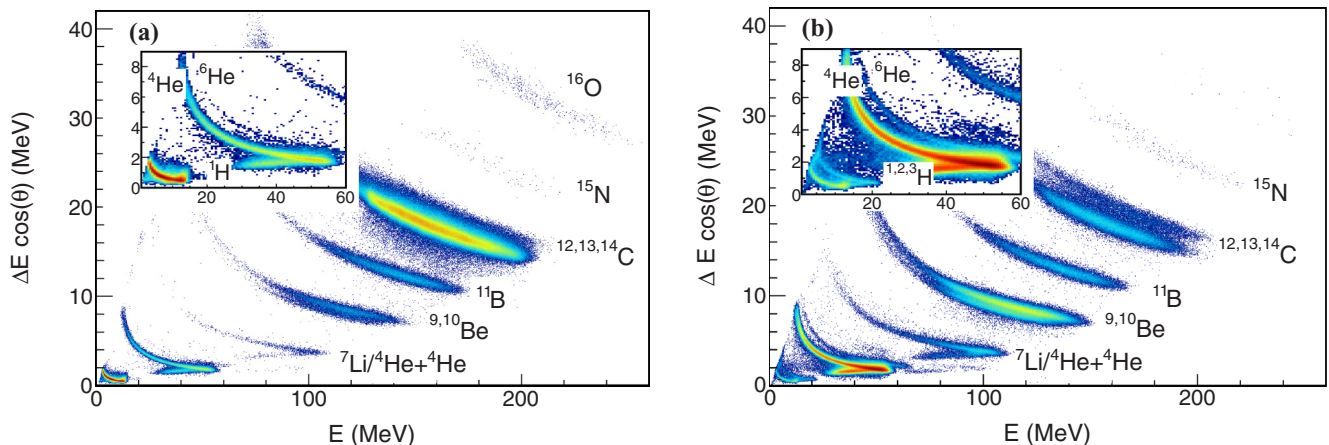


FIG. 4. (Color online) Identification of targetlike nuclei in the SPIDER telescope. The energy loss in the ΔE detector is plotted on the vertical axis. The factor $\cos(\theta)$ accounts for the different effective thickness crossed by targetlike nuclei emitted at different angles θ . The total kinetic energy is represented on the horizontal axis. Experimental data are shown in anticoincidence (a) and coincidence (b) with the detection of a fission fragment in VAMOS.

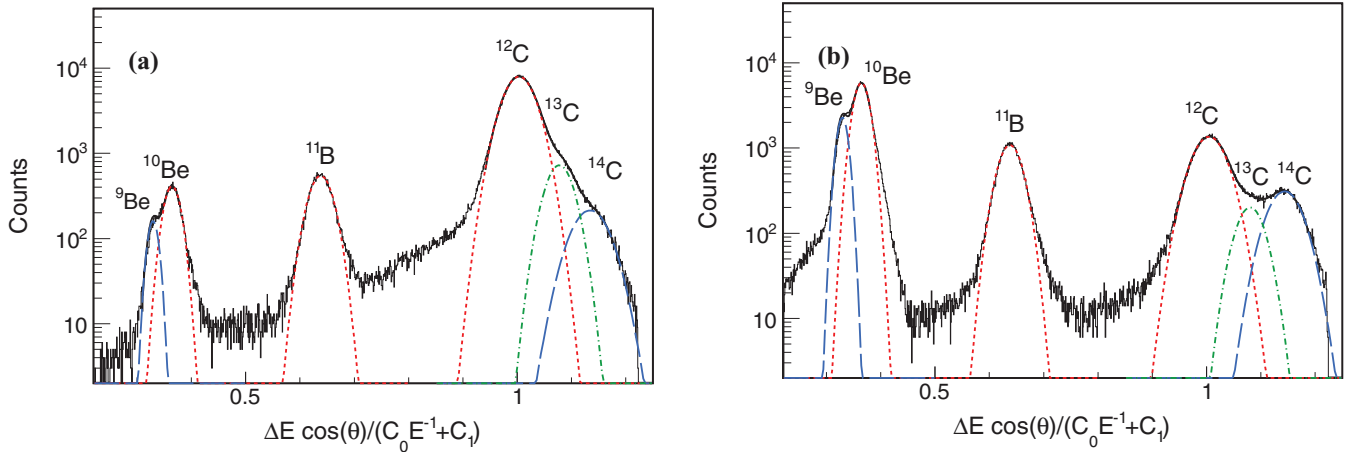


FIG. 5. (Color online) Mass separation of the target nuclei, in anticoincidence (a) and coincidence (b) with the detection of a fission fragment in VAMOS. The coefficients $C_{0,1}$ were determined by fitting the experimental data. The curves therein are the result of multiple-Gaussian fits, which were used to estimate the mixing between $^{9,10}\text{Be}$ and $^{12,13,14}\text{C}$ isotopes.

to illustrate the mass separation achieved for the targetlike nuclei. The coefficients $C_{0,1}$ were obtained from a fit of the experimental data. A mass resolution of approximately 8% (FWHM) was obtained in this work.

Multiple-Gaussian fits of the spectra given in Fig. 5 were used to estimate the mixing between $^{9,10}\text{Be}$ and $^{12,13,14}\text{C}$ isotopes. The contributions to the graphical cuts that we used to select the different targetlike nuclei are given in Table I.

The presence of ^{16}O , ^{15}N , and ^1H was attributed to $^{16}\text{O}(^{238}\text{U}, ^{238}\text{U})^{16}\text{O}$, $^{16}\text{O}(^{238}\text{U}, ^{239}\text{Np})^{15}\text{N}$, and $^1\text{H}(^{238}\text{U}, ^{238}\text{U})^1\text{H}$ reactions, as a result of contaminants in the target. The fact that ^{16}O and ^1H were mostly detected in anticoincidence with a fission fragment, together with the angular distributions measured for these nuclei, suggests that they mainly emerged from elastic reactions. The kinematic E vs θ correlations observed for these nuclei firmly support this hypothesis. Due to the importance of ^1H contaminants, the origin of the observed ^2H and ^3He remains unclear.

Nuclei from C down to He were assigned to $^{238}\text{U} + ^{12}\text{C}$ reactions, leading to the production of actinides between U and Cm. Besides elastic and inelastic scattering, up to nine transfer channels were populated. Our data indicate that nucleons essentially flow from the light to the heavy nucleus. The transfer in the opposite direction was restricted to one or two neutrons, leading to ^{13}C and ^{14}C targetlike nuclei. The presence of ^{13}C was confirmed from γ -ray spectra. Characteristic γ rays at 3684 and 3853 keV, depopulating ^{13}C excited states, are clearly visible in Fig. 10(a).

Furthermore, ^{11}B and $^{10,9}\text{Be}$ nuclei were observed. The transfer of a ^4He nucleus lead to ^8Be , which decayed into two highly correlated α particles, its half life being $T_{1/2} \sim 8 \times 10^{-17}$ s. When they hit the same ring and sector of both ΔE and E_{res} detectors, they populated the same ΔE - E region as ^7Li in the identification matrices of Fig. 4. Although transfer reactions leading to the production of ^8Be are clearly favored by their ground-state to ground-state Q value, possible ^7Li contributions could not be disentangled. For reference, ground-state to ground-state Q values of the observed channels are given in Table II.

The observed ^6He nuclei were attributed to $^{12}\text{C}(^{238}\text{U}, ^{244}\text{Cm})^6\text{He}$. Transfer reactions leading to ^7He , which would decay into $n + ^6\text{He}$ with $T_{1/2} \sim 3 \times 10^{-21}$ s, were neglected, as they would be strongly disfavored by their ground-state to ground-state Q value, $Q_{\text{gg}} = -35.95$ MeV.

The observed ^4He could originate from several channels. Multiplicity-2 events, in which two ^4He nuclei were detected by SPIDER, were unambiguously attributed to $^{12}\text{C}(^{238}\text{U}, ^{242}\text{Pu})^8\text{Be}$ reactions. An example is given in Fig. 2. However, those cases where a single ^4He nucleus was detected represent a more complicated puzzle. They could be as well $^{12}\text{C}(^{238}\text{U}, ^{242}\text{Pu})^8\text{Be}$ events, in which one α particle was lost. Although, from a geometrical point of view, the detection of the two highly correlated α particles was very probable, a limited detection efficiency for such a scenario was found during the experiment, which did not allow to exclude this hypothesis. Additional contributions could originate from $^{12}\text{C}(^{238}\text{U}, ^{246}\text{Cm})^4\text{He}$ reactions. The $^{12}\text{C}(^{238}\text{U}, ^{245}\text{Cm})^5\text{He}$ channel, where ^5He would decay into $n + ^4\text{He}$ with $T_{1/2} \sim 8 \times 10^{-22}$ s, is considered negligible, as it has $Q_{\text{gg}} = -24.93$ MeV.

B. Cross sections and angular distributions

Figure 6 shows the differential cross sections measured for $^{238}\text{U} + ^{12}\text{C}$ elastic, inelastic, and multinucleon transfer reactions, as a function of the angle of the targetlike nuclei, in

TABLE I. Contributions from neighboring isotopes in the selection of targetlike nuclei. The results are given as the proportion of the isotope of interest with respect to the total number of selected nuclei. The labels $\overline{\text{fission}}$ and fission correspond to the anticoincidence and coincidence with the detection of a fission fragment.

	$\overline{\text{fission}}$	fission
$N_{12\text{C}}/(N_{12\text{C}} + N_{13\text{C}})$	0.97 ± 0.02	0.97 ± 0.02
$N_{14\text{C}}/(N_{12\text{C}} + N_{13\text{C}} + N_{14\text{C}})$	0.43 ± 0.12	0.79 ± 0.11
$N_{9\text{Be}}/(N_{9\text{Be}} + N_{10\text{Be}})$	0.71 ± 0.06	0.81 ± 0.09
$N_{10\text{Be}}/(N_{9\text{Be}} + N_{10\text{Be}})$	0.96 ± 0.02	0.95 ± 0.02

TABLE II. Experimental $^{238}\text{U} + ^{12}\text{C}$ cross sections, σ^{expt} , integrated over the angular ranges displayed in Fig. 6. The most probable total excitation energies and the associated optimal Q values are represented by E_x^{expt} and $Q_{\text{opt}}^{\text{expt}}$, respectively. For comparison, the optimal Q values predicted by Eq. (6), $Q_{\text{opt}}^{\text{calc}}$, are indicated, as well as the incident neutron energies leading to the reported E_x^{expt} values in a neutron capture reaction E_n . The ground-state to ground-state Q values Q_{gg} are also indicated.

Reaction	Q_{gg} (MeV)	σ^{expt} (mb)	E_x^{expt} (MeV)	$Q_{\text{opt}}^{\text{expt}}$ (MeV)	$Q_{\text{opt}}^{\text{calc}}$ (MeV)	E_n (MeV)
Elastic	0.00	417 ± 52			0.00	
Inelastic, $E_x > 4$ MeV	0.00	16 ± 2			0.00	thermal
$^{12}\text{C}(^{238}\text{U}, ^{237}\text{U})^{13}\text{C}$	-1.21	30 ± 6			0.00	
$^{12}\text{C}(^{238}\text{U}, ^{236}\text{U})^{14}\text{C}$	1.85	9 ± 4	5.3 ± 0.8	-3.4 ± 0.1	0.00	thermal
$^{12}\text{C}(^{238}\text{U}, ^{239}\text{Np})^{11}\text{B}$	-10.67	22 ± 1	2.9 ± 0.1	-13.7 ± 0.1	-11.00	thermal
$^{12}\text{C}(^{238}\text{U}, ^{240}\text{Pu})^{10}\text{Be}$	-15.43	10.2 ± 0.3	9.9 ± 0.1	-25.3 ± 0.1	-22.26	4
$^{12}\text{C}(^{238}\text{U}, ^{241}\text{Pu})^9\text{Be}$	-17.00	2.2 ± 0.5	10.0 ± 0.2	-26.9 ± 0.1	-22.26	5
$^{12}\text{C}(^{238}\text{U}, ^{242}\text{Pu})^8\text{Be}$	-12.35	$>6.7 \pm 1.4$	18.8 ± 0.1	-31.5 ± 0.1	-22.26	12
$^{12}\text{C}(^{238}\text{U}, ^{243}\text{Am})^7\text{Li}$	-24.77	$<1.4 \pm 0.1$			-33.76	
$^{12}\text{C}(^{238}\text{U}, ^{244}\text{Cm})^6\text{He}$	-28.74	$>0.5 \pm 0.1$	21.3 ± 0.5	-50.0 ± 0.1	-45.52	14
$^{12}\text{C}(^{238}\text{U}, ^{246}\text{Cm})^4\text{He}$	-17.73	$<17 \pm 1$			-45.52	

the center-of-mass reference frame. The distinction between elastic and inelastic events was done by means of the reconstructed excitation energy, given in Fig. 8. A maximum excitation energy of 2 MeV was requested in the selection of the elastic channel. This limit corresponds to two times the width of the elastic peak, approximately. A correction factor of $\sim 3\%$, which accounted for the elastic events beyond this value, was considered for the determination of the cross section. In order to minimize possible contamination from the

elastic channel, only inelastic events with excitation energies above 4 MeV were considered in this analysis.

As illustrated in Fig. 6, the elastic data were normalized to the Rutherford formula at the highest scattering angles. On this basis, they were used to perform the beam normalization and obtain absolute cross sections.

In general, the inelastic and transfer channels show similar bell shapes peaked between 75 and 80° , in the center of mass. The maxima are above the calculated grazing angle,

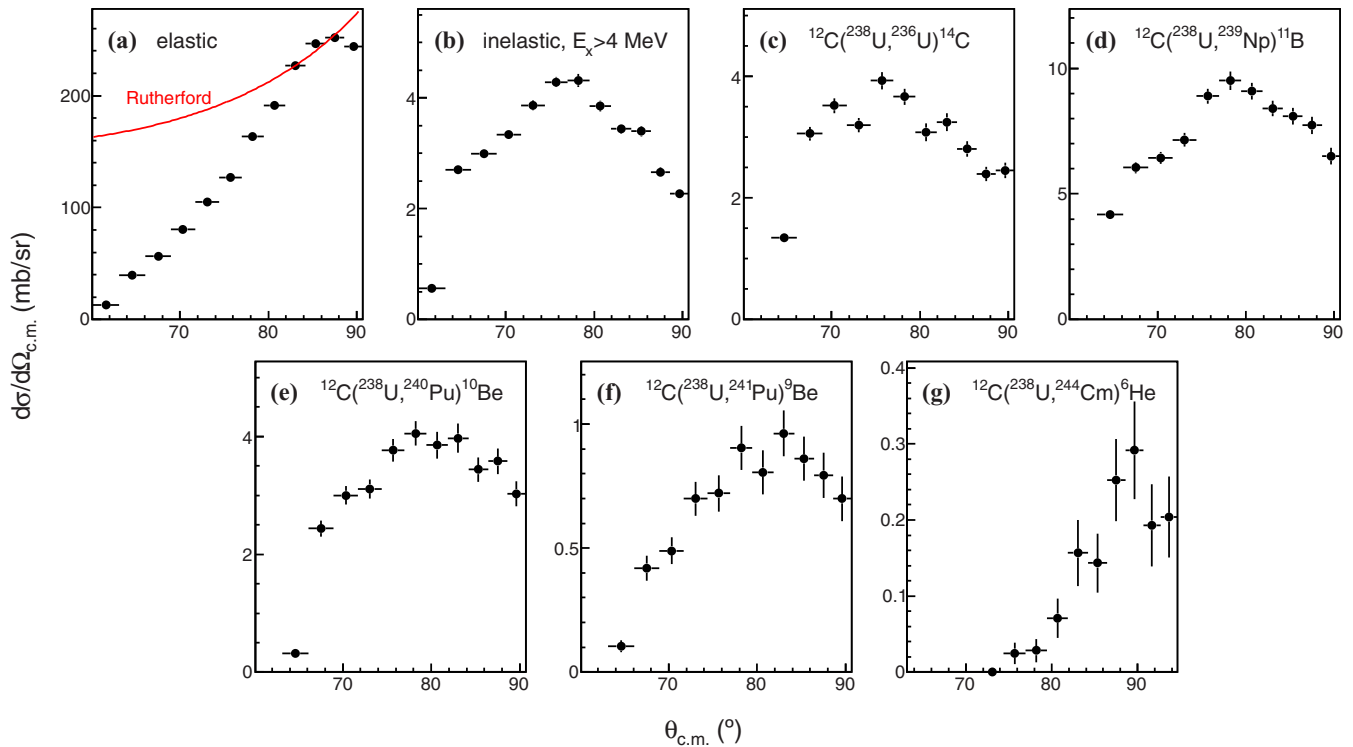


FIG. 6. (Color online) Differential cross sections measured for $^{238}\text{U} + ^{12}\text{C}$ elastic, inelastic, and multinucleon transfer reactions. The angles of the targetlike nuclei, in the center of mass reference frame, are plotted on the abscissa axis. The horizontal error bars reflect the ring size of the SPIDER ΔE detector and the vertical error bars are statistical. In (a), the Rutherford cross section is shown together with the elastic data, in order to illustrate the normalization applied to the experimental cross sections.

which has a value of 60° [31]. No clear evolution of the width of these distributions with the number of transferred nucleons was noted. Moreover, the distributions resulting from $^{12}\text{C}(^{238}\text{U}, ^{244}\text{Cm})^6\text{He}$ should be interpreted carefully. They do not include *punch-through* events, associated with ^6He nuclei that were not stopped in the SPIDER telescope, because they are mixed with ^4He in the identification matrix of Fig. 4 (inset). As these events correspond to the smallest ^6He angles, their exclusion modifies the differential cross section.

The integrated cross sections corresponding to the angular ranges displayed in Fig. 6 are given in Table II. They were corrected for the contamination between $^{9,10}\text{Be}$ and $^{12,13,14}\text{C}$ isotopes, which is represented in Fig. 5 and Table I.

The lower bound given for the $^{12}\text{C}(^{238}\text{U}, ^{242}\text{Pu})^8\text{Be}$ reactions was obtained from multiplicity-2 events, in which two ^4He nuclei were simultaneously detected by SPIDER. In our calculations, all the events lying in the region of Fig. 4 labeled as $^7\text{Li}/^4\text{He}+^4\text{He}$ were assigned to $^{12}\text{C}(^{238}\text{U}, ^{243}\text{Am})^7\text{Li}$ reactions. As possible contributions are expected from the $^{12}\text{C}(^{238}\text{U}, ^{242}\text{Pu})^8\text{Be}$ channel, incompletely reconstructed, our result must be regarded as an upper bound. In a similar way, the events associated with the detection of a single ^4He nucleus were assigned to $^{12}\text{C}(^{238}\text{U}, ^{246}\text{Cm})^4\text{He}$ reactions, despite contaminants from other channels, such as $^{12}\text{C}(^{238}\text{U}, ^{242}\text{Pu})^8\text{Be}$, are expected. Finally, the amount of ^6He nuclei that were not stopped in SPIDER could not be estimated. Therefore, the reported $^{12}\text{C}(^{238}\text{U}, ^{244}\text{Cm})^6\text{He}$ cross section should be considered as a lower limit.

The ensemble of the measured cross sections are represented in Fig. 7(a) as a function of the number of transferred nucleons. In general, they decrease exponentially with the number of transferred nucleons, following the systematics observed in earlier works [38]. The transfer of two neutrons and two protons, which leads to ^8Be and ^{242}Pu , is favored by at least a factor of 4 with respect to the reported tendency, reflecting the clustering of the transferred nucleons into an α particle. Similar results were found in Refs. [38,39] for different projectile-target systems.

The evolution of the cross sections with the ground-state to ground-state Q value, Q_{gg} , is displayed in Fig. 7(b). An exponential decrease with $-Q_{gg}$ has been observed in earlier works [38,39], although large deviations have been found and the need of a Coulomb correction term has been pointed out [38]. Systematics for the different charges are limited in our case to $^{13,14}\text{C}$ and $^{8,9,10}\text{Be}$ isotopes. The first clearly confirms that the number of transferred nucleons is a more appropriate ordering parameter for the description of the cross sections.

C. Total excitation-energy distributions

The most probable total excitation-energy, E_x , in the exit channel can be obtained as the difference between the ground-state to ground-state and an effective Q value, Q_{opt} ,

$$E_x = Q_{gg} - Q_{\text{opt}}, \quad Q_{\text{opt}} = E_{i,\text{c.m.}} - E_{f,\text{c.m.}} \quad (5)$$

In this expression, $E_{i,\text{c.m.}}$ and $E_{f,\text{c.m.}}$ are the initial and the most probable final kinetic energies for a given transfer channel, in the center of mass.

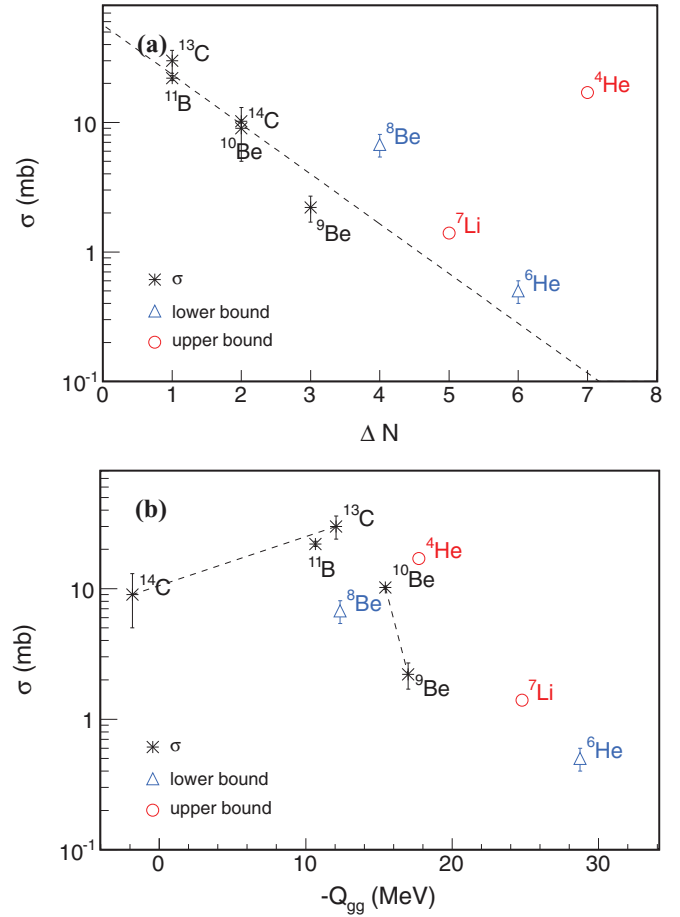


FIG. 7. (Color online) Variation of the transfer cross sections with the number of transferred nucleons (a) and the ground-state to ground-state Q value (b). In those cases where the actual cross sections were not accessible, lower (upper) bounds are represented by empty triangles (circles). The dashed line in (a) corresponds to an exponential fit of the experimental data. The lines in (b) simply indicate the evolution of the cross sections for C and Be isotopes.

If the transfer of nucleons is assumed to take place at the distance of closest approach, D , and collisions near the grazing angle are dominated by Coulomb forces, the orbit matching condition, $D_i = D_f$, leads to [40]

$$Q_{\text{opt}} = \frac{Z_3 Z_4 - Z_1 Z_2}{Z_1 Z_2} E_{i,\text{c.m.}}, \quad (6)$$

where $Z_{1,2,3,4}$ are the atomic numbers of the projectile, target, actinide, and targetlike nucleus, respectively. This expression is used in our work as reference, although it must be taken into account that no energy loss, i.e., $E_{i,\text{c.m.}} = E_{f,\text{c.m.}}$, is predicted when a mass transfer is not accompanied by a charge transfer.

In the present work, the total excitation energy available in the exit channel was derived from Eq. (5), using the energy and the angle of the targetlike nucleus measured by the SPIDER telescope.

Figure 8 shows the reconstructed total excitation-energy distributions for $^{238}\text{U} + ^{12}\text{C}$ inelastic and multinucleon transfer reactions. The elastic peak was fitted to a Gaussian function in order to separate the inelastic component, which rapidly decreases with the excitation energy.

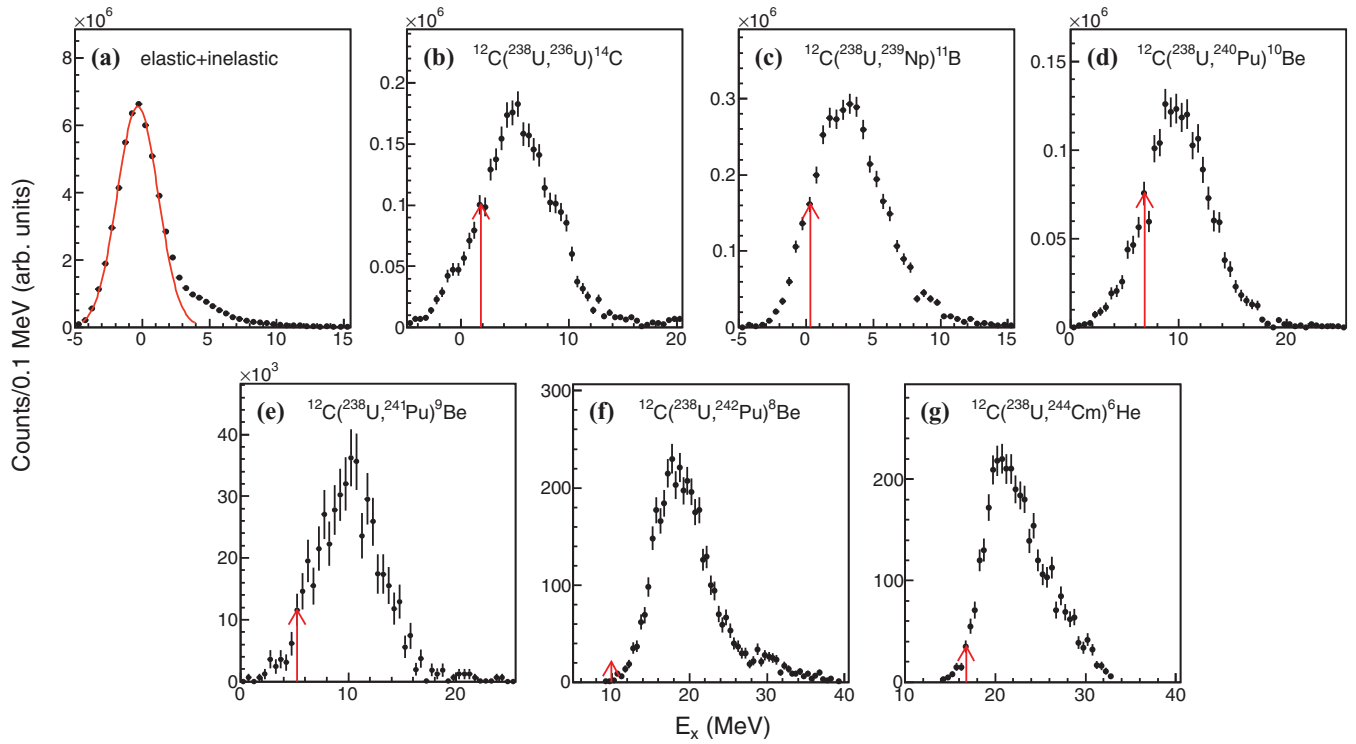


FIG. 8. (Color online) Total excitation-energy distributions measured for $^{238}\text{U} + ^{12}\text{C}$ elastic+inelastic scattering and multinucleon transfer reactions. The elastic peak was fitted to a Gaussian function in order to separate the inelastic component. The most probable excitation energies predicted by Eqs. (5) and (6) are indicated by arrows. The vertical error bars are statistical.

The total excitation-energy distributions populated in transfer reactions show a typical peaked shape, with a full width at half maximum of approximately 8 MeV. The maxima of the experimental distributions are systematically above the values given by Eqs. (6) and (5), reflecting that optimal Q values are smaller than the calculated ones. These results are shown in Table II, which compiles the most probable excitation energies found in this work and the associated optimal Q values. The first were obtained from Gaussian fits around the maxima of the distributions, which also provided the associated errors.

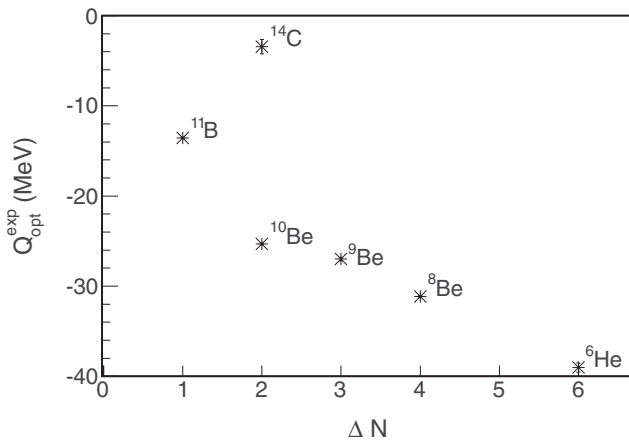


FIG. 9. Optimal Q values determined from the maxima of the total excitation energy distributions. Error bars are smaller than the point size.

Figure 9 shows that, contrary to the predictions of Eq. (6), the amount of energy dissipated increases with the number of transferred nucleons, even if no charge exchange takes place. This result, which is in good agreement with Ref. [38], can be interpreted as an increase of the contact time when more nucleons are transferred. However, the transfer of two neutrons from the beam to the target, leading to ^{14}C and ^{236}U , considerably deviates from the reported trend, showing the influence of the Coulomb interaction in the kinetic energies of the emerging transfer partners, as expressed by Eq. (6).

Finally, the incident neutron energies that, in a neutron capture reaction, would produce a compound nucleus with an excitation energy equal to the reported experimental values are also given in Table II. Both quantities are related to each other through Eq. (2). The previous comparison assumes that the total excitation energy available from inelastic or transfer reactions is exclusively carried by the beamlike partner. Under this assumption, which will be discussed in Sec. V, our data nicely show that transfer reactions cover different excitation-energy regimes. A wide region between equivalent thermal- and fast-neutron induced fission was sampled, the latter being of interest for fast-neutron reactor applications [41], where experimental data are strongly demanded [42].

V. EXCITATION OF LIGHT TRANSFER PARTNERS

Possible excitations of the light inelastic and transfer partners were investigated by means of γ -ray measurements. Figure 10 shows the γ -ray spectra obtained for ^{12}C , ^{11}B , and ^{10}Be , in anticoincidence and coincidence (inset) with the

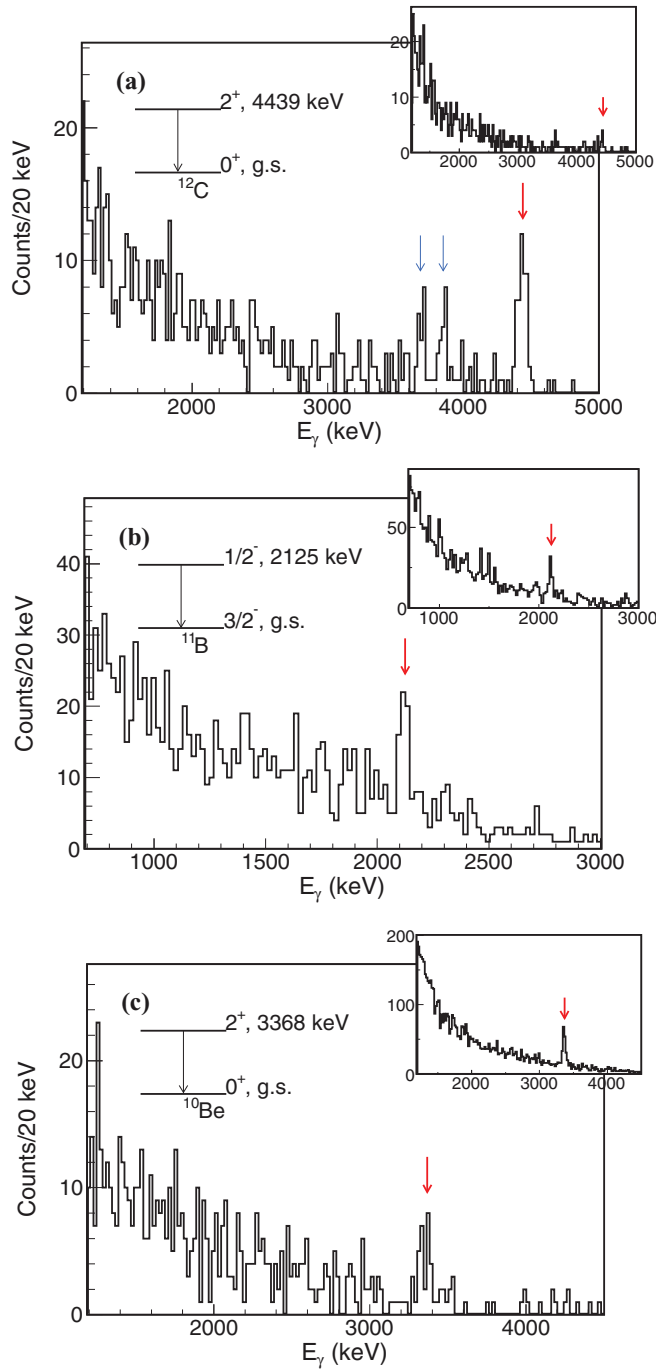


FIG. 10. (Color online) γ -ray spectra measured for ^{12}C (a), ^{11}B (b), and ^{10}Be (c) targetlike nuclei. The diagrams therein represent the deexcitation from the first excited state. The associated γ peaks are marked by thick (red) arrows. Thin (blue) arrows in the spectrum (a) show emissions from excited states of ^{13}C . Main and inset plots correspond to trigger conditions that select anticoincidences and coincidences with the detection of a fission fragment in VAMOS, respectively.

detection of a fission fragment. A scaling factor of 600 should be applied in the first case, as only 1 of 600 anticoincidence events was accepted by the data-acquisition system during the experiment.

Clear γ -ray peaks were observed for ^{12}C , ^{11}B , and ^{10}Be , originating from the deexcitation of their first excited states. Regarding ^{11}B and ^{10}Be , similar results were found for events in anticoincidence and coincidence with the detection of a fission fragment. However, the characteristic γ rays from the first excited state of ^{12}C were hardly visible in fission events. This fact can be explained by the excitation-energy spectra populated in inelastic scattering. As shown in Fig. 8, the 10 MeV of excitation energy required to overcome the fission barrier of ^{238}U and populate at the same time the first excited state of ^{12}C are rarely reached.

No γ -ray emission was observed for ^9Be nuclei, for which the first excited state is placed above the neutron-separation threshold. This result is in good agreement with the high probability of deexcitation through neutron emission reported for this state [43].

A quantitative analysis of the γ -energy spectra provided the probabilities of γ -ray emission from the first excited states of our targetlike nuclei. They were calculated for each reaction channel as follows:

$$P_\gamma = \frac{N_{f+\gamma} + DN_{\bar{f}+\gamma}}{\varepsilon_\gamma(N_f + DN_{\bar{f}})}. \quad (7)$$

The numbers of inelastic/transfer events in coincidence and anticoincidence with the detection of a fission fragment in VAMOS are given by N_f and $N_{\bar{f}}$, respectively, while $N_{f+\gamma}$ and $N_{\bar{f}+\gamma}$ correspond to the numbers of γ rays observed in those events. The factor D is the reduction applied in our acquisition system to $N_{\bar{f}}$ and $N_{\bar{f}+\gamma}$, $D = 600$.

A minimum excitation energy was required for the computation of $N_{\bar{f}}$, which corresponds to the energy of the first excited state of the targetlike nucleus. This threshold was increased by the height of the actinide fission barrier for the computation of N_f , as this is the minimum excitation energy required for a simultaneous fission event. Finally, the parameter ε_γ represents the efficiency of our γ -ray measurements. Its determination was explained in Sec. III.

The obtained probabilities are given in Table III. Similar results were found for ^{12}C , ^{11}B , and ^{10}Be , for which $P_\gamma = 0.12$ – 0.14 . A relative error between 15% and 20% must be considered for this result.

Those events where the targetlike partner is excited require an additional total excitation energy in the exit channel to overcome the actinide fission barrier. A similar effect arises at the onset of the second-chance fission. As a consequence, the fission probability represented as a function of the total excitation energy is expected to differ from the results obtained

TABLE III. Probabilities of γ -ray emissions from the first excited states of targetlike nuclei P_γ^{expt} determined in this work. The γ -ray energies E_γ and the associated detection efficiencies ε_γ are also shown.

Targetlike nucleus	E_γ (keV)	ε_γ	P_γ^{expt}
^{12}C	4439	0.0059 ± 0.0004	0.14 ± 0.03
^{11}B	2125	0.0103 ± 0.0006	0.12 ± 0.02
^{10}Be	3368	0.0073 ± 0.0004	0.14 ± 0.04

in absence of targetlike partner excitation and its average value will be slightly smaller.

VI. FISSION PROBABILITIES

A. Technical considerations

For each transfer channel, the fission probabilities are defined as the proportion of events where the produced compound nuclei decay by fission. In our work, they can be derived as a function of the total excitation energy E_x , using the following expression:

$$P_f(E_x) = \frac{N_f(E_x)}{a[N_f(E_x) + DN_{\bar{f}}(E_x)]}, \quad (8)$$

in which $N_f(E_x)$ and $N_{\bar{f}}(E_x)$ represent the number of events in coincidence and anticoincidence with the detection of a fission fragment in the VAMOS spectrometer. The total number of events for a given excitation energy is given by $N_f(E_x) + DN_{\bar{f}}(E_x)$, where D is the reduction factor applied in our acquisition system to $N_{\bar{f}}$, $D = 600$.

The term a accounts for the acceptance of the VAMOS spectrometer. It was determined as the product of the acceptances in the azimuthal and polar angles of the fission fragments, in the center-of-mass reference frame, denoted by $a_{\phi_{c.m.}}$ and $a_{\theta_{c.m.}}$ in the equation below.

$$a = a_{\phi_{c.m.}} a_{\theta_{c.m.}} \quad (9)$$

The azimuthal and polar angle distributions were obtained by reconstructing the trajectories of the fission fragments through the VAMOS spectrometer [44]. A transformation into the center-of-mass reference frame was applied afterwards, using the kinematics of the transfer reaction.

Four configurations of VAMOS were used, in which the spectrometer was located at 20° and 14° with respect to the beam axis. The measurements at 20° were done with reference magnetic rigidities of 1.1 and 1.2 Tm, while 1.2 and 1.3 Tm were applied at 14° . They allowed to scan polar angles of the fission fragments from 60° to 95° , in the center of mass.

For each VAMOS configuration, the acceptance $a_{\phi_{c.m.}}$ was calculated as the ratio between the FWHM of the reconstructed $\phi_{c.m.}$ distribution and the total 2π range.

The calculation of $a_{\theta_{c.m.}}$ was based on the $\theta_{c.m.}$ distributions measured with the four VAMOS configurations. They were normalized by the acceptance in $\phi_{c.m.}$ and the beam intensity, which was accounted for by the number of elastic events detected in the SPIDER telescope. Then, few points around the maxima of these distributions were fitted to a function of

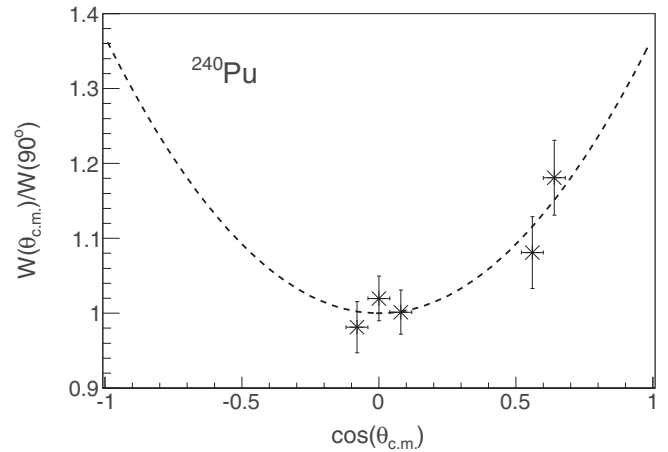


FIG. 11. Anisotropy of the fission fragments in transfer-induced fission of ^{240}Pu . The curve is a fit of the experimental points. It follows Eq. (10).

the form

$$W(\theta_{c.m.}) = W(90^\circ)(1 + \alpha \cos^2 \theta_{c.m.}), \quad (10)$$

providing a first-order description of the anisotropy of the fission fragments [45], $W(0^\circ)/W(90^\circ)$. An example is given in Fig. 11.

This procedure was applied to each transfer-induced fission channel. The resulting anisotropies are given in Table IV. The associated errors account for statistical fluctuations, as well as for movements of the beam of up to 2 mm influencing the rate of elastic events that we use for normalization. The uncertainty introduced by the simplification proposed in Eq. (9), which neglects the interdependence between the magnetic rigidity and the polar and the azimuthal angles [15], and the method given in Fig. 11 was evaluated by means of a simulation and is included as well in the error bars.

For each configuration of VAMOS, the acceptance $a_{\theta_{c.m.}}$ was determined as the ratio between the integral of the $\theta_{c.m.}$ distribution and the integral of the anisotropy curve, in the range $0-180^\circ$. By combining these results with the previously calculated $a_{\phi_{c.m.}}$, total acceptances between 3% and 5% were obtained, depending on the derived anisotropy and the VAMOS configuration. The same error sources affect both the anisotropy and the acceptance calculations, although the latter are less sensitive to them, with a relative error $\delta_a/a = 0.2$. The final fission probabilities were obtained as the average of those measured with the four VAMOS configurations.

TABLE IV. Average fission probabilities obtained in this work, $\langle P_f \rangle^{\text{expt}}$. For reference, the neutron separation energies S_n and the fission barriers [47] are included. Because the fission barriers are double humped, they are given as $B_{f,A}$, and $B_{f,B}$.

Actinide	S_n (MeV)	$B_{f,A}$ (MeV)	$B_{f,B}$ (MeV)	$W(0^\circ)/W(90^\circ)$	E_x range (MeV)	$\langle P_f \rangle^{\text{expt}}$
^{238}U	6.15	5.7	5.7	1.48 ± 0.39	6–12	0.22 ± 0.06
^{239}Np	6.22	6.1	5.6	1.66 ± 0.46	6–12	0.38 ± 0.11
^{240}Pu	6.53	5.6	5.1	1.37 ± 0.36	5–18	0.72 ± 0.23
^{241}Pu	5.24	6.1	5.5	2.19 ± 0.59	6–18	0.80 ± 0.24
^{242}Pu	6.31	5.6	5.1	1.00 ± 0.28	14–28	0.86 ± 0.33
^{244}Cm	6.80	5.8	4.3	1.43 ± 0.43	16–32	1.0 ± 0.4

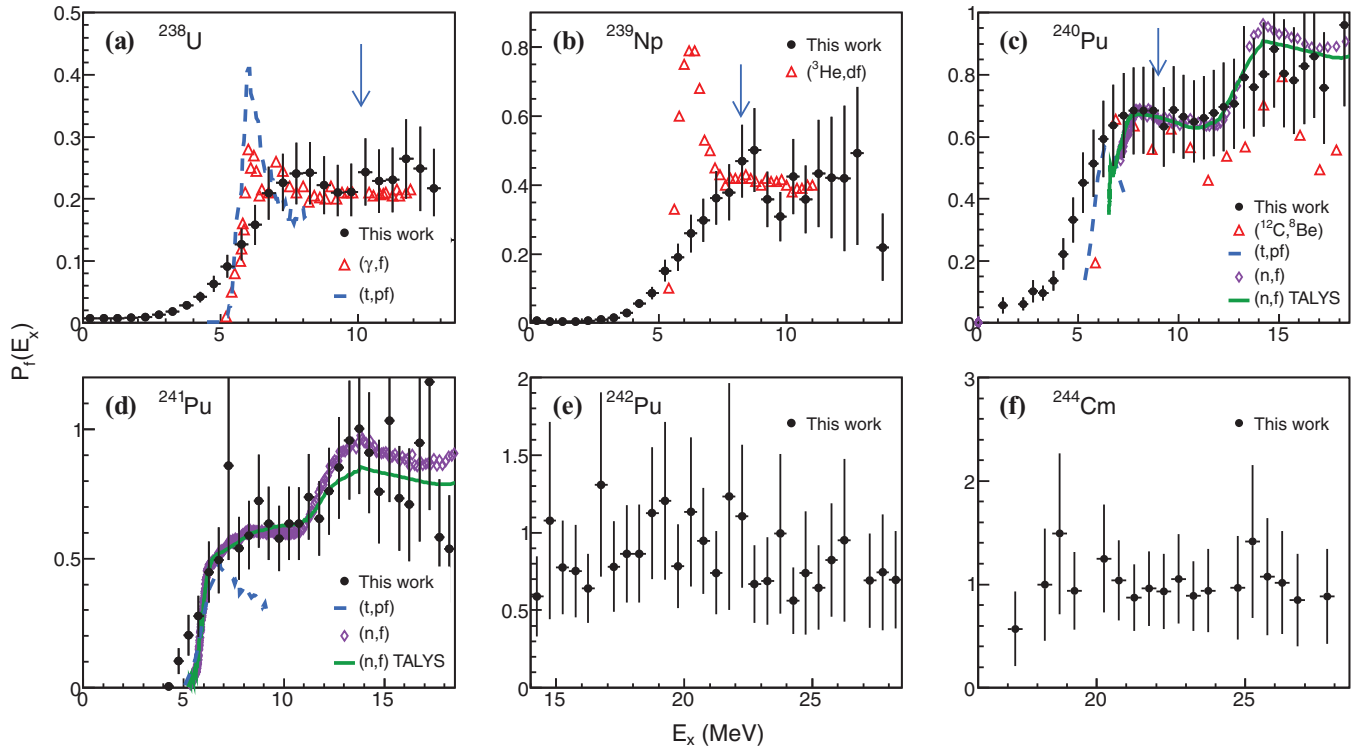


FIG. 12. (Color online) Fission probabilities as a function of the total excitation energy. Results are presented for $^{238}\text{U} + ^{12}\text{C}$ inelastic scattering (a) and transfer-induced fission reactions (b)–(f). The fissioning nucleus is indicated in each figure. Earlier γ -, transfer-, and neutron-induced fission data are included for comparison, as well as TALYS calculations of neutron-induced fission probabilities. The corresponding references are given in the text. Arrows point to excitation energies equal to the sum of the fission barrier and the energy of the first excited state of the light transfer partner.

The influence of targetlike nuclei scattered off the Al collimator on the excitation-energy distributions, which is discussed in Sec. III and Fig. 3, was observed in the ^{238}U fission probabilities. Even though the particles scattered off the collimator were strongly suppressed by limiting the azimuthal and polar angles of targetlike nuclei, a small fraction of elastic events wrongly reconstructed lead to a significant overestimation of $N_f(E_x)$, as the elastic cross section dominates by an order of magnitude over the inelastic one. The quantity $N_f(E_x)$ remained comparably unaffected, because these events were always in coincidence with the detection of a fission fragment. As a result, the fission probability was underestimated by approximately a factor of 2. A correction factor was thus applied to the ^{238}U fission probability by scaling it to the distributions available from earlier works [2,46], using the points at 9 MeV of excitation energy as reference. However, the effect is still visible in the drop of the fission probability at excitation energies above 12 MeV.

A similar issue affected the ^{239}Np fission probabilities, due to the wrong reconstruction of transfer events with excitation energies below the fission barrier. The scattering of these particles off the collimator lead to an overestimation of $N_f(E_x)$, although with a more moderated effect because, in this case, the cross section below the fission barrier is considerably smaller than for ^{238}U . The restrictions applied in the azimuthal and polar angles of targetlike nuclei provided ^{239}Np fission probabilities in good agreement with previous data [1], without need for a correction factor.

B. Results

Figure 12 shows the fission probabilities obtained in this work, as a function of the total excitation energy available in the exit channel. Six fissioning systems, between U and Cm, were investigated. The average values along the given excitation-energy ranges are compiled in Table IV.

The rising of ^{238}U , ^{239}Np , and $^{240,241}\text{Pu}$ fission probabilities at excitation energies of approximately 6 MeV shows the onset of first-chance fission, reflecting the heights of the fission barriers. The general agreement with earlier works and the recommended values of fission barriers [47], which are summarized in Table IV, brings confidence in the present technique, given the moderate excitation-energy resolution achievable in inverse-kinematics measurements.

Second-chance fission, for which a neutron is emitted before fission occurs, takes place at higher excitation energies. For reference, neutron separation energies are also reported in Table IV. This region could not be explored for ^{238}U and ^{239}Np because the minimum excitation energies required were barely reached, as shown in Fig. 8. However, the onset of second-chance fission is clearly visible for $^{240,241}\text{Pu}$, for which the fission probabilities suddenly increase when excitation energies of approximately 12 MeV are reached.

The fission of ^{242}Pu and ^{244}Cm was investigated at higher excitation energies. For the latter, a minor actinide of interest in accelerator-driven systems for the recycling of radioactive waste [17], the region above 15 MeV has been addressed for the first time [18].

The comparison of our results with other works provides some hints concerning the influence of the reaction mechanism used to produce the fissioning system. One of these indications may be the disappearance of the peak near the onset of first chance fission for ^{238}U and ^{239}Np , which diverges from earlier (γ, f) [46], (t, pf) [2] and $(^3\text{He}, df)$ [1] measurements. The angular momentum transferred into the fissioning system in heavy-ion collisions may be inferred to explicate this difference.

Earlier transfer-induced fission measurements for $^{240,241}\text{Pu}$ are limited to a single $(^{12}\text{C}, ^8\text{Be})^{240}\text{Pu}$ experiment [10] and (t, pf) reactions [2], where only the region around the fission barrier was addressed. For both $^{240,241}\text{Pu}$ isotopes, neutron-induced fission probabilities were calculated with the TALYS code [48], as well the cross sections of the formation of a compound nucleus in $^{239}\text{Pu}+n$ and $^{240}\text{Pu}+n$ reactions [49]. They were corrected with a pre-equilibrium component, which becomes important at excitation energies above 10 MeV. Following Eq. (1), the compound nucleus cross sections were used to convert the experimental neutron-induced fission cross sections provided in Ref. [50] into fission probabilities. Figures 12(c) and 12(d) show the comparisons with our results.

Because the neutron separation energy of ^{240}Pu is approximately 1 MeV above the fission barrier, excitation energies around the barrier cannot be addressed in neutron capture reactions. The apparent threshold shown by $^{239}\text{Pu}(n, f)$ data is actually related to the neutron-separation energy and reflects the importance of the so-called compound-nucleus elastic channel, where the compound nucleus decays to its ground state by emitting a neutron. A discrepancy of up to a factor of 2 is observed in this region with respect to our results. It may be explained by the higher angular momenta populated in our work, which forbids the compound-nucleus elastic branch.

The fission probabilities measured in $(^{12}\text{C}, ^8\text{Be})^{240}\text{Pu}$ [10] were obtained in a transfer reaction similar to the one that we use. However, these data present lower fission probabilities than measured in the present work. Surprisingly, they are also lower than the ones obtained in neutron-induced fission, whereas more angular momentum is supposed to be inferred in the fissioning system, increasing the fission probability. This fact may indicate an underestimation of the results given in Ref. [10].

In the case of ^{241}Pu , the fission barrier is above the neutron-separation energy, where the compound-nucleus elastic cross section becomes less important. As a consequence, a better agreement with neutron-induced fission data is observed, both by our data and earlier (t, pf) measurements [2].

A careful observation of Fig. 12 also provides some indications of the role played by excitations of the light transfer partner in the fission probabilities. Although tiny, our results for ^{238}U , ^{239}Np , and ^{240}Pu show structures at excitation energies corresponding to the sum of the height of the fission barrier and the energy of the first excited state of the light transfer partner, indicated by arrows in Fig. 12. They are presumably produced by this effect.

The sensitivity of our results is limited by the background produced in the collimator behind the target, the excitation-

energy resolution, and the uncertainty in the determination of the acceptance of the VAMOS spectrometer, making it difficult to draw firmer conclusions on the interpretation of the measured fission probabilities and the differences with respect to the measurements in lighter transfer reactions. Further experiments of this type, as well as comparisons of the fission probabilities measured through different techniques, would be important in order to deepen into this discussion, which is crucial for the application of the surrogate-reaction technique.

VII. SUMMARY AND CONCLUSIONS

The present work approaches transfer-induced fission measurements from an innovative perspective, which implies the use of heavier transfer reactions and inverse kinematics. Inelastic and multinucleon transfer reactions between a ^{238}U beam and a ^{12}C target, at energies of approximately 10% above the Coulomb barrier, have been considered for the investigation of fission properties of exotic actinides, hardly accessible with standard neutron-irradiation measurements. As a result, up to nine different transfer channels were identified, leading to a variety of neutron-rich, excited actinides, from U to Cm.

The characterization of $^{238}\text{U} + ^{12}\text{C}$ inelastic and transfer channels was obtained from the detection and the identification of the final targetlike partner in a Si telescope. Nucleons were mainly transferred from the light target to the heavy beam, while the flux in the opposite direction was limited to one and two neutrons. Cross sections between few and few tens of mb were measured for the populated channels. The inelastic and transfer differential cross sections were found to be bell shaped and peaked above the calculated grazing angle.

In addition, total excitation-energy distributions were determined. The use of inverse kinematics, which makes possible the isotopic identification of the heavy and light fission fragments [15], limited our excitation-energy resolution to 2.7 MeV (FWHM).

Our results show that higher excitation energies are reached as the number of nucleons transferred increases, even if no charge exchange takes places. The investigated reactions sample different excitation-energy regimes up to 30 MeV. In equivalent neutron-induced fission experiments, they would cover the region between thermal and fast neutron-induced fission.

Furthermore, possible excitations of the targetlike nuclei in the exit channel were explored for the first time. The decays from the first excited states of ^{12}C , ^{11}B , and ^{10}Be were observed with probabilities of 0.12–0.14, by means of in-flight γ -ray measurements, performed in the target region. Beyond its importance for the characterization of transfer reactions, this result shows that the excitation of the light transfer partner has to be taken into account for an accurate description of the fission probabilities within the surrogate-reaction technique.

Fission probabilities were determined for ^{238}U , ^{239}Np , and $^{240,241}\text{Pu}$, as well as for ^{242}Pu and ^{244}Cm , where excitation energies above 15 MeV were addressed. The distributions obtained for the first reflect the height of the fission barrier,

the role of the angular momentum populated in the fissioning system, and the effect of the reported excitation of targetlike nuclei, although the limited sensitivity of the results did not allow a deeper investigation of these features. Further investigations of this subject are advisable and will be of strong interest for the framework of the surrogate-reaction technique. The reduction of the uncertainty in the fission probabilities would require a dedicated experimental setup. In the present work, the use of the VAMOS spectrometer for the identification of the fission fragments [15] results in a small solid angle ($\sim 4\%$) for the fission detection. A less sensitive, but more efficient, setup would allow for a detailed study of the fission-fragment anisotropy, with a complete coverage of the fission-fragment angular distribution. The latter will benefit from the kinematical boost of the fragments in inverse kinematics, which makes possible a 4π coverage in the center-of-mass reference frame with a relatively small setup.

This work shows the potential of transfer-induced fission experiments with heavier targets, such as ^{12}C , to access more neutron-rich and short-lived actinides, heavier than the beam. Future inverse-kinematics experiments based on this technique will benefit from the radioactive beams provided by ISOL facilities [51,52], expanding the body of fissioning systems available to the experimental research.

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