

Research Article

Eigenvalue Asymptotics of the Even-Dimensional Exterior Landau-Neumann Hamiltonian

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We study the Schrödinger operator with a constant magnetic field in the exterior of a compact domain in \mathbb{R}^{2d} , $d \geq 1$. The spectrum of this operator consists of clusters of eigenvalues around the Landau levels. We give asymptotic formulas for the rate of accumulation of eigenvalues in these clusters. When the compact is a Reinhardt domain we are able to show a more precise asymptotic formula.

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1. Introduction

The Landau Hamiltonian describes a charged particle moving in a plane, influenced by a constant magnetic field of strength $B > 0$ orthogonal to the plane. It is a classical result, see [1, 2], that the spectrum of the Landau Hamiltonian consists of infinitely degenerate eigenvalues $B(2q + 1)$, $q = 0, 1, 2, \dots$, called Landau levels.

In this paper, we will study the even-dimensional Landau Hamiltonian outside a compact obstacle, imposing magnetic Neumann conditions at the boundary. Our motivation to study this operator comes mainly from the papers [3, 4]. Spectral properties of the exterior Landau Hamiltonian in the plane are discussed in [3], under both Dirichlet and Neumann conditions at the boundary, with focus mainly on properties of the eigenfunctions. A more qualitative study of the spectrum is done in [4], where the authors fix an interval around a Landau level and describe how fast the eigenvalues in that cluster converge to that Landau level. They work in the plane and with Dirichlet boundary conditions only. The goal of this paper is to perform the same qualitative description when we impose magnetic Neumann conditions at the boundary. Moreover, we do not limit ourself to the plane but work in arbitrary even-dimensional Euclidean space.

The result is that the eigenvalues do accumulate with the same rate to the Landau levels for both types of boundary conditions; see Theorem 3.2 for the details. However, the eigenvalues can only accumulate to a Landau level from below in the Neumann setting. In the Dirichlet case they accumulate only from above.

It should be mentioned that we suppose that the compact set removed has no holes and that its boundary is smooth. This is far more restrictive than the conditions imposed on the compact set in [4].

Several different perturbations of the Landau Hamiltonian have been studied in the last years; see [4–11]. They all share the common idea of making a reduction to a certain Toeplitz-type operator whose spectral asymptotics are known. We also do this kind of reduction. The method we use is based on the theory for pseudodifferential operators and boundary PDE methods, which we have not seen in any of the mentioned papers.

In Section 2, we define the Landau Hamiltonian and give some auxiliary results about its spectrum, eigenspaces, and Green function.

We begin Section 3 by defining the exterior Landau Hamiltonian with magnetic Neumann boundary condition and formulating and proving the main theorems (Theorems 3.1 and 3.2) about the spectral asymptotics of the operator. The main part of the proof, the reduction step, is quite technical and therefore moved to Section 4. When the reduction step is done we use the asymptotic formulas of the spectrum of the Toeplitz-type operators, given in [8, 10], to obtain the asymptotic formulas in Theorem 3.2.

In the higher dimensional case (\mathbb{R}^{2d} , $d > 1$), we also consider the case when the compact obstacle is a Reinhardt domain. We use some ideas from [12] to prove a more precise asymptotic formula for the eigenvalues. This is done in Section 5.

2. The Landau Hamiltonian in \mathbb{R}^{2d}

We denote by $x = (x^1, \dots, x^{2d})$ a point in \mathbb{R}^{2d} . Let $B > 0$ and denote by \vec{a} the magnetic vector potential

$$\vec{a}(x) = (a_1(x), \dots, a_{2d}(x)) = \frac{B}{2}(-x^2, x^1, -x^4, x^3, \dots, -x^{2d}, x^{2d-1}). \quad (2.1)$$

It corresponds to an isotropic magnetic field of constant strength B . The Landau Hamiltonian L in \mathbb{R}^{2d} describes a charged, spinless particle in this homogeneous magnetic field. It is given by

$$L = (-i\nabla - \vec{a})^2 \quad (2.2)$$

and is essentially self-adjoint on the set $C_0^\infty(\mathbb{R}^{2d})$ in the usual Hilbert space $\mathcal{H} = L_2(\mathbb{R}^{2d})$. For $j = 1, \dots, d$, we also introduce the self-adjoint operators

$$L_j = \left(-i \left(\frac{\partial}{\partial x^{2j-1}}, \frac{\partial}{\partial x^{2j}} \right) - (a_{2j-1}, a_{2j}) \right)^2, \quad (2.3)$$

in the Hilbert spaces $\mathcal{H}_j = L_2(\mathbb{R}^2)$. Note that $\mathcal{H} = \otimes_{j=1}^d \mathcal{H}_j$, and

$$L = L_1 \otimes I^{\otimes(d-1)} + I \otimes L_2 \otimes I^{\otimes(d-2)} + \dots + I^{\otimes(d-1)} \otimes L_d. \quad (2.4)$$

2.1. Landau Levels

The spectrum of each two-dimensional Landau Hamiltonian L_j consists of the so-called Landau levels, eigenvalues $B(2q + 1)$, $q \in \mathbb{N} := \{0, 1, 2, \dots\}$, each of infinite multiplicity. Let $\hat{\kappa} = (\kappa_1, \dots, \kappa_d) \in \mathbb{N}^d$ be a multi-index. We denote by $|\hat{\kappa}| = \kappa_1 + \dots + \kappa_d$ the length of the multi-index $\hat{\kappa}$ and also set $\hat{\kappa}! = \kappa_1! \dots \kappa_d!$. From (2.4) it follows that the spectrum of L consists of the infinitely degenerate eigenvalues

$$\Lambda_{\hat{\kappa}} = B \sum_{j=1}^d (2\kappa_j + 1), \quad \kappa_j \in \mathbb{N}. \quad (2.5)$$

Note that $\Lambda_{\hat{\kappa}} = \Lambda_{\hat{\kappa}'}$ if $|\hat{\kappa}| = |\hat{\kappa}'|$. Hence the spectrum of L consists of eigenvalues of the form $\Lambda_{\mu} = B(2\mu + d)$, $\mu \in \mathbb{N}$.

2.2. Creation and Annihilation Operators

The structure of the eigenspaces of L has been described before in [8]. We give the results without proofs. It is convenient to introduce complex notation. Let $z = (z^1, \dots, z^d) \in \mathbb{C}^d$, where $z^j = x^{2j-1} + ix^{2j}$. Also, we use the scalar potential $W(z) = -(B/4)|z|^2$ and the complex derivatives

$$\frac{\partial}{\partial z^j} = \frac{1}{2} \left(\frac{\partial}{\partial x^{2j-1}} - i \frac{\partial}{\partial x^{2j}} \right), \quad \frac{\partial}{\partial \bar{z}^j} = \frac{1}{2} \left(\frac{\partial}{\partial x^{2j-1}} + i \frac{\partial}{\partial x^{2j}} \right). \quad (2.6)$$

We define creation and annihilation operators Q_j^* , Q_j as

$$Q_j^* = -2ie^{-W} \frac{\partial}{\partial z^j} e^W, \quad Q_j = -2ie^W \frac{\partial}{\partial \bar{z}^j} e^{-W}, \quad (2.7)$$

and note that

$$[Q_j^*, Q_k^*] = [Q_j, Q_k] = [Q_j^*, Q_k] = 0, \quad \text{if } j \neq k. \quad (2.8)$$

The notation Q_j^* for the creation operators is motivated by the fact that it is the formal adjoint of Q_j in \mathcal{H} .

A function u belongs to the lowest Landau level Λ_0 if and only if $Q_j u = 0$ for $j = 1, \dots, d$. This means that the function $f = e^{-W} u$ is an entire function, so via multiplication by e^{-W} the eigenspace \mathcal{L}_{Λ_0} corresponding to Λ_0 is equivalent to the Fock space

$$\mathfrak{F}_B^2 = \left\{ f \mid f \text{ is entire and } \int_{\mathbb{C}^d} |f|^2 e^{-(B/2)|z|^2} dm(z) < \infty \right\}. \quad (2.9)$$

Here, and elsewhere, dm denotes the Lebesgue measure. A function u belongs to the eigenspace $\mathcal{L}_{\Lambda_\mu}$ of the Landau level Λ_μ if and only if it can be written in the form

$$u = \sum_{|\bar{k}|=\mu} c_{\bar{k}} (Q^*)^{\bar{k}} (e^W f_{\bar{k}}), \quad (2.10)$$

where $(Q^*)^{\bar{k}} = (Q_1^*)^{k_1} \dots (Q_d^*)^{k_d}$ and $f_{\bar{k}}$ all belong to \mathfrak{F}_B^2 . The multiplicity of each eigenvalue Λ_μ is infinite. We denote by $\rho_{\Lambda_{\bar{k}}}$ and ρ_{Λ_μ} the projection onto the eigenspaces $\mathcal{L}_{\Lambda_{\bar{k}}}$ and $\mathcal{L}_{\Lambda_\mu}$, respectively, and note by (2.16) that the orthogonal decompositions

$$\mathcal{L}_{\Lambda_\mu} = \bigoplus_{|\bar{k}|=\mu} \mathcal{L}_{\Lambda_{\bar{k}}}, \quad \rho_{\Lambda_\mu} = \bigoplus_{|\bar{k}|=\mu} \rho_{\Lambda_{\bar{k}}} \quad (2.11)$$

hold in \mathcal{L} .

2.3. The Resolvent

Let $R_\rho = (L + \rho I)^{-1}$ be the resolvent of L , $\rho \geq 0$. An explicit formula of the kernel $G_\rho(x, y)$ of R_ρ was given in [3] for $d = 1$. In Section 4.2, we will use the behavior of $G_\rho(x, y)$ near the diagonal $x = y$, given in the following lemma.

Lemma 2.1. *R_ρ is an integral operator with kernel $G_\rho(x, y)$ that has the following singularity at the diagonal:*

$$G_\rho(x, y) \sim \begin{cases} \frac{1}{\pi} \log \left(\frac{1}{|x-y|} \right) + O(1), & d = 1, \\ \frac{1}{2\pi^2} |x-y|^{-2} + O \left(\log \left(\frac{1}{|x-y|} \right) \right), & d = 2, \\ \frac{\Gamma(d-1)}{2\pi^d} |x-y|^{2-2d} + O(|x-y|^{4-2d}), & d > 2, \end{cases} \quad (2.12)$$

as $|x-y| \rightarrow 0$.

Proof. The kernel $G_\rho(x, y)$ of R_ρ can be written as

$$G_\rho(x, y) = \int_0^\infty e^{-\rho t} e^{-Lt}(x, y) dm(t). \quad (2.13)$$

Now, since the variables separate pairwise, we have

$$e^{-Lt}(x, y) = \prod_{j=1}^d e^{-L_j t}(x^{2j-1}, x^{2j}, y^{2j-1}, y^{2j}). \quad (2.14)$$

The formula for $e^{-L_j t}$ is given in [13]. It reads

$$\begin{aligned} e^{-L_j t} &= \frac{B}{4\pi} \exp\left(-\frac{iB}{2}(x^{2j-1}y^{2j} - x^{2j}y^{2j-1})\right) \frac{1}{\sinh(Bt/2)} \\ &\times \exp\left(-\frac{B}{4} \coth\left(\frac{Bt}{2}\right) ((x^{2j-1} - y^{2j-1})^2 + (x^{2j} - y^{2j})^2)\right). \end{aligned} \quad (2.15)$$

Hence the formula for $G_\rho(x, y)$ becomes

$$G_\rho(x, y) = \left(\frac{B}{4\pi}\right)^d \exp\left(-\frac{iB}{2} \sum_{j=1}^d (x^{2j-1}y^{2j} - x^{2j}y^{2j-1})\right) I(|x - y|^2), \quad (2.16)$$

where

$$I(s) = \int_0^\infty e^{-\rho t} \frac{1}{\sinh^d(Bt/2)} \exp\left(-\frac{B}{4} \coth\left(\frac{Bt}{2}\right) s\right) dm(t). \quad (2.17)$$

An expansion of $I(s)$ shows that

$$I(s) \sim \begin{cases} \left(\frac{2}{B}\right) \log\left(\frac{1}{s}\right) + O(1), & d = 1, \\ \frac{8}{B^2} s^{-1} + O\left(\log\left(\frac{1}{s}\right)\right), & d = 2, \text{ as } s \rightarrow 0, \\ \left(\frac{4}{B}\right)^d \frac{\Gamma(d-1)}{2} s^{1-d} + O(s^{2-d}), & d > 2, \end{cases} \quad (2.18)$$

from which (2.12) follows. □

3. The Exterior Landau-Neumann Hamiltonian in \mathbb{R}^{2d}

Let $K \subset \mathbb{R}^{2d}$ be a simply connected compact domain with smooth boundary Γ and let $\Omega = \mathbb{R}^{2d} \setminus K$. We define the exterior Landau-Neumann Hamiltonian L_Ω in $\mathcal{H}_\Omega = L_2(\Omega)$ by

$$L_\Omega = (-i\nabla - \vec{a})^2, \quad \text{in } \Omega, \quad (3.1)$$

with magnetic Neumann boundary conditions

$$\partial_N u := (-i\nabla - \vec{a})u \cdot \nu = 0, \quad \text{on } \Gamma. \quad (3.2)$$

Here ν denotes the exterior normal to Γ . Our aim is to study how much the spectrum of L_Ω differs from the Landau levels discussed in the previous section. The first theorem below states that the eigenvalues of L_Ω can accumulate to each Landau level only from below. The second theorem says that the eigenvalues do accumulate to the Landau levels from below, and the rate of convergence is given.

Theorem 3.1. *For every $\mu \in \mathbb{N}$ and each ε , $0 < \varepsilon < dB$, the number of eigenvalues of L_Ω in the interval $(\Lambda_\mu, \Lambda_\mu + \varepsilon)$ is finite.*

Denote by $l_1^{(\mu)} \leq l_2^{(\mu)} \leq \dots$ the eigenvalues of L_Ω in the interval $(\Lambda_{\mu-1}, \Lambda_\mu)$ and by $N(a, b, T)$ the number of eigenvalues of the operator T in the interval (a, b) , counting multiplicities. Also, let $\text{Cap}(K)$ denote the logarithmic capacity of K ; see [14, Chapter 2].

Theorem 3.2. *Let $\mu \in \mathbb{N}$.*

- (a) *If $d = 1$ then $\lim_{j \rightarrow \infty} (j!(\Lambda_\mu - l_j^{(\mu)}))^{1/j} = (B/2)(\text{Cap}(K))^2$.*
- (b) *If $d > 1$ then $N(\Lambda_{\mu-1}, \Lambda_\mu - \lambda, L_\Omega) \sim (\frac{\mu+d-1}{d-1})(1/d!)(|\ln \lambda| / \ln |\ln \lambda|)^d$ as $\lambda \searrow 0$.*

3.1. Proof of the Theorems

We want to compare the spectrum of the operators L and L_Ω . However, the expression $L - L_\Omega$ has no meaning since L and L_Ω act in different Hilbert spaces. We introduce the Hilbert space $\mathcal{H}_K = L_2(K)$ and define the interior Landau-Neumann Hamiltonian L_K in \mathcal{H}_K by the same formulas as in (3.1) and (3.2) but with Ω replaced by K . We note that $\mathcal{H} = \mathcal{H}_K \oplus \mathcal{H}_\Omega$ and define \tilde{L} as

$$\tilde{L} = L_K \oplus L_\Omega, \quad \text{in } \mathcal{H}_K \oplus \mathcal{H}_\Omega. \quad (3.3)$$

The inverse of L_K is compact, so L_K has at most a finite number of eigenvalues in each interval $(\Lambda_{\mu-1}, \Lambda_\mu)$. The operators L_K and L_Ω act in orthogonal subspaces of \mathcal{H} , so $\sigma(\tilde{L}) = \sigma(L_K) \cup \sigma(L_\Omega)$. This means that \tilde{L} has the same spectral asymptotics as L_Ω in each interval $(\Lambda_{\mu-1}, \Lambda_\mu)$, so it is enough to prove the statements in Theorems 3.1 and 3.2 for the operator \tilde{L} instead of L_Ω .

Since the unbounded operators L and \tilde{L} have different domains, we cannot compare them directly. However, they act in the same Hilbert space, so we can compare their inverses. Let

$$R = R_0 = L^{-1}, \quad \tilde{R} = \tilde{L}^{-1} = L_K^{-1} \oplus L_\Omega^{-1}, \quad (3.4)$$

and set

$$V = \tilde{R} - R, \quad T_\mu = \rho_{\Lambda_\mu} V \rho_{\Lambda_\mu}, \quad \text{for } \mu \in \mathbb{N}. \quad (3.5)$$

Lemma 3.3. *V is nonnegative and compact.*

Proof. See Section 4.1. □

By Weyl's theorem, the essential spectrum of R and \tilde{R} coincides. Since $\tilde{R} = R + V$ and $V \geq 0$, Theorem 3.1 follows immediately from [15, Theorem 9.4.7] and the fact that $\sigma(R) = \sigma_{\text{ess}}(R) = \{\Lambda_\mu^{-1}\}$. We continue with the proof of Theorem 3.2.

Let $\tau > 0$ be such that $((\Lambda_\mu^{-1} - 2\tau, \Lambda_\mu^{-1} + 2\tau) \setminus \{\Lambda_\mu^{-1}\}) \cap \sigma_{\text{ess}}(R) = \emptyset$. Denote the eigenvalues of T_μ by

$$t_1^{(\mu)} \geq t_2^{(\mu)} \geq \dots, \quad (3.6)$$

and the eigenvalues of \tilde{R} in the interval $(\Lambda_\mu^{-1}, \Lambda_\mu^{-1} + \tau)$ by

$$r_1^{(\mu)} \geq r_2^{(\mu)} \geq \dots. \quad (3.7)$$

Lemma 3.4. *Given $\varepsilon > 0$, there exists an integer l such that*

$$(1 - \varepsilon)t_{j+l}^{(\mu)} \leq r_j^{(\mu)} - \Lambda_\mu^{-1} \leq (1 + \varepsilon)t_{j-l}^{(\mu)}, \quad \text{for all sufficiently large } j. \quad (3.8)$$

Proof. See [4, Proposition 2.2]. □

Hence the study of the asymptotics of the eigenvalues of \tilde{R} is reduced to the study of the eigenvalues of the Toeplitz-type operator T_μ . For a bounded simply connected set U in \mathbb{R}^{2d} , we define the Toeplitz operator S_μ^U as

$$S_\mu^U = \mathcal{D}_{\Lambda_\mu} \chi_U \mathcal{D}_{\Lambda_\mu}, \quad (3.9)$$

where χ_U denotes the characteristic function of U . The following lemma reduces our problem to the study of these Toeplitz operators, which are easier to study than T_μ .

Lemma 3.5. *Let $K_0 \Subset K \Subset K_1$ be compact domains such that $\partial K_i \cap \Gamma = \emptyset$. There exist a constant $C > 0$ and a subspace $\mathcal{S} \subset \mathcal{L}$ of finite codimension such that*

$$\frac{1}{C} \langle f, S_\mu^{K_0} f \rangle \leq \langle f, T_\mu f \rangle \leq C \langle f, S_\mu^{K_1} f \rangle \quad (3.10)$$

for all $f \in \mathcal{S}$.

Proof. See Section 4.2. □

The asymptotic expansion of the spectrum of S_μ^U is given in the following lemma.

Lemma 3.6. Denote by $s_1^{(\mu)} \geq s_2^{(\mu)} \geq \dots$ the eigenvalues of S_μ^U and by $n(\lambda, S_\mu^U)$ the number of eigenvalues of S_μ^U greater than λ (counting multiplicity). Then

$$(a) \text{ if } d = 1 \text{ we have } \lim_{j \rightarrow \infty} (j! s_j^{(\mu)})^{1/j} = (B/2)(\text{Cap}(U))^2,$$

$$(b) \text{ if } d > 1 \text{ we have } n(\lambda, S_\mu^U) \sim \binom{\mu+d-1}{d-1} (1/d!) (|\ln \lambda| / \ln |\ln \lambda|)^d \text{ as } \lambda \searrow 0.$$

Proof. See [10, Lemma 3.2] for part (a) and [8, Proposition 7.1] for part (b). \square

Proof of Theorem 3.2. We are now able to finish the proof of Theorem 3.2. By letting K_0 and K_1 in Lemma 3.5 get closer and closer to our compact K we see that the eigenvalues $\{t_j^{(\mu)}\}$ of T_μ satisfy

$$\lim_{j \rightarrow \infty} (j! t_j^{(\mu)})^{1/j} = \frac{B}{2} (\text{Cap}(K))^2 \quad (3.11)$$

if $d = 1$, and

$$n(\lambda, T_\mu) \sim \binom{\mu+d-1}{d-1} \frac{1}{d!} \left(\frac{|\ln \lambda|}{\ln |\ln \lambda|} \right)^d, \text{ as } \lambda \searrow 0 \quad (3.12)$$

if $d > 1$. Since neither formula (3.11) nor (3.12) is sensitive for finite shifts in the indices, it follows from Lemma 3.4 that the eigenvalues of $\{r_j^{(\mu)}\} \tilde{R}$ satisfy

$$\lim_{j \rightarrow \infty} (j! (r_j^{(\mu)} - \Lambda_\mu^{-1}))^{1/j} = \frac{B}{2} (\text{Cap}(K))^2 \quad (3.13)$$

if $d = 1$, and

$$N(\Lambda_\mu^{-1} + \lambda, \Lambda_{\mu-1}^{-1}, \tilde{R}) \sim \binom{\mu+d-1}{d-1} \frac{1}{d!} \left(\frac{|\ln \lambda|}{\ln |\ln \lambda|} \right)^d, \text{ as } \lambda \searrow 0. \quad (3.14)$$

If we translate this in terms of \tilde{L} we get

$$\lim_{j \rightarrow \infty} (j! (\Lambda_\mu - l_j^{(\mu)}))^{1/j} = \frac{B}{2} (\text{Cap}(K))^2 \quad (3.15)$$

for $d = 1$, and

$$\begin{aligned} N(\Lambda_{\mu-1}, \Lambda_\mu - \lambda, \tilde{L}) &\sim \binom{\mu+d-1}{d-1} \frac{1}{d!} \left(\frac{|\ln(\lambda/\Lambda_\mu(\Lambda_\mu - \lambda))|}{\ln |\ln(\lambda/\Lambda_\mu(\Lambda_\mu - \lambda))|} \right)^d \\ &\sim \binom{\mu+d-1}{d-1} \frac{1}{d!} \left(\frac{|\ln \lambda|}{\ln |\ln \lambda|} \right), \text{ as } \lambda \searrow 0, \end{aligned} \quad (3.16)$$

for $d > 1$. This completes the proof of Theorem 3.2. \square

4. Proof of the Lemmas

In this section, we prove Lemmas 3.3 and 3.5. We will use the theory of pseudodifferential operators and boundary layer potentials. More details about these tools can be found in [16] and [17, Chapter 5].

4.1. Proof of Lemma 3.3

The operators L and \tilde{L} are defined by the same expression, but the domain of \tilde{L} is contained in the domain of L . It follows from [4, Proposition 2.1] that $L - \tilde{L} \geq 0$, and hence $V = \tilde{R} - R \geq 0$.

Next we prove the compactness of V . Let f and g belong to \mathcal{A} . Also, let $u = Rf$ and $v = \tilde{R}g$. Then u belongs to the domain of L and v belongs to the domain of \tilde{L} , so $v = v_K \oplus v_\Omega$, and $L_K v_K \oplus L_\Omega v_\Omega = g$. Integrating by parts and using (3.2) for v_K and v_Ω , we get

$$\begin{aligned} \langle f, Vg \rangle &= \langle f, \tilde{R}g \rangle - \langle Rf, g \rangle \\ &= \int_K Lu \cdot \overline{v_K} dm(x) + \int_\Omega Lu \cdot \overline{v_\Omega} dm(x) - \int_K u \cdot \overline{L_K v_K} dm(x) - \int_\Omega u \cdot \overline{L_\Omega v_\Omega} dm(x) \\ &= \int_\Gamma \partial_N u \cdot \overline{(v_\Omega - v_K)} dS. \end{aligned} \tag{4.1}$$

Here dS denotes the surface measure on Γ .

Take a smooth cutoff function $\chi \in C_0^\infty(\mathbb{R}^{2d})$ such that $\chi(x) = 1$ in a neighborhood of K . Then we can replace u and v by $\tilde{u} = \chi u$ and $\tilde{v} = \chi v$ in the right-hand side of (4.1). By local elliptic regularity we have that $\tilde{u} \in H^2(\mathbb{R}^{2d})$ and $\tilde{v} \in H^2(\mathbb{R}^{2d} \setminus \Gamma)$. However, the operator $\tilde{u} \mapsto \partial_N \tilde{u}|_\Gamma$ is compact as considered from $H^2(\mathbb{R}^{2d})$ to $L_2(\Gamma)$ and both $\tilde{v} \mapsto \tilde{v}_\Omega|_\Gamma$ and $\tilde{v} \mapsto \tilde{v}_K|_\Gamma$ are compact as considered from $H^2(\mathbb{R}^{2d} \setminus \Gamma)$ to $L_2(\Gamma)$, so it follows that V is compact.

4.2. Proof of Lemma 3.5

We start by showing that T_μ , originally defined in $L_2(\mathbb{R}^{2d})$, can be reduced to an operator in $L_2(\Gamma)$. More precisely we show that T_μ can be realized as an elliptic pseudodifferential operator of order 1 on some subspace of $L_2(\Gamma)$ of finite codimension, and hence there exists a constant $C > 0$ such that

$$\frac{1}{C} \|f\|_{L_2(\Gamma)} \|f\|_{H^1(\Gamma)} \leq \langle f, T_\mu f \rangle \leq C \|f\|_{L_2(\Gamma)} \|f\|_{H^1(\Gamma)} \tag{4.2}$$

for all f in that subspace.

Let f and g belong to \mathcal{A} . Also, let $u = Rf$, $v = \tilde{R}g$, and $w = Rg$. We saw in (4.1) that

$$\langle f, Vg \rangle = \int_\Gamma \partial_N u \cdot \overline{(v_\Omega - v_K)} dS. \tag{4.3}$$

To go further we will introduce the Neumann-to-Dirichlet and Dirichlet-to-Neumann operators. Let $G_\rho(x, y)$ be as in (2.16). We start with the single- and double-layer integral operators, defined by

$$\begin{aligned}
\mathcal{A}\alpha(x) &= \int_{\Gamma} G_0(x-y)\alpha(y)dS(y), \quad x \in \mathbb{R}^{2d}, \\
\mathcal{B}\alpha(x) &= \int_{\Gamma} \partial_{Ny}G_0(x-y)\alpha(y)dS(y), \quad x \in \mathbb{R}^{2d} \setminus \Gamma, \\
A\alpha(x) &= \int_{\Gamma} G_0(x-y)\alpha(y)dS(y), \quad x \in \Gamma, \\
B\alpha(x) &= \int_{\Gamma} \partial_{Ny}G_0(x-y)\alpha(y)dS(y), \quad x \in \Gamma.
\end{aligned} \tag{4.4}$$

The last two operators are compact on $L_2(\Gamma)$, since, by Lemma 2.1, their kernels have weak singularities. Moreover, since the kernel G_0 has the same singularity as the Green kernel for the Laplace operator in \mathbb{R}^{2d} (see [18, Chapter 7, Section 11]), we have the following limit relations on Γ

$$\begin{aligned}
\mathcal{A}\alpha_K &= A\alpha_K, \\
\mathcal{B}\alpha_K &= \frac{1}{2}\alpha + B\alpha, \\
\mathcal{A}\alpha_\Omega &= A\alpha_\Omega, \\
\mathcal{B}\alpha_\Omega &= -\frac{1}{2}\alpha + B\alpha.
\end{aligned} \tag{4.5}$$

Using a Green-type formula for L in K we see that

$$\beta = \mathcal{B}\beta_K - \mathcal{A}(\partial_N\beta_K). \tag{4.6}$$

If we combine this with the limit relations (4.5), we get

$$\left(B - \frac{1}{2}I\right)\beta_K = A(\partial_N\beta_K), \quad \text{on } \Gamma. \tag{4.7}$$

A similar calculation for Ω gives

$$\left(B + \frac{1}{2}I\right)\beta_\Omega = A(\partial_N\beta_\Omega), \quad \text{on } \Gamma. \tag{4.8}$$

It seems natural to do the following definitions.

Definition 4.1. We define the Dirichlet-to-Neumann and Neumann-to-Dirichlet operators in K and Ω as

$$\begin{aligned}
(DN)_K &= A^{-1} \left(B - \frac{1}{2} I \right), \\
(ND)_K &= \left(B - \frac{1}{2} I \right)^{-1} A, \\
(DN)_\Omega &= A^{-1} \left(B + \frac{1}{2} I \right), \\
(ND)_\Omega &= \left(B + \frac{1}{2} I \right)^{-1} A.
\end{aligned} \tag{4.9}$$

Remark 4.2. The inverses above exist at least on a space of finite codimension. This follows from the fact that A is elliptic and B is compact.

Lemma 4.3. *The operator $(ND)_K - (ND)_\Omega$ is an elliptic pseudodifferential operator of order -1 .*

Proof. Using a resolvent identity, we see that

$$(ND)_K - (ND)_\Omega = \left(B + \frac{1}{2} I \right)^{-1} \left(B - \frac{1}{2} I \right)^{-1} A. \tag{4.10}$$

It follows from the asymptotic expansion of $G_0(x, y)$ in Lemma 2.1 and the fact that G_0 is a Schwartz kernel (again, see [18, Chapter 7, Section 11]) that A is an elliptic pseudodifferential operator of order -1 . Moreover the operator B is compact, so the other two factors are pseudodifferential operators of order 0 which do not change the principal symbol noticeably. \square

Let us now return to the expression of V . We have

$$\begin{aligned}
\langle f, Vg \rangle &= \int_\Gamma \partial_N u \cdot \overline{(v_\Omega - v_K)} dS \\
&= \int_\Gamma \partial_N u \cdot \overline{(v_\Omega - w + w - v_K)} dS \\
&= \int_\Gamma \partial_N u \cdot \overline{((ND)_\Omega(\partial_N(v_\Omega - w)) + (ND)_K(\partial_N(w - v_K)))} dS \\
&= \int_\Gamma \partial_N u \cdot \overline{(((ND)_K - (ND)_\Omega)(\partial_N w))} dS.
\end{aligned} \tag{4.11}$$

Since we are interested in T_μ and not V , we may assume that f and g belong to $\mathcal{L}_{\Lambda_\mu}$. Then $u = Rf = \Lambda_\mu^{-1} f$ and $w = Rg = \Lambda_\mu^{-1} g$. For such f and g we get

$$\langle f, Vg \rangle = (\Lambda_\mu)^{-2} \int_\Gamma \partial_N f \cdot \overline{(((ND)_K - (ND)_\Omega)(\partial_N g))} dS \tag{4.12}$$

or with the introduced operators above

$$\langle f, Vg \rangle = (\Lambda_\mu)^{-2} \int_{\Gamma} \overline{f \cdot ((DN)_K^* ((ND)_K - (ND)_\Omega) ((DN)_K g))} dS. \quad (4.13)$$

Moreover, $(DN)_K$ is an elliptic pseudodifferential operator of order 1. This follows from the identity $A(DN)_K = B - (1/2)I$, and the fact that A is an elliptic pseudodifferential operator of order -1 . It follows from (4.13) that T_μ is an elliptic pseudodifferential operator of order 1.

Next, we prove the inequality (3.10). Because of the projections, it is enough to show it for functions f in $\mathcal{L}_{\Lambda_\mu}$.

The lower bound. We prove that there exists a subspace $\tilde{\mathcal{S}} \subset \mathcal{L}_{\Lambda_\mu}$ of finite codimension such that the lower bound in (3.10) is valid for all $f \in \tilde{\mathcal{S}}$. Since $f \in \mathcal{L}_{\Lambda_\mu}$ we have $L_\mu f := (L - \Lambda_\mu)f = 0$ so f belongs to the kernel of the second-order elliptic operator L_μ . Let $\varphi = f|_{\Gamma}$. We study the problem

$$\begin{aligned} L_\mu f &= 0 \quad \text{in } K^\circ, \\ f &= \varphi \quad \text{on } \Gamma. \end{aligned} \quad (4.14)$$

Let $E(x, y)$ be the Schwartz kernel for L_μ . It is smooth away from the diagonal $x = y$. One can repeat the theory with the single- and double-layer potentials for L_μ and write the solution f in the case it exists.

Let B_μ be the double-layer operator evaluated at the boundary,

$$B_\mu \alpha(x) = \int_{\Gamma} \partial_{N_y} E(x, y) \alpha(y) dS(y), \quad x \in \Gamma. \quad (4.15)$$

The operator B_μ is compact, since the kernel $\partial_{N_y} E(x, y)$ has a weak singularity at the diagonal $x = y$. Thus there exists a subspace $\mathcal{S}_1 \subset L_2(\Gamma)$ of finite codimension such that the operator $(1/2)I + B_\mu$ is invertible on \mathcal{S}_1 . Hence, there exists a subspace $\tilde{\mathcal{S}} \subset \mathcal{L}_{\Lambda_\mu}$ of finite codimension where we have the representation formula

$$f(x) = \int_{\Gamma} \frac{\partial E(x, y)}{\partial \nu_y} \left(\left(\frac{1}{2}I + B_\mu \right)^{-1} \varphi \right) (y) dS(y), \quad x \in K^\circ \quad (4.16)$$

for all $f \in \tilde{\mathcal{S}}$. The inequality $\|f\|_{L_2(K_0)} \leq C\|f\|_{L_2(\Gamma)}$ follows easily from (4.16) for all such functions f .

Since we also have $\|f\|_{L_2(\Gamma)} \leq C\|f\|_{H^1(\Gamma)}$ the lower bound in (3.10) follows via the lower bound in (4.2).

The upper bound. By the upper bound in (4.2) it is enough to show the following inequalities

$$\|f\|_{L_2(\Gamma)} \|f\|_{H^1(\Gamma)} \leq C\|f\|_{H^{1/2}(K)} \|f\|_{H^{3/2}(K)} \leq C\|f\|_{H^2(K)}^2 \leq C\|f\|_{L_2(K_1)}^2. \quad (4.17)$$

However, the first inequality is just the Trace theorem, the second is the Sobolev-Rellich embedding theorem. We note that $L_\mu f = 0$, so the third inequality is a standard estimate for elliptic operators.

5. Spectrum of Toeplitz Operators in a Reinhardt Domain

In the case when K is a Reinhardt domain one can strengthen part (b) of Lemma 3.6. Assume that K° , the interior of K , is a Reinhardt domain. This means that $0 \in K^\circ$ and if $z \in K^\circ$, then the set

$$\{(w^1, \dots, w^d) \mid w^j = tz^j, t \in \mathbb{C}, |t| < 1\} \quad (5.1)$$

is a subset of K° . If the set

$$\log |K| = \{(y^1, \dots, y^d) \mid y^j = \log |z^j|, z \in K^\circ\} \quad (5.2)$$

is convex in the usual sense, then K° is said to be logarithmically convex, and K° is a domain of holomorphy. Denote by $V_K : \mathbb{R}^d \rightarrow \mathbb{R}$ the function defined by

$$V_K(x) = \sup_{y \in \log |K|} \langle x, y \rangle. \quad (5.3)$$

We denote by $J : \mathcal{F}_B^2 \rightarrow \widetilde{\mathcal{L}} := L_2(K, e^{-(B/2)|z|^2} dm(z))$ the embedding operator. The s -values $s_{\widehat{\kappa}}$, $\widehat{\kappa} \in \mathbb{N}^d$, of J coincide with the numbers

$$\left\{ \frac{\|z^{\widehat{\kappa}}\|_{\widetilde{\mathcal{L}}}^2}{\|z^{\widehat{\kappa}}\|_{\mathcal{F}_B^2}^2} \right\}_{\widehat{\kappa} \geq 0} \quad (5.4)$$

(we remind the reader of the notation of eigenvalues in Lemma 3.6). Unlike the case $d = 1$, see [10], it is natural to numerate the eigenvalues by the d -tuples $\widehat{\kappa} = (\kappa_1, \dots, \kappa_d)$, just as for the eigenvalues of the Laplace operator in the unit cube $[0, 1]^d$, where the eigenvalues are given by $(2\pi)^{-d} |\widehat{\kappa}|_2^2 = (2\pi)^{-d} (\kappa_1^2 + \dots + \kappa_d^2)$.

Lemma 5.1. *Let $d > 1$ and $\omega = \widehat{\kappa}/|\widehat{\kappa}|$. Then*

$$(\widehat{\kappa}! s_{\widehat{\kappa}})^{1/|\widehat{\kappa}|} \sim \frac{B}{2} \exp(2V_K(\omega))(1 + o(1)), \quad \text{as } |\widehat{\kappa}| \rightarrow \infty. \quad (5.5)$$

Proof. The denominator in (5.4) is easily calculated to be

$$\|z^{\widehat{\kappa}}\|_{\mathcal{F}_B^2}^2 = \left(\frac{2\pi}{B}\right)^d \left(\frac{2}{B}\right)^{|\widehat{\kappa}|} \widehat{\kappa}!. \quad (5.6)$$

For the numerator, we do estimations from above and below, as in [12]. First, note that

$$I_{\hat{\kappa}} = \|z^{\hat{\kappa}}\|_{\mathcal{L}}^2 = \int_{\log|K|} \exp(2\langle \hat{\kappa}, x \rangle) d\tilde{m}(x), \quad (5.7)$$

where $d\tilde{m}(x)$ is the transformed measure. It is clear that

$$I_{\hat{\kappa}} \leq \exp(2|\hat{\kappa}|V_K(\omega))m(K). \quad (5.8)$$

For the inequality in the other direction, fix $\delta > 0$. The hyperplane

$$\langle \hat{\kappa}, x \rangle = (1 - \delta)V_K(\hat{\kappa}) \quad (5.9)$$

cuts $\log|K|$ in two components. Let P_δ be the component for which the inequality $\langle \hat{\kappa}, x \rangle \geq (1 - \delta)V_K(\hat{\kappa})$ holds. Then we have

$$I_{\hat{\kappa}} \geq \int_{P_\delta} \exp(2|\hat{\kappa}|(1 - \delta)V_K(\omega)) d\tilde{m}(x) \geq C_\delta \exp(2|\hat{\kappa}|(1 - \delta)V_K(\omega)), \quad (5.10)$$

where $C_\delta = \int_{P_\delta} d\tilde{m}(x) > 0$. It follows that

$$\begin{aligned} (\hat{\kappa}!S_{\hat{\kappa}})^{1/|\hat{\kappa}|} &\leq \left(m(K) \left(\frac{B}{2\pi} \right)^d \right)^{1/|\hat{\kappa}|} \frac{B}{2} \exp(2V_K(\omega)), \\ (\hat{\kappa}!S_{\hat{\kappa}})^{1/|\hat{\kappa}|} &\geq \left(C_\delta \left(\frac{B}{2\pi} \right)^d \right)^{1/|\hat{\kappa}|} \frac{B}{2} \exp(2(1 - \delta)V_K(\omega)), \end{aligned} \quad (5.11)$$

from which (5.5) follows. □

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